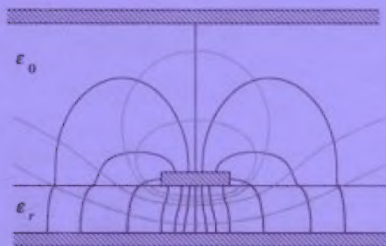


Transmission Line Design Handbook



Brian C. Wadell

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Preface

This book is a reference to designing transmission lines and designing with transmission lines. It has evolved from my interest in and desire to understand the design process and the many exceptions to the standard design procedures. As I explored this, I found that the standard equations were not valid for the standard configurations. For example, it is very common to have embedded microstripline and off-center stripline, yet the commonly used equations have significant errors for these configurations. What some may consider a mathematical exception, in fact is common and introduces significant errors in the design.

I found that the necessary references are spread out over dozens of texts and journals over seventy years. As a result, many engineers are using equations that are nearest at hand rather than the most correct available for their purpose. My goal became to find the most accurate and authoritative articles and extract the equations one needs for design. This seemed simple; however, it became a huge task of collecting references, analyzing them, and attempting calculations to verify their completeness and accuracy. The result, I think, is one of the most complete and easy-to-use references available.

This handbook compiles these equations and references in a uniform format for easy reference. It is meant to be a guide to making the calculations involved in high-frequency design and as such does not directly cover any EM field theory. There are several good EM texts listed in the references.

The equations collected herein are so-called closed-form equations. Many of these are exact solutions of the problem or of transformed versions of the problem. A few are the result of curve-fitting exact numerical solutions of Maxwell's equations. Where only curves are found in the references, CAE equations have been fit. These equations apply only in the ranges specified and are supplied for convenience in computer designs.

I have included material for the design of certain inductors and capacitors, which are not normally considered transmission lines. This is important material to have available when designing in the 200 to 500 MHz region where transmission line lengths may be unwieldy and off-the-shelf lumped elements self-resonate.

In most cases, the major source of error or uncertainty in designing with transmission lines is the fabrication process. I have gathered tables of the available

dielectrics and metals together with the expected tolerances. These can be used together with the closed-form equations to predict for yield.

Finally, I have included the special functions required for calculation of transmission lines. In the past, transmission line equations were often approximated for simple hand calculation or reduced to graphs, tables, and nomographs for common cases. However, with modern computers and calculators readily available to every engineer there is no reason to use anything other than the most accurate equations. Subroutines for complicated functions are relatively straightforward to write and libraries of these functions are readily available on many computers. Because of this, my philosophy has been to present the exact equations wherever available. The computer itself takes the place of the graphs, tables, and nomographs formerly required. This allows a great number of equations to be solved quickly to evaluate design trade-offs. The computer also does the calculation to the source accuracy without introducing any additional uncertainty due to interpolation of the table or reading of graphs. For many of the structures, the significant errors remaining will be attributable to fabrication process variations.

If you wish to write your own programs, you are now free to make any necessary trade-offs in accuracy and speed as suits your requirements and that of your computing platform. If your interest is solely in using the equations for design work, software is available to solve these equations.

As with any significant task, this book is the product of a number of people's contributions. I would like to acknowledge the assistance and encouragement of my friends and colleagues at Teradyne. In particular, Bobbi Bernard was instrumental in tracking down the references used in early versions of this book and Dan Rosenthal reviewed the draft. The anonymous reviewer and the editors and staff of Artech House, especially Mark Walsh, provided timely aid and encouragement. I would also like to thank the libraries of the Massachusetts Institute of Technology and Cornell University. My wife, Andrea C. Wadell, wrote software to check and solve the equations found in this book and to generate some of the more difficult figures. She also provided a great deal of moral support, which I greatly appreciate. Most of all, I would like to thank all of the authors referenced herein for the pleasure I have had in reading, studying, and using their articles and books.

Brian C. Wadell
Reading, MA
March 25, 1990

Chapter 1

Introduction

1.1 NOTATION USED IN THIS BOOK

An attempt has been made to be consistent with the notation throughout this book. This was done whenever multiple notations existed in the literature for consistency, clarity, and ease of comparison. In the process of doing this, the notation herein may be different than that of a referenced paper. This may seem a bit nitpicky to some, but anyone who has tried to compare several equations where the authors arbitrarily interchanged w and h , or pulled k out of the air, or referred to the relative dielectric constant as the dielectric constant (two completely different numbers!) will understand. The frequent figures are clearly labeled with these variables. In general, I have adopted the following conventions.

All dimensions are measured as labeled. Wherever possible these dimensions are specified as they would be measured on a given physical feature rather than a calculation from other dimensions or features. Most equations will work with any set of units (as long as all dimensions are the same unit). These are recognizable by the fact that they are expressed as ratios of dimensions. Where this is not true, I explicitly state the dimensions used and try to stick to SI units, which are used as summarized in Table 1.1.1 below.

MKS Units	Unit	Unit Symbol	Variables
length	meter	m	$l, d, D, w, t, h, b, \text{etc.}$
weight	kilogram	kg	
time	second	s	
current	amp	A	I, i
temperature	degrees Kelvin	°K	
frequency	hertz	Hz	f
power	watt	W	
charge	coulomb	C	q
resistance	ohm	Ω	R
conductance	mho, siemen	$1/\Omega, S$	G
resistivity	ohm-meter	$\Omega\text{-m}$	ρ
conductance	mho-meter, Siemen-meter	$\text{m}/\Omega, \text{S-m}$	σ
voltage	volt	V	V
capacitance	farad	F	C
inductance	henry	H	L
permittivity	farad / m	F/m	ϵ
permeability	henry / meter	H/m	μ

Table 1.1.1: MKS Units

The three-dimensional physical structure of the lines is represented as two-dimensional views either in the x-z plane as shown in Figure 1.1.1, or in the x-y plane as shown in Figure 1.1.2. The x-y view is used mainly for planar discontinuities. Top-view figures may have equations that implicitly reference the cross-sectional figure's dimensions. For example, a y-direction length dimension, l , is assumed for every physical transmission line structure.

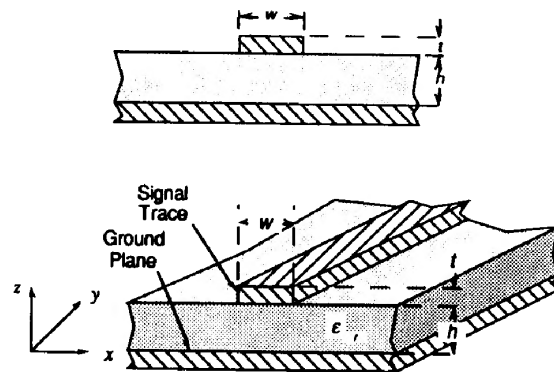


Figure 1.1.1: Cross-Section Figure Notation

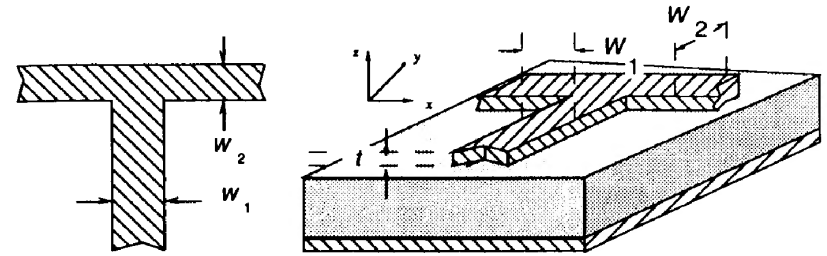


Figure 1.1.2: Top-View Figure Notation

The *dielectric constant* is $\epsilon = \epsilon_0 \epsilon_r$. The *relative dielectric constant* of the material is ϵ_r . The relative dielectric constant (ϵ_r) of air¹ is 1; the dielectric constant of air is ϵ_0 . Dielectrics are shown with a shaded (dotted) pattern. If several dielectrics are involved, each is shown with a different shading density. Air dielectric is shown without any pattern (white) and is labeled ϵ_0 or ϵ_{air} (as a reminder). The effective relative dielectric constant, ϵ_{eff} , is used to represent the effect of a mixed dielectric for quasi-TEM lines. Wherever possible ϵ_r and ϵ_{eff} are explicitly shown in the equations (some authors imply a $1/\sqrt{\epsilon_r}$ term in their equations for Z_0).

Dielectric thicknesses are h and b . The thickness of the microstrip line dielectric is h , and b is used for striplines. A single variable h could have been used for both, however, b is commonly used for striplines. For circular and square dielectrics, the corresponding dimension is D .

Center or signal conductor width and thickness are w and t , respectively. Conductors are drawn with a diagonal hatching pattern. For circular and square conductors, the corresponding dimension is d .

Coupled lines are separated by a spacing s . This is used strictly for distances between nongrounded conductors. A confusing situation occurs in the references for coplanar waveguide and similar structures where the width of the trace and its distance from the adjacent ground plane are both specified. In terms of fabrication, the relevant dimensions are the trace width and the spacing to the ground plane, however many references specify the trace width and the distance from the ground on one side of the trace to the ground on the other side of the trace. Some may even specify the centerline distances. In this book the latter a, b notation (e.g., Figure 3.4.1.1) is used, although I reserve the right to come up with a more convenient notation in the future.

The odd mode characteristic impedance is $Z_{0,o}$. That is, Z_0 (read "zee zero") subscript o for odd. Similarly $Z_{0,e}$ ("zee zero even") is the even mode characteristic

¹When an air dielectric is present in a structure and you wish to calculate the impedance for a similar structure with another dielectric ($\epsilon_r, 2$) replacing the air, make the following equation substitutions: replace ϵ_r with $\epsilon_r / \epsilon_r, 2$; replace η_0 with $\eta_0 / \epsilon_r, 2$.

impedance. Admittances are Y s with the same conventions. In the past (partly due to lack of computer-aided typesetting), some authors got sloppy and omitted, substituted, or interchanged zeroes and O s. (It really wouldn't matter if only it were consistent and clear.)

Coupling, k , is a unitless ratio. When coupling in dB is intended, k_{dB} is used for clarity. The dielectric constant will never be k .

The wavelength is represented by λ . The symbol λ should only be used for wavelength. It is confusing to use it as a function parameter, particularly if it should end up inside a hyperbolic function. There are usually two wavelengths involved; λ_0 is the wavelength without a dielectric (free-space wavelength) and λ_g is the guide wavelength (wavelength of the actual structure). We will avoid omitting the subscript unless the context is clear.

We have tried to use natural logarithms (logarithm base e) throughout for ease of comparison. Many equations appear in the references with the base 10 logarithms. The conversion to natural logarithms is quickly accomplished via:

$$\frac{\ln_e x}{\log_{10} x} = \ln_e 10 = 1 / \log_{10} e = 2.3026... \quad (1.1)$$

The logarithms are explicitly subscripted here for clarity, but this will not be done anywhere else.

Losses that are expressed in nepers can be converted to dB with the following relations:

$$\text{Loss} = \ln x \quad (\text{nepers}) \quad (1.2)$$

$$\text{Loss} = 20 \log_{10} x \quad (\text{dB}) = \frac{20 \ln x}{\log 10} = 8.6858 \ln x = 8.6858 \times \text{nepers} \quad (1.3)$$

For other mathematical functions the notation conventions of the *Handbook of Mathematical Functions* by Abramowitz and Stegun [1] are followed. This is most convenient because it is an excellent all-in-one reference containing definitions, equations, tables of values, graphs, and approximation formulas for virtually every special function we could encounter.

The notation commonly used for expressing impedance or admittance is:

$$Z = R + j X \quad \Leftrightarrow \quad \text{Impedance} = \text{Resistance} + j \text{Reactance} \quad (1.4)$$

$$Y = G + j B \quad \Leftrightarrow \quad \text{Admittance} = \text{Conductance} + j \text{Susceptance} \quad (1.5)$$

The impedance (admittance) is often normalized (divided by) a reference impedance (admittance) which is denoted Z_0 (Y_0). Impedance is expressed in Ohms (Ω), and Admittance ($1 / \text{Impedance}$) is expressed in Siemens (S) or mhos. When susceptances

or reactances are given in a reference they can be converted to inductance and capacitance with

$$L = \frac{-1.0}{\omega B_L} \quad (1.6)$$

$$L = \frac{X_L}{\omega} \quad (1.7)$$

$$C = \frac{B_C}{\omega} \quad (1.8)$$

$$C = \frac{-1.0}{X \omega} \quad (1.9)$$

1.2 HOW TO USE THIS BOOK

As a guide to using this book, I would just like to outline briefly the expected design procedure depicted in Figure 1.2.1.

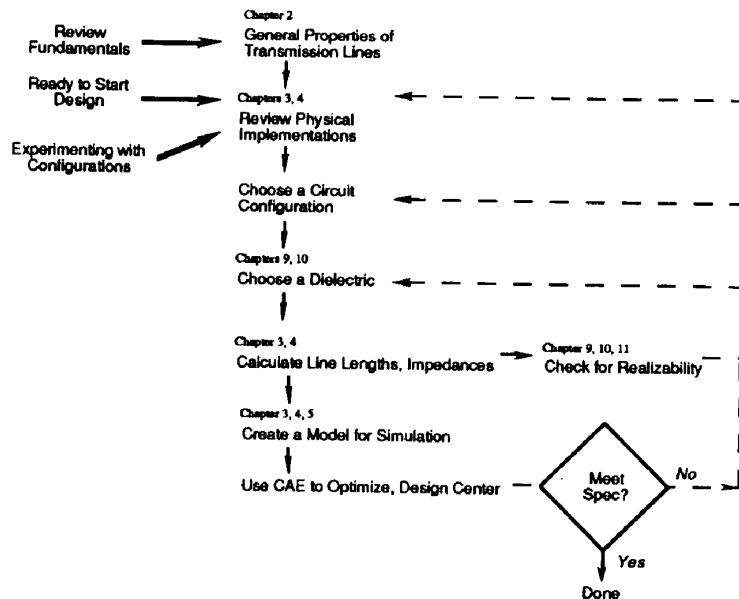


Figure 1.2.1 Design Flowchart

Generally, the designer starts with a general requirement such as “design a 15 dB coupler at 3 GHz.” Using the equations of Chapter 2, the design is converted to a pair of transmission lines with specified impedances and electrical lengths. We then look at *Chapter 4, Coupled Lines*, to get a feel for the physical implementations possible. Choose a substrate or dielectric (*Chapter 9, 10*) and try a few designs with the programs. Decide if the specs can be met with simple circuits or whether a cascaded or compensated design is required.

With this information, choose the type of line—microstrip line, stripline, coplanar waveguide, or other from *Chapter 4*. Synthesize the physical lines with equations from *Chapter 4*. Check for realizability of these dimensions using *Chapter 9, 10, 11* and iterate as necessary until the design is achievable. You may have to change to a different type of coupler or to a different physical realization.

Now using the dimensions and the physical layout add the discontinuities (*Chapter 5*) to your model. Enter the model into an analysis program (e.g., Touchstone™ SuperCompact™², etc.). Use the program’s optimization and yield analysis features to correct for the effect of the strays and center the design. Design centering assures

you that the values in use are the best for the expected parameter variations. Now you are ready to cut metal! At this point we may have designed dozens of transmission lines and tested them—all in software.

I have also encountered a second type of design process. In this case, the designer is familiar with the capability and limitations of a particular structure and is beginning a new design. Due to mechanical constraints of the system it is desired to make modifications to a standard structure. Electrically the requirements are relatively straightforward. This text and a transmission line program will allow the various alternatives to be tried without the delays of multiple prototypes.

Most designs will proceed in a similar fashion from spec to design, to physical realization, to layout with discontinuities, and finally to optimization and design centering. The use of computer tools allows this process to be fast and accurate and gives the designer the freedom to play “what if” with the goal of improving the design’s reproducibility.

REFERENCES

- [1] Abramowitz, Milton, and Irene A. Stegun, ed., *Handbook of Mathematical Functions with Formulas, Graphs, and Mathematical Tables*, Dover Publications, 1972.
- [2] ANSI/IEEE Std 1004-1987, “IEEE Standard Definitions of Planar Transmission Lines,” September 8, 1988.

²Touchstone is a registered trademark of EEsof, Inc. SuperCompact is a registered trademark of Compact Software, Inc.

Chapter 2

Generalized Transmission

Lines

There are only a few properties necessary to characterize completely the behavior of transmission lines for RF and microwave designs. When these are understood we will have an intuitive feel for their use in transmission line or *distributed element* circuits. The goal of this chapter is to develop this base.

2.1 DISTRIBUTED ELEMENT LINES

A transmission line circuit can be modeled in different ways. One particularly useful way is to classify a transmission line by the number of orthogonal physical dimensions, which are large relative to a wavelength in the transmission line (See Figures 1.1.1 and 1.1.2). Thus, a coaxial line might be considered one dimensional (y, its length), a right-angle bend in stripline might be considered two dimensional (x-y), and a transition from a coaxial line to a microstrip line would be considered three dimensional. A zero-dimensional transmission line is a lumped element. Remember that this classification is based on both the wavelength of the propagating signal and the discontinuity size. Depending on the frequency range and the accuracy desired, the choice of model may be made differently in seemingly similar circumstances.

No matter how many dimensions are actually significant, we will almost always try to reduce the problem to one of a circuit made up of RLC networks, ideal transformers, and ideal transmission lines. This type of calculation is usually much faster than the available field-solution techniques and it provides us with models in terms of more familiar concepts. All of the models presented in this book were developed from solutions of Maxwell's equations by various analytic tools. Unfortunately, field-solving programs can require hours to calculate results for even simple structures and have other limitations that restrict their use to special two- and three-dimensional problems. Field solvers are also not of much use when synthesizing a design. Fortunately, closed-form equations have been

developed by researchers from direct solutions and EM programs. Using these we have a highly accurate set of easily used models available for many designs. Certainly as hardware and software continues to become more powerful we will be hearing more about field-solving programs in the future.

A one-dimensional transmission line is modeled as a lumped element circuit with a very large number of elements as shown in Figure 2.1.1. The total line length, l , has been divided into an infinite number of sections of length Δy .

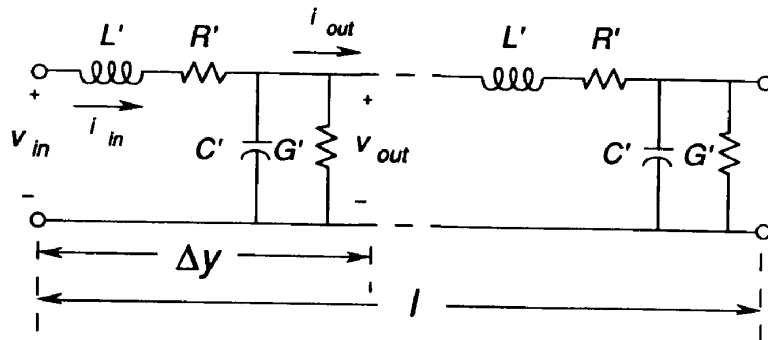


Figure 2.1.1: Transmission Line Distributed Parameters

Each lumped element represents an infinitesimal piece of the physical transmission line. The values L' , C' , R' , and G' are the inductance per length, capacitance per length, resistance per length, and conductance per length, respectively. These are the line's distributed parameters. For example, in a coaxial wire, L' represents the series inductance of the center conductor, C' represents the center conductor's capacitance to the shield, R' represents the resistance of the center conductor and shield, and G' represents the dielectric's leakage resistance. Table 2.1.1 below summarizes many of the contributions to the model parameters:

Model	Parameter	Origin	Comment
R	Conductance	Conductance	atomic property of metal used
		Geometry	larger cross section = lower R longer = higher R
	Radiation	Radiation	fields not contained are losses
		Skin Effect	current flows in smaller cross-section as frequency increases
	Proximity Affect	Proximity Affect	adjacent conductor's field forces current into smaller cross-section
L	Permeability	Permeability	property of metal used
		Geometry	larger cross section = lower L longer = higher L
	Skin Effect	Skin Effect	currents at different depths in conductor have different phases, net effect is inductive
		Proximity	adjacent conductor's field forces current into smaller cross-section
C	Geometry	Radiation	larger area = higher C
		Geometry	closer conductors = higher C
G	Permittivity	Permittivity	higher ϵ_r = higher C
	Loss Tangent	Loss Tangent	higher loss tan = higher losses in dielectric
	Conductance	Conductance	semiconductor dielectrics
	Geometry	Geometry	

Table 2.1.1: Transmission Line Model Parameters and Sources

2.2 WAVE PROPAGATION AND CHARACTERISTIC IMPEDANCE

In this section we show the solution of wave propagation in the distributed element line. It is very informative to understand how the simple model in Figure 2.1.1 creates the wave propagation equations and a definition of characteristic impedance; however, one can successfully undertake designs and make use of most of this book even if this derivation is skipped.

The line we will analyze is shown in Figure 2.1.1. The length of the line is the y dimension. By inspection, and replacing differences with differentials (by taking the limit as Δy approaches zero) we can write equations for the signal on our infinitesimal section of line as:

$$\frac{\partial}{\partial y} = -v G' - \frac{\partial v_{out}}{\partial t} C' \quad (2.2.1)$$

$$\frac{\partial v}{\partial y} = -i R' - L' \frac{\partial i}{\partial t} \quad (2.2.2)$$

That is, each section ∂y of the line creates the changes in the voltage and current described above. Equations (2.2.1) and (2.2.2) are referred to as the *Telegrapher's equations*.

If we differentiate each of these equations with respect to t and y , we obtain:

$$\frac{\partial^2 v}{\partial y^2} = -R' \frac{\partial i}{\partial y} - L' \frac{\partial^2 i}{\partial y \partial t} \quad (2.2.3)$$

$$\frac{\partial^2 v}{\partial y \partial t} = -R' \frac{\partial i}{\partial t} - L' \frac{\partial^2 i}{\partial t^2} \quad (2.2.4)$$

$$\frac{\partial^2 i}{\partial y \partial t} = -G' \frac{\partial v}{\partial t} - C' \frac{\partial^2 v}{\partial t^2} \quad (2.2.5)$$

$$\frac{\partial^2 i}{\partial y^2} = -G' \frac{\partial v}{\partial y} - C' \frac{\partial^2 v}{\partial y \partial t} \quad (2.2.6)$$

We can now combine equations to eliminate variables. Substituting (2.2.1) and (2.2.5) into (2.2.3) and (2.2.2), and (2.2.4) into (2.2.6):

$$\begin{aligned} \frac{\partial^2 v}{\partial y^2} &= -R' \left(-v G' - \frac{\partial v}{\partial t} C' \right) - L' \left(-G' \frac{\partial v}{\partial t} - C' \frac{\partial^2 v}{\partial t^2} \right) \\ &= v (R' G') + \frac{\partial v}{\partial t} (R' C' + L' G') + \frac{\partial^2 v}{\partial t^2} (L' C') \end{aligned} \quad (2.2.7)$$

$$\begin{aligned} \frac{\partial^2 i}{\partial y^2} &= -G' \left(-i R' - L' \frac{\partial i}{\partial t} \right) - C' \left(-R' \frac{\partial i}{\partial t} - L' \frac{\partial^2 i}{\partial t^2} \right) \\ &= i (G' R') + \frac{\partial i}{\partial t} (L' G' + C' R') + \frac{\partial^2 i}{\partial t^2} (L' C') \end{aligned} \quad (2.2.8)$$

Now we have two differential equations that describe our transmission line, but can we gain understanding of the wave's propagation from it? (Much of the following discussion will be for v , but the reasoning for i is parallel.) Examining (2.2.7) we see an equation for a voltage that is a function of its time derivatives. Any function v that is a solution of (2.2.7) describes a possible method of propagation on our line. There may be many such

equations, so how can we proceed? An equation that is a solution and is fortunately quite general is the complex exponential function, which has derivatives [note the similarities to (2.2.7) and (2.2.8)]:

$$\frac{\partial e^u}{\partial x} = e^u \frac{\partial u}{\partial x} \quad (2.2.9)$$

$$\frac{\partial^2 e^u}{\partial x^2} = e^u \frac{\partial^2 u}{\partial x^2} + e^u \left(\frac{\partial u}{\partial x} \right)^2 = e^u \left[\frac{\partial^2 u}{\partial x^2} + \left(\frac{\partial u}{\partial x} \right)^2 \right] \quad (2.2.10)$$

We also have one other problem: v is a function of y and t . To solve this, we make the assumption that we can separate the variables in our solution. This means that we can write our function v as the product of two functions v' and v'' . Similarly, we rewrite i and obtain

$$v(y, t) = v'(y) v''(t) \quad (2.2.11)$$

$$i(y, t) = i'(y) i''(t) \quad (2.2.12)$$

Now we are ready to solve (2.2.7) and (2.2.8). We choose

$$v(y, t) = v'(y) e^{j \omega t} \quad (2.2.13)$$

$$\frac{\partial v}{\partial t} = j \omega v(y, t) \quad (2.2.14)$$

$$\frac{\partial^2 v}{\partial t^2} = -\omega^2 v(y, t) \quad (2.2.15)$$

$$i(y, t) = i'(y) e^{j \omega t} \quad (2.2.16)$$

$$\frac{\partial i}{\partial t} = i(y, t) j \omega \quad (2.2.17)$$

$$\frac{\partial^2 i}{\partial t^2} = -\omega^2 i(y, t) \quad (2.2.18)$$

Rewriting (2.2.1) and (2.2.2) with (2.2.13) and (2.2.14)

$$\frac{\partial v}{\partial y} = v(y, t) (G' + j \omega C') \quad (2.2.19)$$

$$\frac{\partial v}{\partial y} = -i(y, t) (R' + j \omega L') \quad (2.2.20)$$

and substituting into (2.2.7)

$$\begin{aligned} \frac{\partial^2 v}{\partial y^2} &= v(R'G' + j\omega v(y, t)(R'C' + L'G') - \omega^2 v(y, t)(L'C')) \\ &= v(y, t) [R'G' + j\omega(R'C' + L'G') - \omega^2 L'C'] \\ &= v(y, t) [(R + j\omega L)(G + j\omega C)] \end{aligned} \quad (2.2.21)$$

Again we recognize that this is an equation whose solution is of the exponential form:

$$v(y) = K_0 e^{-\gamma y} + K_1 e^{+\gamma y} \quad (2.2.22)$$

$$\gamma = \sqrt{(R' + j\omega L')(G' + j\omega C')} \quad (2.2.23)$$

It is worth examining Equation (2.2.22) further. First we see that there are two pieces of the solution for generality. Both need not exist, however, they may. Each term represents a sinusoidal wave propagating in the y or $-y$ direction. This means that our lumped-element line will support propagation in either direction alone or both directions simultaneously. Each wave may have its own amplitude. This means that the voltage on the line at location y will be the sum of two sine waves having the same period in distance.

We can substitute (2.2.22) and (2.2.23) into our rewritten (2.2.2)

$$-\gamma K_0 e^{-\gamma y} + \gamma K_1 e^{+\gamma y} = -i(y, t) (R' + j\omega L')$$

$$i(y, t) = \frac{-\gamma K_0 e^{-\gamma y} + \gamma K_1 e^{+\gamma y}}{(R' + j\omega L')} \quad (2.2.24)$$

The ratio of v to i or the impedance at a point y on the line is

$$\frac{v(y, t)}{i(y, t)} = Z_0 = \sqrt{\frac{R' + j\omega L'}{G' + j\omega C'}} \quad (2.2.25)$$

This is known as the **characteristic impedance** of the transmission line. Again, notice that for our line this is not a function of the propagating waveform, the position on the line, or time. For this uniform line, it is a function of the model parameters only. In general, these parameters may be a function of y which would make Z_0 a function of the position in the line. If $R' = G' = 0$, a lossless line, Z_0 is not a function of frequency either.

$$Z_0 = \sqrt{\frac{L'}{C'}} \quad (2.2.26)$$

In the above equations we defined a parameter γ . This is the **propagation constant**

$$\gamma = \alpha + j\beta = \sqrt{j\omega\mu(\sigma + j\omega\epsilon)} \quad (2.2.27)$$

which is made up of the **attenuation constant**, α ,

$$\alpha = \omega \sqrt{\frac{\mu\epsilon}{2} \left(\sqrt{1.0 + \frac{\sigma^2}{\omega^2\epsilon^2}} - 1.0 \right)} \quad (\text{Np/m}) \quad (2.2.28)$$

and the **phase constant**, β

$$\beta = \omega \sqrt{\frac{\mu\epsilon}{2} \left(\sqrt{1.0 + \frac{\sigma^2}{\omega^2\epsilon^2}} + 1.0 \right)} \quad (\text{rad/m}) \quad (2.2.29)$$

The phase constant is sometimes written as k . We will avoid this usage to prevent confusion with a number of other k 's that occur in transmission line problems, *i.e.*, the modulus of the complete elliptic integral, the dielectric constant, and coupling. The attenuation constant can be written in dB/m by multiplying by the factor, $20/\ln 10$.

2.3 MAXWELL'S EQUATIONS

Although we have described our transmission line as a lumped equivalent, this solution happily corresponds to the solution of Maxwell's equations. We will not discuss Maxwell's equations in depth, however, for completeness, in differential form they are:

$$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t} \quad \mathbf{D} = \epsilon \mathbf{E}$$

$$\nabla \times \mathbf{H} = \mathbf{J} + \frac{\partial \mathbf{D}}{\partial t} \quad \mathbf{B} = \mu \mathbf{H}$$

$$\nabla \cdot \mathbf{D} = q_e \quad \mathbf{J} = \sigma \mathbf{E}$$

$$\nabla \cdot \mathbf{B} = q_m$$

(2.3.1)

where

$$\nabla \times \mathbf{V} = \left(\frac{\partial V_z}{\partial y} - \frac{\partial V_y}{\partial z} \right) \hat{x} + \left(\frac{\partial V_x}{\partial z} - \frac{\partial V_z}{\partial x} \right) \hat{y} + \left(\frac{\partial V_y}{\partial x} - \frac{\partial V_x}{\partial y} \right) \hat{z} = \text{curl } \mathbf{V}$$

$$= \begin{bmatrix} \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ V_x & V_y & V_z \\ \hat{x} & \hat{y} & \hat{z} \end{bmatrix} \quad (2.3.2)$$

$$\mathbf{u} = \begin{bmatrix} \hat{x} \\ \hat{y} \\ \hat{z} \end{bmatrix} = \text{unit vector} \quad (2.3.3)$$

where \mathbf{E} and \mathbf{H} are the electric and magnetic field intensities, and \mathbf{D} and \mathbf{B} are the electric and magnetic flux densities. The current density is \mathbf{J} , and q_e and q_m are the electric and magnetic charge densities. The magnetic current density is \mathbf{M} . All bold-faced variables are functions of x, y, z , and t .

In solving Maxwell's equations for a lossless TEM line and equating them to the distributed parameter solution, we find other relations in terms of the parameters of the dielectric and the conductors. The characteristic impedance is

$$Z_0 = \frac{\sqrt{\mu_0 \epsilon_0 \epsilon_r}}{C} \quad (\Omega) \quad (2.3.4)$$

$$Z_0 = \frac{L}{\sqrt{\mu_0 \epsilon_0 \epsilon_r}} \quad (\Omega) \quad (2.3.5)$$

By combining (2.3.4) and (2.3.5) we define the **wave impedance**:

$$\eta = \sqrt{\frac{\mu}{\epsilon}} \quad (\Omega) \quad (2.3.6)$$

The **characteristic impedance of free space** is:

$$\eta_0 \equiv \sqrt{\frac{\mu_0}{\epsilon_0}} = 120 \pi \quad (\Omega) \equiv 377 \quad (\Omega) \quad (2.3.7)$$

The **speed of light** is the propagation velocity or phase velocity in free space (a vacuum) and is defined as:

$$c = \frac{1.0}{\sqrt{\mu_0 \epsilon_0}} = 2.997925 \times 10^8 \quad (\text{m/s}) \quad (2.3.8)$$

From this, we define the **phase velocity** or the velocity of light in the medium:

$$v_p = \frac{1.0}{\sqrt{\mu \epsilon}} = \frac{1.0}{\sqrt{\mu_0 \mu_r \epsilon_0 \epsilon_r}} = \frac{c}{\sqrt{\mu_r \epsilon_r}} \quad (2.3.9)$$

In most materials $\mu_r = 1$ (nonmagnetic), so we often write (2.3.9) as:

$$v_p = \frac{c}{\sqrt{\epsilon_r}} \quad (2.3.10)$$

The phase velocity can be used to calculate the wavelength in the medium, also known as the **guide wavelength**:

$$\lambda_g = \frac{c}{f \sqrt{\mu_r \epsilon_r}} = \frac{\lambda_0}{\sqrt{\mu_r \epsilon_r}} \quad (2.3.11)$$

where f is the frequency of the wave and λ_0 is the wavelength in free space. Again for nonmagnetic media we can simplify:

$$\lambda_g = \frac{c}{f \sqrt{\epsilon_r}} \quad (2.3.12)$$

and in free space:

$$\lambda_0 = \frac{c}{f} \quad (2.3.13)$$

For a nonhomogenous, nonmagnetic dielectric,

$$\lambda_g = \frac{c}{f \sqrt{\epsilon_{eff}}} = \frac{\lambda_0}{\sqrt{\epsilon_{eff}}} \quad (2.3.14)$$

The **electrical length** θ of a transmission line having physical length l at frequency f is

$$\begin{aligned} \theta &= 2.0 \pi \frac{l}{\lambda_g} \quad (\text{radians}) = 360.0 \frac{l}{\lambda_g} \quad (\text{degrees}) \\ &= 2.0 \pi l f \frac{\sqrt{\epsilon_{eff}}}{c} \quad (\text{radians}) = 360.0 l f \frac{\sqrt{\epsilon_{eff}}}{c} \quad (\text{degrees}) \end{aligned} \quad (2.3.15)$$

2.3.1 Example: Calculating Impedance from Capacitance

We are given the incremental capacitance of a coaxial line as:

$$C = \frac{2 \pi \epsilon_0 \epsilon_r}{\ln D / d}$$

What is the equation for Z_0 ?

Solution:

Since the dielectric is homogenous and the shield surrounds the center conductor, we know that $\epsilon_{eff} = \epsilon_r$. Using (2.3.4):

$$Z_0 = \frac{\sqrt{\mu_0 \epsilon_0 \epsilon_r \ln D / d}}{2 \pi \epsilon_0 \epsilon_r} = \frac{\sqrt{\mu_0 \ln D / d}}{2 \pi \sqrt{\epsilon_0 \epsilon_r}} = \boxed{\frac{\eta_0 \ln D / d}{2 \pi \sqrt{\epsilon_r}}}$$

2.4 MATERIAL PROPERTIES AND LUMPED MODEL

As we saw in the previous section, there are several properties of the material which are important to signal propagation. The properties of the medium in which we are propagating are expressed in terms of the *permeability*:

$$\mu = \mu_0 \mu_r \quad (2.4.1)$$

the *permittivity*,

$$\epsilon = \epsilon_0 \epsilon_r \quad (2.4.2)$$

and the *conductivity*.

$$\sigma = \frac{1.0}{\rho} \quad (2.4.3)$$

Conductivity is also encountered as its reciprocal, *resistivity*, ρ . The permeability and permittivity of free space have the values:

$$\mu_0 \equiv 4.0 \pi \times 10^{-7} \text{ (H/m)} \quad (2.4.4)$$

$$\epsilon_0 \equiv 8.854 \times 10^{-12} \text{ (F/m)} \equiv \frac{1.0}{36\pi \times 10^{-9}} \text{ (F/m)} \quad (2.4.5)$$

The use of μ_0 and ϵ_0 allows us to express all of our equations in terms of the unitless relative permeability and permittivities, μ_r and ϵ_r , for convenience. Resistivity is often similarly normalized to the resistivity of copper or gold.

The specific values of ϵ_r , μ_r , and σ are determined by interactions at the atomic level. For example, the conductivity is a function of the mobility and the number of free charge carriers (electrons and holes). In nonconductive materials the charges are bound or locked in position. They cannot move freely when a field is applied. The permittivity¹ is a measure of the ability to store energy by shifting the center of the bound charges against the atomic forces (*polarizing* the material) to cancel the applied field. There are three mechanisms for this: dipole, ionic, and electronic. Each of these can create resonance-like behavior in the values of ϵ_r and $\tan \delta$ at high frequencies. Similarly, the permeability is a measure of the ability of the material to store magnetic energy. An impressed magnetic field stores energy in the material by displacing the magnetic moments of the material. Again, there are several atomic mechanisms that can be involved: electric loops, spin, and nuclear spin.

¹The *index of refraction* is the square root of the permittivity.

Values of relative permeability range from slightly below one to as high as a million (supermalloy). It is a measure of the ability of the medium to store energy in the magnetic field. It is frequency dependent and has both a real and imaginary component. Permeabilities of some metals are tabulated in Chapter 9.

Relative permittivity is 1.0 for free space and can be as high as 6000. It is also a function of frequency with both real and imaginary components. Permittivities are tabulated in Chapters 7 and 9.

Resistivity varies over a very wide range (10^{25}). It is a function of frequency. Resistivities are tabulated for metals in Chapter 9.

2.4.1 Complex Permeability and Permittivity

Permeability and permittivity are complex numbers so we write

$$\epsilon_r^* = \epsilon_r' + j\epsilon_r'' \quad (2.4.1.1)$$

$$\mu_r^* = \mu_r' + j\mu_r'' \quad (2.4.1.2)$$

The real part represents the ability of the medium to store electric and magnetic energy while the imaginary part represents losses. Usually ϵ_r'' , and μ_r'' are small so we drop the primed and starred notations when referring to the real part.

The meaning of the complex values of permittivity and permeability can be seen by comparison with the more common power factor. The power factor is defined by circuit theory to be the ratio of the average power dissipated to the voltage \times current product. This is the cosine of the angle between the actual or real power dissipated to the instantaneous VI product as shown in Figure 2.4.1.

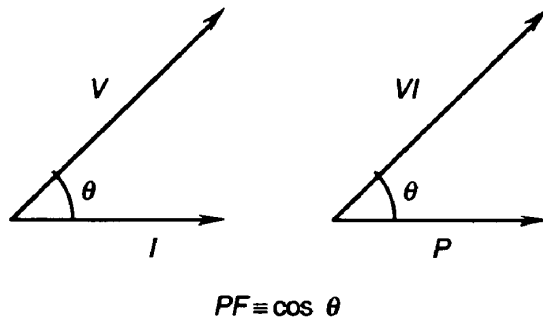


Figure 2.4.1: Power Factor

For dielectric and magnetic losses, we define a similar and related function, the *loss tangent*. The loss tangent is depicted in Figure 2.4.2.

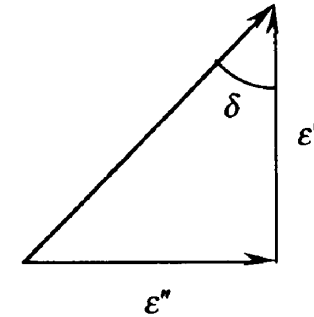


Figure 2.4.2: Loss Tangent

In both Figures 2.4.1 and 2.4.2 the horizontal axis represents lost energy and the vertical axis represents stored energy. The power factor of the dielectric and the loss tangent are related ways of expressing the material's losses. For small power factors (less than 0.15, which is generally the case for dielectrics we use), the loss tangent approximates the power factor within 1%. The relationships are:

$$PF = \text{power factor} = \sin \left(\frac{\text{energy loss per cycle}}{\text{energy stored per cycle}} \right) = \sin [\tan^{-1} D] \quad (2.4.8)$$

$$D = \text{dissipation factor} = \frac{1.0}{Q} = \frac{\epsilon''}{\epsilon'} + \frac{\sigma}{\omega \epsilon'} \quad (2.4.9)$$

For a dielectric, the last term is much less than one so

$$D \equiv \tan \delta \quad (2.4.10)$$

and for small D

$$PF = \sin [\tan^{-1} D] \approx \sin [\sin^{-1} D] = D \equiv \tan \delta \quad (2.4.11)$$

When we use ϵ_r we will mean ϵ_r' , and similarly we replace μ_r' with the notation μ_r . The imaginary parts are then represented by the dielectric and magnetic loss tangents

$$\epsilon_r'' \equiv \epsilon_r' \tan \delta_d \quad (2.4.12)$$

$$\mu_r'' \equiv \mu_r' \tan \delta_m \quad (2.4.13)$$

These are referred to as dielectric or magnetic “loss tan” or “tan delta.” You will most often see $\tan \delta$ used without any subscript where the dielectric loss tan is implicit since $\tan \delta_m$ is less often encountered in actual design.

2.5 MATERIAL CLASSIFICATIONS

The parameters we have defined create a number of natural classifications. These classes allow us to simplify our analysis of transmission line characteristics. Although we often assume that the materials making up our transmission lines have constant parameters, this is not true in all cases.

If $[\sigma / (\omega \epsilon)]^2 \gg 1.0$ we have a **conductor**. When $[\sigma / (\omega \epsilon)]^2 \ll 1.0$ the material is a **dielectric**. Semiconductors lie in between. **Superconductors** are conductors for which $\sigma = \infty$.

The permittivity, conductivity, and permeability are all functions of frequency and direction in the medium. When these parameters vary with frequency, the phase velocity also varies with frequency and the line is said to have **dispersion**. A dispersionless line has constant signal delay over frequency. Such a line propagates wideband signals (e.g., pulses) without distortion.

Woven dielectrics are **anisotropic**; they will have slightly different ϵ_r values in different directions within the material. We also encounter **nonhomogeneous** dielectrics in semiconductor dielectrics, multilayer printed circuit boards, and in microstrip line. A nonhomogeneous dielectric has an ϵ_r that is a function of position (x, y, z) in the line.

A material may also be **nonlinear**—that is, ϵ_r or μ_r is a function of any applied fields. Permeability is commonly a function of the previous magnetization of the material.

2.6 QUASI-TEM TRANSMISSION LINE PARAMETERS

Inhomogeneous dielectrics occur when we have several dielectrics or when the ϵ_r , μ_r varies with position in the dielectric. In this case, the wave is not propagating in a strictly TEM fashion. However, it is convenient and useful to assume that for low frequencies we can substitute an **effective** parameter calculated from the geometry and the parameters of the individual regions making up the medium. These parameters—the effective permeability, μ_{eff} , and effective permittivity, ϵ_{eff} —are relative numbers, like μ_r and ϵ_r .

For structures with mixed dielectrics (e.g., air and substrate as in microstrip line) we define the **effective relative dielectric constant**:

$$\epsilon_{eff} = \frac{C_{actual}}{C_{air}} \quad (2.6.1)$$

which is the ratio of the actual structure’s capacitance (Figure 2.6.1a) to the capacitance when the dielectric is replaced with air (Figure 2.6.1b). This effective dielectric constant of a homogeneous structure (Figure 2.6.1c) is used wherever ϵ_r would be used in a

homogenous dielectric calculation. This definition allows us to treat propagation as propagation in an equivalent or **quasi-TEM** line.

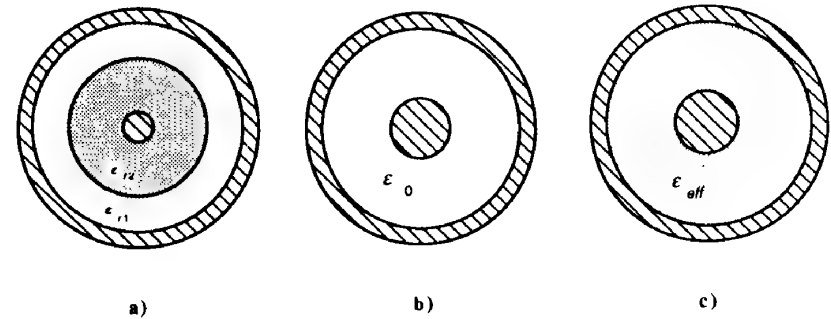


Figure 2.6.1: Effective Dielectric Constant

Another commonly used number is the **filling factor**, q :

$$q = \frac{\epsilon_{eff} - 1}{\epsilon_r - 1} = \begin{cases} 0 & \text{if all air} \\ 1 & \text{if all dielectric} \end{cases} \quad (2.6.2)$$

or rearranging:

$$\epsilon_{eff} = q(\epsilon_r - 1) + 1 \quad (2.6.3)$$

2.6.1 Example: Time Delay, Electrical Length and ϵ_{eff}

What is the time delay of a 3 cm length of microstrip line having an ϵ_{eff} of 3.1? What is the phase delay of this line at the fundamental and first two harmonics of a 500 MHz square wave?

Solution: We can calculate the phase velocity in the medium with (2.3.10):

$$v_p = \frac{3.0 \times 10^8 \text{ m/s}}{\sqrt{3.1}} = 1.7039 \times 10^8 \text{ m/s}$$

so the time delay between input and output is

$$t = \frac{3 \text{ cm} \times 1 \text{ m} / 100 \text{ cm}}{1.7039 \times 10^8 \text{ m/s}} = 176 \text{ ps}$$

and we calculate the guide wavelengths of the harmonics as (2.3.12):

$$\lambda_{g1} = \frac{3.0 \times 10^8 \text{ m/s}}{500 \times 10^6 \sqrt{3.1}} = 0.3407 \text{ m}$$

$$\lambda_{g2} = \frac{3.0 \times 10^8 \text{ m/s}}{1000 \times 10^6 \sqrt{3.1}} = 0.1704 \text{ m}$$

$$\lambda_{g3} = \frac{3.0 \times 10^8 \text{ m/s}}{1500 \times 10^6 \sqrt{3.1}} = 0.1136 \text{ m}$$

Thus, the 3 cm represents a phase delay ϕ_n at the n th harmonic of

$$\phi_1 = 360^\circ \times .03 \text{ m} / .3407 \text{ m} = 31.6994^\circ$$

$$\phi_2 = 360^\circ \times .03 \text{ m} / .1704 \text{ m} = 63.3803^\circ$$

$$\phi_3 = 360^\circ \times .03 \text{ m} / .1136 \text{ m} = 95.0704^\circ$$

2.7 TRANSMISSION LINE LOSSES

The total losses in a transmission line are described by:

$$\alpha = \alpha_c + \alpha_d + \alpha_r + \alpha_l \quad (2.7.1)$$

where

$$\alpha_c = \text{conductor losses} \quad (2.7.2)$$

$$\alpha_d = \text{dielectric losses} \quad (2.7.3)$$

$$\alpha_r = \text{radiation losses} \quad (2.7.4)$$

$$\alpha_l = \text{leakage losses} \quad (2.7.5)$$

Losses are converted to and from series resistances with the equation

$$\alpha \equiv \frac{R_{\text{series}}}{2.0 Z_0} \quad (2.7.6)$$

which is valid for small R_{series} .

2.7.1 Conductor Losses

A number of variables and concepts commonly occur when calculating conductor losses. For a complete derivation of the following equations see, *e.g.*, Ramo *et al.* [17]. The dc resistance of a conductor is calculated with

$$R_{dc} = \frac{\rho l}{A} \quad (\Omega) \quad (2.7.1.1)$$

where ρ is the resistivity, l is the conductor length, and A is the cross-sectional area.

The internal impedance of the conductor is a function of frequency

$$Z_s = R_s + j \omega L_i \quad (2.7.1.2)$$

where

$$R_s = \frac{1.0}{\sigma \delta} \quad (2.7.1.3)$$

$$= \sqrt{\frac{\pi f \mu}{\sigma}} = \text{surface resistivity or sheet resistivity} \quad (\Omega / \text{square})$$

$$L_i = R_s / \omega \quad (2.7.1.4)$$

and

$$\delta = \frac{1.0}{\sqrt{\omega \mu \sigma}} = \text{skin depth} \quad (2.7.1.5)$$

A *normalized frequency* is defined as

$$P = \sqrt{2.0 \mu \sigma f A} \quad (2.7.1.6)$$

The variable R_s is the *skin effect* resistance. As frequency increases our conductor losses also increase. The skin depth is the depth at which the current density is reduced to $1/e$. It can also be thought of as the equivalent thickness of a conductor at dc having the same resistance. The skin depth for several conductors is shown in Figure 2.7.1.1.

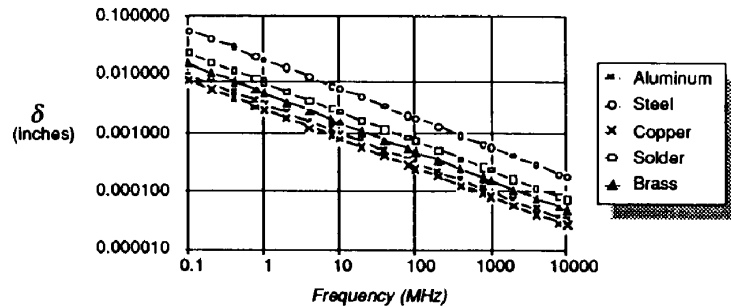


Figure 2.7.1.1: Skin Depth in Conductors

The magnitude of the real and imaginary parts of the internal impedance are equal at all frequencies and the phase angle of Z_s is therefore a constant 45° . This fact is the origin of the **incremental inductance** rule used to calculate conductor losses. We can use the incremental inductance rule to calculate the conductor losses of any transmission line when

- losses are low
- losses are due to skin effect
- internal inductance is small
- conductor thickness \gg skin depth
- conductor radii \gg skin depth

Using the equivalence of (2.7.7) we can calculate the loss resistance as

$$R = \omega [L (\text{nominal dimensions}) - L (\text{conductors reduced by } \delta/2)] \quad (2.7.1.7)$$

The meaning of the inductances is shown in Figure 2.7.1.2.

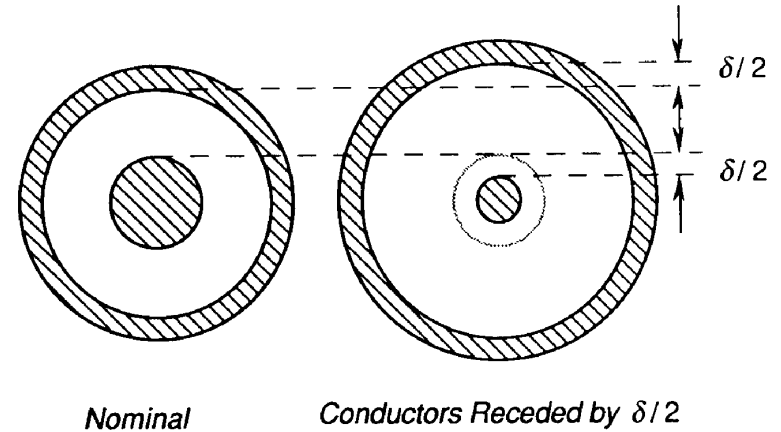


Figure 2.7.1.2: Incremental Inductance Defined

The presence of the dielectric has no effect on the inductance. Because we can use this relation for virtually any transmission line for which we have an expression for the inductance or impedance, it will be of great value in loss calculations.

2.7.2 Dielectric Losses

Dielectric losses are created by the atomic mechanisms described earlier. Inhomogeneous (in the x-z plane of Figure 1.1.1 and 1.1.2) dielectric losses are reduced by the presence of the (lossless) air in the cross-section. The loss of these structures is modified by a dielectric loss filling factor [10]

$$q_{\tan \delta} = \frac{\epsilon_{eff} - 1.0}{\epsilon_{eff} - \epsilon_{eff} / \epsilon_r} \quad (2.7.2.1)$$

and

$$\alpha_d = \frac{\pi f q_{\tan \delta} \tan \delta \sqrt{\epsilon_{eff}}}{c_0} \quad (2.7.2.2)$$

2.7.3 Radiation Losses and Leakage Losses

Signal power which is radiated into space or to another conductor is not available in the load. This loss is heavily dependent on the physical structure. Discontinuities tend to increase the radiation losses. Usually the design keeps these losses small and they are

neglected. Some structures (e.g., open ends) have been characterized for calculation purposes.

Leakage losses are present in transmission lines having dielectrics with finite resistivity. An example of this is transmission lines fabricated in MMICs which have semiconductor material as the dielectric. The propagation is complicated in this situation; however, Hoffman [10] gives an approximation useful in some cases.

When the electric field lines lie along (or approximately along) the dielectric boundary the following may be used

$$\alpha_l = 1.64 \times 10^3 \sigma \frac{\sqrt{\epsilon_{eff}} - 1 / \sqrt{\epsilon_{eff}}}{\epsilon_r - 1.0} \quad (\text{dB/cm}) \quad (2.7.3.1)$$

where conductivity is in $(\Omega\text{-cm})^{-1}$.

2.8 TRANSMISSION LINE CALCULATIONS

The solution of Telegrapher's equations found above is only one possible form. Another useful form is expressed in terms of hyperbolic functions as shown in the following sections. A number of calculations making use of this form will be illustrated in this section.

2.8.1 Line Terminated with an Arbitrary Load

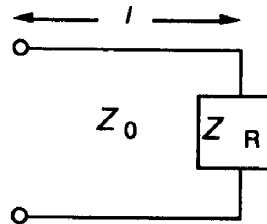


Figure 2.8.1.1: Line Terminated With Z_R

The impedance seen looking into a generalized transmission line terminated in Z_R (as shown in Figure 2.8.1.1) is:

$$Z_{in} = Z_0 \left(\frac{Z_R \cosh \gamma l + Z_0 \sinh \gamma l}{Z_0 \cosh \gamma l + Z_R \sinh \gamma l} \right) \quad (\Omega) \quad (2.8.1.1)$$

This equation comes from the solution for the RLCG line (2.2.25) where we have chosen hyperbolic trigonometric functions rather than e^x when simplifying. For a lossless line, the above reduces to:

$$Z_{in} = Z_0 \left(\frac{Z_R \cos \beta l + j Z_0 \sin \beta l}{Z_0 \cos \beta l + j Z_R \sin \beta l} \right) \quad (2.8.1.2)$$

where

$$\alpha = 0$$

$$\beta = \frac{2\pi}{\lambda}$$

$$\gamma = j\beta = \frac{2\pi j}{\lambda}$$

Note that in both (2.8.1.1) and (2.8.1.2) the impedance seen is a function of the load impedance, Z_R , the line parameters, α and β , and the distance from the load, l . Note that the form of the equations for α and β is in units per wavelength. That is, the loss per wavelength is a constant. Note that βl is the electrical length of a line with length l in radians.

2.8.1.1 Example: Quarter-Wave Line with Open Circuit Load

What is the impedance looking into a lossless line one-quarter of a wavelength long terminated in an open circuit?

Solution: We have been given that

$$l = \lambda / 4$$

$$Z_R = \infty \Omega$$

So from Equation (2.8.1.1):

$$Z_{in} = Z_0 \left(\frac{Z_R \cos \beta l + j Z_0 \sin \beta l}{Z_0 \cos \beta l + j Z_R \sin \beta l} \right) = Z_0 \left(\frac{\infty \cos \frac{\pi}{2} + j Z_0 \sin \frac{\pi}{2}}{Z_0 \cos \frac{\pi}{2} + j \infty \sin \frac{\pi}{2}} \right)$$

$$Z_{in} = \frac{Z_0}{\infty} = 0 \Omega$$

The quarter-wavelength line has transformed the output open circuit into a short circuit at the input of the line.

2.8.1.2 Example: Quarter-Wave Line Terminated with a Generalized Load

What happens when we put a general impedance Z_L at the end of a quarter-wavelength line?

Solution: We have been given that

$$l = \lambda / 4$$

$$Z_R = Z_L$$

So again using (2.8.1.1):

$$Z_{in} = Z_0 \left(\frac{Z_R \cos \beta l + j Z_0 \sin \beta l}{Z_0 \cos \beta l + j Z_R \sin \beta l} \right) = Z_0 \left(\frac{j Z_0}{j Z_L} \right) = \frac{Z_0^2}{Z_L}$$

This is often expressed as the quarter-wave line impedance, Z_0 required to match a desired Z_{in} to a specified Z_L as:

$$Z_0 = \sqrt{Z_{in} Z_L} \quad (2.8.1.2.1)$$

For example, if our source impedance is 50Ω and we wish to match it to a 68Ω load, we could use a quarter-wave line having impedance:

$$Z_0 = \sqrt{50 \times 68} = \sqrt{3400} = 58.31 \Omega$$

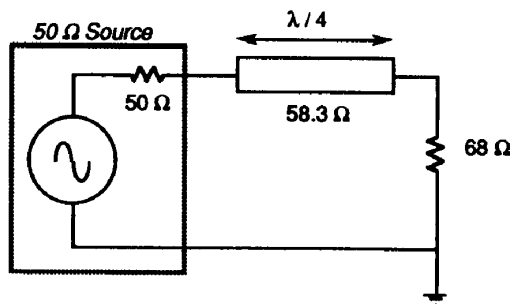


Figure 2.8.1.2.1: Quarter-Wave Transformer

The quarter-wavelength required would be calculated from the frequency of the source and the ϵ_{eff} of the transmission line. As shown in Figure 2.8.1.2.2, the impedance is matched exactly at the design frequency (in this case 1 GHz). The impedance looking from the source into the line is 50Ω , and the impedance looking back from the load is 68Ω at 1 GHz. The match is good to within some tolerance over a wider bandwidth and the response repeats.

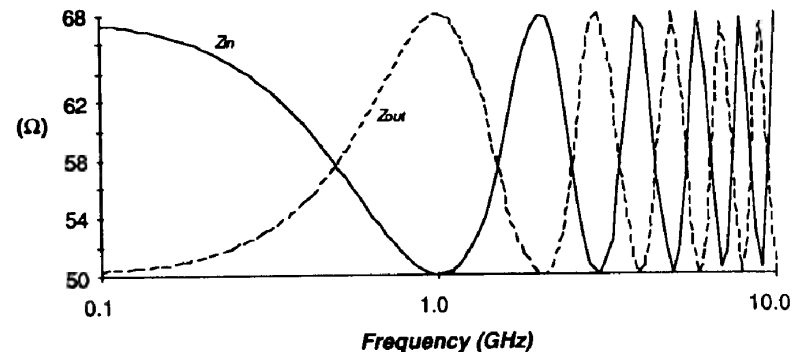


Figure 2.8.1.2.2: Input and Output Impedance of Quarter-Wavelength Transformer

This property is the basis for multiple-step transformers and tapered lines used for matching unequal impedances.

2.8.2 Line Terminated with Z_0

If we let the arbitrary load in Equation (2.8.1.1), Z_R , equal the characteristic impedance of the line, we get:

$$Z_{in} = Z_0 \left(\frac{Z_0 \cosh \gamma l + Z_0 \sinh \gamma l}{Z_0 \cosh \gamma l + Z_0 \sinh \gamma l} \right) = Z_0 \quad (2.8.2.1)$$

which is no longer a function of l , the distance from the load. This condition is known as *matched*.

2.8.3 Line Terminated with Short Circuit

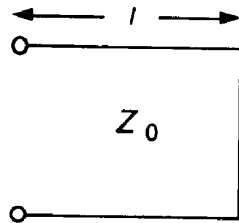


Figure 2.8.3.1: Short-Circuited Line

If we terminate the line in a short circuit ($Z_R = 0$), equation (2.8.1.1) gives:

$$Z_{in} = Z_{sc} = Z_0 \left(\frac{Z_0 \sinh \gamma l}{Z_0 \cosh \gamma l} \right) = Z_0 \tanh \gamma l \quad (2.8.3.1)$$

For lossless lines, Equation (2.8.1.1) reduces to:

$$Z_{sc} = j Z_0 \tan \beta l = j Z_0 \tan \left(\frac{2\pi l}{\lambda} \right) \quad (2.8.3.2)$$

If $\tan x$ is now approximated with x (true for small x):

$$Z_{sc} \approx j Z_0 \beta l = j Z_0 \left(\frac{2\pi l}{\lambda} \right) = j \omega \frac{(Z_0 l)}{c} \quad (2.8.3.3)$$

This equation looks a great deal like the impedance of an inductor ($Z_{ind} = j \omega L$) with inductance $Z_0 l$. In fact, physically short, short-circuited lines are used in this fashion as inductors in frequency ranges where lumped inductors become unusable. We can increase the apparent inductance by either increasing the line length or by increasing its characteristic impedance, Z_0 .²

²In designing transmission line filters the second part of (2.8.3.2) is known as *Richards' variable*, Ω

$$\Omega = \tan \left(\frac{2\pi l}{\lambda} \right) = \tan \left(\frac{\pi f_c}{2.0 f_0} \right) \quad (2.8.3.4)$$

For example, a low pass filter could be designed by using $\lambda/4$ lines (at f_0 in the reject band). The cutoff frequency, f_c is then used with (2.8.3.2), (2.8.3.4), and the filter g_k s to calculate the required line lengths.

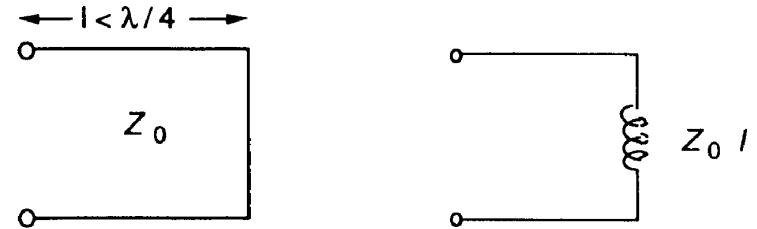


Figure 2.8.3.2: Short-Circuited Line Used as an Inductor

Notice that the tangent function in (2.8.3.2) is periodic in its parameter and Z_{sc} will therefore also repeat. Also, the goodness of our lumped (2.8.3.3) approximation to the tangent function degrades as the line's length relative to a wavelength increases, which effectively makes the line's equivalent inductance a function of frequency.

Similarly, for:

$$\pi/2 < \left(\frac{2\pi l}{\lambda} \right) < \pi$$

the tangent function is negative and Z_{sc} is approximately:

$$Z_{sc} \approx -j Z_0 \beta l = -j Z_0 \left(\frac{2\pi l}{\lambda} \right) = -j \omega \frac{Z_0 l}{c} = -j \frac{1.0}{\left(\frac{c}{\omega Z_0 l} \right)} \quad (2.8.3.5)$$

which is the impedance of a capacitor of value

$$C = \frac{1.0}{\omega^2 Z_0 l} \quad (2.8.3.6)$$

Finally, for

$$2\pi l / \lambda = \pi/2 \quad (2.8.3.7)$$

or equivalently,

$$l = \lambda/4$$

the tangent function of (2.8.3.2) and the line's impedance is infinite. This property is used to create series resonant circuits.

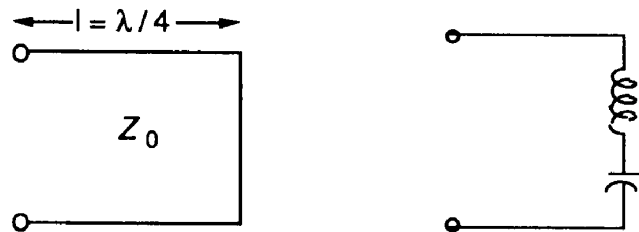


Figure 2.8.3.3: Short-Circuited Line Used as a Resonant Circuit

Summarizing our results for the short-circuited line:

$$\text{"inductive"} \begin{cases} l < \frac{\lambda}{4} \\ \frac{\lambda}{2} < l < \frac{3\lambda}{4} \end{cases}$$

$$\text{"capacitive"} \begin{cases} \frac{\lambda}{4} < l < \frac{\lambda}{2} \\ \frac{3\lambda}{4} < l < \lambda \end{cases}$$

$$\text{"series LC resonator"} \left\{ l = \frac{n\lambda}{4} \text{ where } n \in \mathbb{I} \text{ (the integers)} \right.$$

2.8.3.1 Example: Resonant Line

Design a series resonator from a shorted 50 Ω coaxial cable having a PTFE dielectric at 402 MHz.

Solution: We have been given:

$$Z_0 = 50 \Omega$$

$$\epsilon_r = 2.10$$

$$Z_R = 0 \Omega$$

We calculate the wavelength of 402 MHz in PTFE from Equation (2.3.12):

$$\lambda_g = \frac{c}{f \sqrt{\epsilon_r}} = \frac{3 \times 10^8 \text{ m/s}}{402 \times 10^6 \text{ /s} \sqrt{2.10}} = 0.5149 \text{ m} = 20.27 \text{ in}$$

$$\lambda_g / 4 = 0.5149 \text{ m} / 4 = 0.1287 \text{ m} = 5.07 \text{ in}$$

2.8.3.2 Example: Line Used as a Capacitor

Design a capacitive reactance from a shorted 50 Ω coaxial cable with an air dielectric equivalent to a 1 pF capacitor at 1.5 GHz.

Solution: We have been given:

$$Z_0 = 50 \Omega$$

$$\epsilon_r = 1$$

$$Z_{in} = \frac{1}{j \omega C} = \frac{1}{j 2 \pi \times 1.5 \times 10^9 \text{ /s} \times 1.0 \times 10^{-12} \text{ F}} = -j 106.103 \Omega$$

$$Z_R = 0 \Omega$$

We assume:

$$\frac{\lambda}{4} < l < \frac{\lambda}{2}$$

We calculate the required length with (2.8.3.2):

$$Z_{sc} = j Z_0 \tan \beta l = j Z_0 \tan \left(\frac{2\pi}{\lambda} l \right) = -j 106.103 \Omega$$

and

$$\lambda_g = \frac{c}{f \sqrt{\epsilon_r}} = \frac{3 \times 10^8 \text{ m/s}}{1.5 \times 10^9 \text{ /s}} = 0.2 \text{ m} = 7.874 \text{ in}$$

$$\lambda_g / 4 = 0.2 \text{ m} / 4 = 0.05 \text{ m} = 1.969 \text{ in (our answer will be longer than this)}$$

Solving,

$$\tan^{-1} \left(\frac{-50 \Omega}{106.103 \Omega} \right) = -0.440 \text{ radians} = -25.216^\circ$$

The calculator's inverse tangent function gives us a value in the fourth quadrant but we know we want an answer in the second quadrant so we add the above to π :

$$\pi + (-0.440) = 2.701 \text{ radians} = 154.768^\circ$$

$$l = \frac{2.701 \text{ radians } \lambda_g}{2 \pi}$$

$$l = 0.086 \text{ m} = 3.385 \text{ in}$$

2.8.4 Line Terminated with Open Circuit

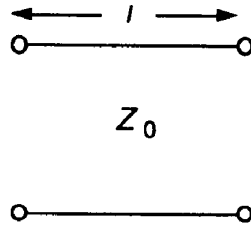


Figure 2.8.4.1: Open-Circuited Line

As we did for the short-circuited line in the last section, we simplify (2.8.1.1). The equation for the input impedance of an open circuited line is:

$$Z_{oc} = Z_0 \coth \gamma l \quad (2.8.4.1)$$

for a lossless line:

$$Z_{oc} = -j Z_0 \cot \beta l = -j Z_0 \cot \left(\frac{2 \pi l}{\lambda} \right) \quad (2.8.4.2)$$

and we can similarly derive that for an open circuited line:

$$\text{"capacitive"} \begin{cases} l < \frac{\lambda}{4} \\ \frac{\lambda}{2} < l < \frac{3\lambda}{4} \end{cases}$$

$$\text{"inductive"} \begin{cases} \frac{\lambda}{4} < l < \frac{\lambda}{2} \\ \frac{3\lambda}{4} < l < \lambda \end{cases}$$

$$\text{"parallel LC resonator"} \left\{ l = \frac{n \lambda}{4} \text{ where } n \in \mathbb{I} \text{ (the integers)} \right.$$

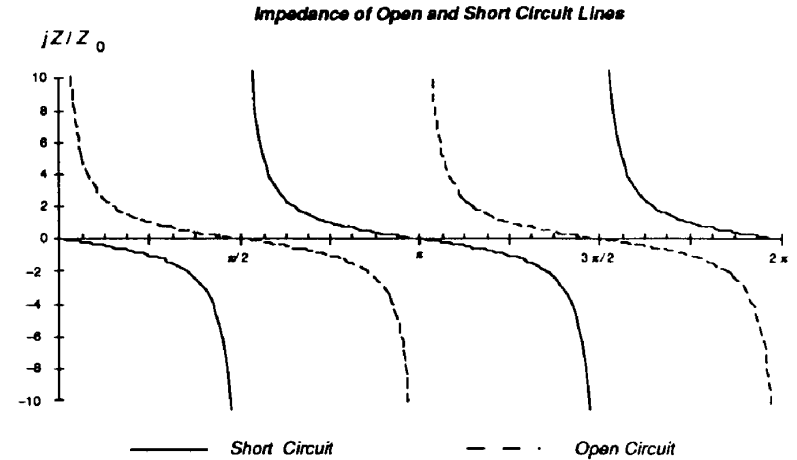


Figure 2.8.4.2: Impedance of Open- and Short-Circuited Lines

2.8.4.1 Example: Chebyshev LPF

From the specification it is determined that an $n = 5$, 0.5 dB Chebyshev LPF is required. The required corner frequency is 4.0 GHz and the reject band is 8.0 to 10.0 GHz. Design the lumped filters. Using the relations of the previous two sections, convert this to a distributed equivalent.

Solution: Using the usual filter tables or equations we find the g_k s required

$$g_1 = g_5 = 1.81$$

$$g_2 = g_4 = 1.3$$

$$g_3 = 2.69$$

The lumped element values are then

$$C_2 = C_4 = \frac{g_k}{2.0 \pi f_c Z_0} = 3.60 \text{ nH}$$

$$L_1 = L_5 = \frac{gk Z_0}{2.0 \pi f_c} = 1.03 \text{ pF}$$

$$L_3 = 5.35 \text{ nH}$$

The distributed element filter is calculated by using lines which are $\lambda/4$ at the center of the reject band or 9.0 GHz. The inductors of the lumped element design will be replaced with short-circuited lines and the capacitors with open-circuited lines. For physical reasons, it is most convenient to use series short-circuited lines and shunt open-circuited lines.

Richards' variable at the corner frequency is

$$\Omega_c = \tan\left(\frac{\pi 4.0 \text{ GHz}}{2.0 \times 9.0 \text{ GHz}}\right) = 0.8391$$

We calculate the line impedances to be

$$Z_1 = Z_5 = \frac{gk Z_0}{\Omega_c} = 107.9 \Omega \quad \text{series short-circuit line}$$

$$Z_2 = Z_4 = \frac{gk}{Z_0 \Omega_c} = 32.3 \Omega \quad \text{shunt open-circuit line}$$

$$Z_3 = 160.3 \Omega \quad \text{series short-circuit line}$$

The line dimensions (lengths and widths) are calculated with the equations of Chapter 3 for the desired method of construction. In an actual design we would then add step junction discontinuities to our model, and correct the open-circuited line lengths for fringing effects.

The response of the two filters constructed with ideal elements is shown in Figure 2.8.4.3 below. The limited number of simulation points creates the finite depth of the nulls and the roughness in the distributed element plot. As expected, the distributed filter response repeats limiting its useful bandwidth (a lumped element filter constructed with real-world components would also show returns due to strays, etc.). This design is an approximation to the lumped design because there is no broadband equality between lumped and distributed elements.

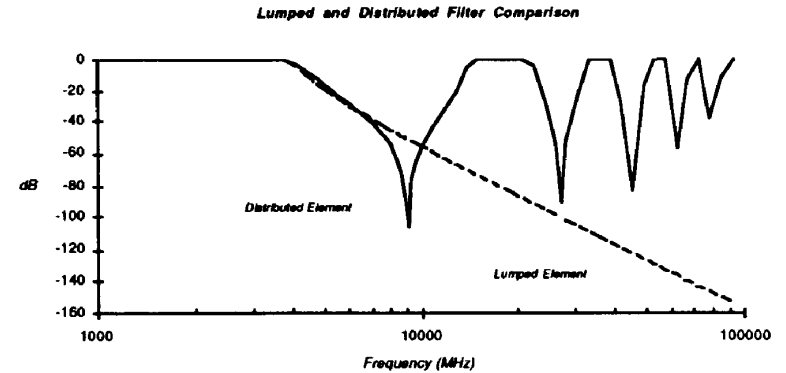


Figure 2.8.4.3: Comparison of Lumped and Distributed Filters

2.8.5 Line Terminated with a Lumped Element

In addition to open circuits and short circuits we can place lumped elements on the end of our transmission lines. This effectively changes the length of the line as we shall see in the next example.

2.8.5.1 Example: Varactor Tuned Line Length

Calculate the value of a tunable capacitor (varactor) required to vary an open-circuited 50 Ω PTFE dielectric coaxial line's apparent length from 0.1 to 0.2 λ at 2.5 GHz.

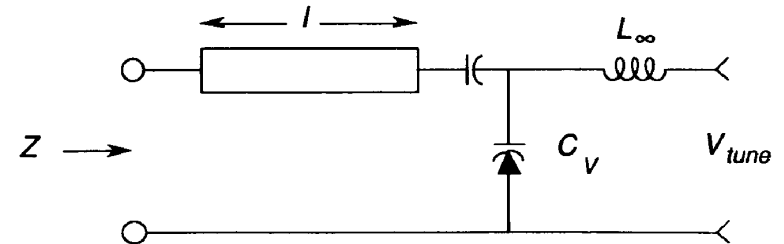


Figure 2.8.5.1.1: Varactor Tuned Line

Solution: We have been given

$$\epsilon_r = 2.10$$

$$f = 2.5 \times 10^9 \text{ /s}$$

$$Z_0 = 50 \Omega$$

$$Z_R = \frac{1}{j \omega C}$$

$$\lambda_g = \frac{c}{f \sqrt{\epsilon_r}} = \frac{3 \times 10^8 \text{ m/s}}{2.5 \times 10^9 \text{ /s} \times \sqrt{2.10}} = 0.083 \text{ m} = 3.260 \text{ in}$$

We know that we wish to tune the capacitance to increase the length. As the length of an open-circuited line ($l < \lambda/4$) is increased, the impedance decreases and the effective capacitance therefore is increased. So, we know that as the capacitor is tuned from its minimum to its maximum capacitance, the line's apparent length will increase. We now calculate the effective capacitance of an open-circuited line at the two requested line lengths:

$$Z_{oc}(0.1 \lambda) = -j Z_0 \cot\left(\frac{2 \pi l}{\lambda}\right) = -j 68.819 \Omega$$

$$Z_{oc}(0.2 \lambda) = -j Z_0 \cot\left(\frac{2 \pi l}{\lambda}\right) = -j 16.246 \Omega$$

We will assume that we can obtain a capacitor that can be set to 0 pF for this example, although in real life we would need to shorten the line to compensate for the adjustable capacitor's minimum capacitance. In our case, the line's physical length will be 0.1λ .

We now use the general form equation for the impedance of an arbitrarily terminated line (2.8.1.1):

$$l = 0.1 \lambda$$

$$Z_{in} = Z_0 \left(\frac{Z_R \cos 0.2 \pi + j Z_0 \sin 0.2 \pi}{Z_0 \cos 0.2 \pi + j Z_R \sin 0.2 \pi} \right) = -j 16.246 \Omega$$

Solving,

$$-j 0.325 = \left(\frac{Z_R \cos 0.2 \pi + j Z_0 \sin 0.2 \pi}{Z_0 \cos 0.2 \pi + j Z_R \sin 0.2 \pi} \right)$$

$$-j 0.325 = \left(\frac{Z_R 0.8090 + j 29.3893}{40.4508 + j Z_R 0.5878} \right)$$

$$(-j 0.325)(40.4508 + j Z_R 0.5878) = Z_R 0.8090 + j 29.3893$$

$$-13.1465 j + 0.1910 Z_R = Z_R 0.8090 + j 29.3893$$

$$Z_R = \frac{42.5358 j}{-0.6180} = -j 68.8282 \Omega = \frac{-j}{\omega C}$$

$$C = 0.9249 \text{ pF}$$

A similar analysis could have been used to tune the line's $\lambda/4$ resonant frequency, giving us a tunable resonator.

2.9 ABCD MATRIX OF A GENERALIZED TRANSMISSION LINE

Use of the ABCD matrix is described in Chapter 12. The ABCD matrix for a series transmission line is:

$$\begin{bmatrix} \cosh \gamma l & Z_0 \sinh \gamma l \\ \frac{\sinh \gamma l}{Z_0} & \cosh \gamma l \end{bmatrix} \quad \text{general case}$$

$$\begin{bmatrix} \cos \beta l & j Z_0 \sin \beta l \\ \frac{j \sin \beta l}{Z_0} & \cos \beta l \end{bmatrix} \quad \text{lossless case}$$

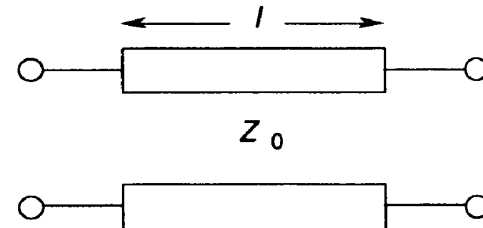


Figure 2.9.1: Series Transmission Line

The ABCD matrix for a shunt open-circuited transmission line is:

$$\begin{bmatrix} 1.0 & 0.0 \\ \frac{j \tan \beta l}{Z_0} & 1.0 \end{bmatrix}$$

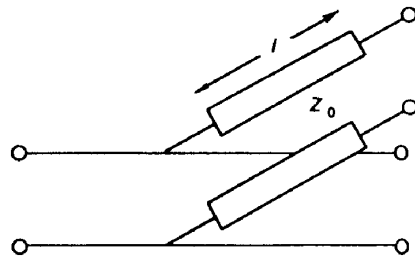


Figure 2.9.2: Shunt Open-Circuited Transmission Line

The ABCD matrix for a shunt short-circuited transmission line is:

$$\begin{bmatrix} 1.0 & 0.0 \\ \frac{-j}{Z_0 \tan \beta l} & 1.0 \end{bmatrix}$$

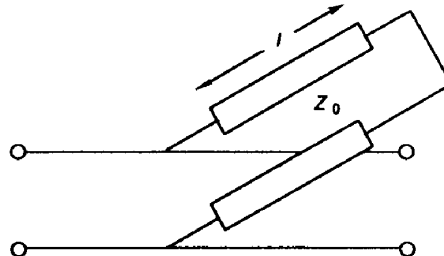


Figure 2.9.3: Shunt Short-Circuited Transmission Line

2.10 SUMMARY

In this chapter, we have seen, in a general way, how transmission line calculations are made and how transmission lines can be used as circuit elements. We have seen that they can be used to emulate lumped elements and resonators. We have learned that the impedance of a transmission line element does not vary with frequency the way that a lumped element does. In fact, it is clear that a good definition of a *lumped* element is a section of transmission line that is electrically short. As we saw, when the assumption of short electrical length was made in the above examples, the characteristics of lines reduce to that of lumped elements.

In Chapter 3 we will see how to translate these general line equations into actual physical realizations.

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Chapter 3

Physical Transmission Lines

3.1 INTRODUCTION

In this chapter we transform our knowledge about the generalized transmission line calculations into actual physical realizations for a wide variety of configurations. Structures are grouped into major categories: coaxial, paired lines, coplanar, microstrip line, stripline, waveguide, and others. Within each category, a number of nonstandard configurations are included to illustrate the effects of variations from the standard structure. Tables and graphs are generally excluded because it is assumed that the reader has access to and will use the *Transmission Line Design Software* available as part of this text.

ABCD matrices for the lines in this chapter are found in Chapter 2. The equations of this chapter allow calculation of Z_0 in the matrices from the line's physical dimensions. The electrical length is calculated from the operating frequency, f , and the physical length, l , using ϵ_{eff} .

$$\theta = 2.0 \pi l f \frac{\sqrt{\epsilon_{eff}}}{c} \text{ (radians)} = 360.0 l f \frac{\sqrt{\epsilon_{eff}}}{c} \text{ (degrees)} \quad (3.1.1)$$

A number of the older and commonly referenced articles are available as part of reprint collections listed below. Two of these, [2] and [4] are out of print but they may be found in libraries. Consulting these can save considerable time.

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3.2 COAXIAL STRUCTURES

3.2.1 Round Coaxial Cable

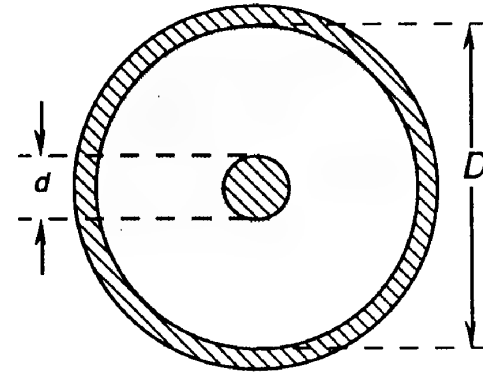


Figure 3.2.1.1: Round Coaxial Cable

The equations for this transmission line can be derived easily and are exact for the case of an infinitely long line (no fringing fields):

$$Z_0 = \frac{\eta_0}{2.0 \pi \sqrt{\epsilon_r}} \ln \left(\frac{D}{d} \right) \quad (\Omega) \quad (3.2.1.1)$$

$$\alpha_c = \left(\frac{0.014272 \sqrt{f}}{Z_0} \right) \left(\frac{1}{d} + \frac{1}{D} \right) \quad (\text{dB} / \text{m}) \quad (3.2.1.2)$$

$$\alpha_d = 0.091207 f \sqrt{\epsilon_r} \tan \delta \quad (\text{dB} / \text{m}) \quad (3.2.1.3)$$

For inner and outer conductors having equal conductivity, the optimal D/d for breakdown voltage is 2.718; for optimal power transfer, 1.65; and for minimum attenuation, 3.59 [7, 13]. See [7] for information on optimizing these parameters.

Coaxial connectors have mechanical tolerances of ± 0.0025 mm and ± 0.005 mm for laboratory and general precision, respectively. For a 50 Ω connector, this corresponds to about 0.5% and 0.8% tolerance on the impedance.

To reduce losses, we might increase d and D , keeping their ratio constant to lower α_c . However, the coaxial wire of Figure 3.2.1.1 will propagate higher order modes.

The cutoff frequency of these modes decreases with increasing d , and D placing a limitation on the cable size for a desired useful frequency range. Marcuvitz [11] gives the relations for the higher order modes, which begin at approximately 1.5 times the TE₁₁ mode. The first higher order mode begins to propagate when

$$\lambda_{C11} = \frac{2.0 \pi}{\beta_{z1}} \quad (3.2.1.4)$$

where β_{zm} is the first ($m = 1$) solution of

$$\frac{J_1'(\beta_{zm} R)}{N_1'(\beta_{zm} R)} = \frac{J_1'(\beta_{zm} r)}{N_1'(\beta_{zm} r)} \quad (3.2.1.5)$$

where

$$r = d / 2.0 \quad (3.2.1.6)$$

$$R = D / 2.0 \quad (3.2.1.7)$$

Dimitrios [5] rewrites the above by replacing derivative forms of Bessel functions with nonderivative forms. This form is more accurate when used with available tables of Bessel functions:

$$\frac{J_0'(\beta_{zm} R) - \frac{1}{\beta_{zm} R} J_1'(\beta_{zm} R)}{N_0'(\beta_{zm} R) - \frac{1}{\beta_{zm} R} N_1'(\beta_{zm} R)} = \frac{J_0'(\beta_{zm} r) - \frac{1}{\beta_{zm} r} J_1'(\beta_{zm} r)}{N_0'(\beta_{zm} r) - \frac{1}{\beta_{zm} r} N_1'(\beta_{zm} r)} \quad (3.2.1.8)$$

Green [8] derives an expression for the wavelength at which this occurs (TE₁₁ mode begins to propagate):

$$\lambda_c = 2 \pi r_m \left[1.0 - \frac{1.0}{6.0} \left(\frac{t}{2.0 r_m} \right)^2 - \frac{7.0}{120.0} \left(\frac{t}{2.0 r_m} \right)^4 - \dots \right] \quad (3.2.1.9)$$

where

$$r_m = \frac{d + D}{2.0} \quad (3.2.1.10)$$

$$t = D - d \quad (3.2.1.11)$$

The stated agreement with Marcuvitz's Bessel function solution is better than 1% for $D/d < 2.5$, and better than 4% for $D/d < 4.0$. The number of terms used to reach this agreement is not stated.

A simple approximation to the above is:

$$\lambda_c \cong \frac{\pi (D + d)}{2.0} \quad (\text{units of } D, d) \quad (3.2.1.12)$$

which is the circumference of a circle with radius midway between the inner and outer conductors. Dimitrios [5] compared the accuracy of the equation above to the exact solution and found that it is generally accurate to 3% for 50 Ω lines with various dielectrics.

A lesser known fact about real-world coaxial cables is that they can introduce nonlinear distortion at a low level. Amin [3] presents data at the L, S, and C (390 MHz to 10.9 GHz) bands for various commercial cables showing that braided cables can have distortion products down 90 to 115 dB from the carrier depending on construction. Solid shields had no distortion. Nickel-plated, stainless steel, and Al alloy braids had distortion down 90 to 95 dB from the carrier. Distortion increased with power, frequency, and cable length.

3.2.1.1 Example: Coaxial Pogo Pin

We wish to design a 50 Ω controlled impedance path that launches from a microstrip line and makes use of a standard size (0.040" o.d.) pogo pin (a spring-loaded pin used in ATE equipment) as the center conductor. This configuration is of use in test equipment where a controlled Z path must be blind-mated repetitively, as in Figure 3.2.1.2. Calculate the required shield dimensions for both air and PTFE dielectrics.

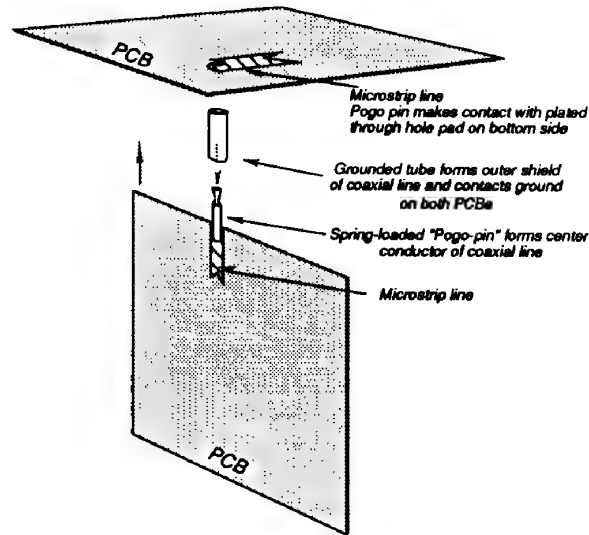


Figure 3.2.1.2: Coaxial Pogo-Pin Example

Solution:

$$Z_0 = 50 \Omega = \frac{\eta_0}{2.0 \pi \sqrt{\epsilon_r}} \ln \left(\frac{D}{d} \right) \Omega$$

For air:

$$50 \Omega = \frac{377}{2 \pi \sqrt{1}} \ln \left(\frac{D}{0.040} \right) \Omega$$

$$D = 0.092"$$

For PTFE:

$$50 \Omega = \frac{377}{2 \pi \sqrt{2.1}} \ln \left(\frac{D}{0.040} \right) \Omega$$

$$D = 0.134"$$

A number of companies have attempted to combine the 360° shielding benefits of the coaxial structure with the fabrication advantages of printed-circuit boards (PCBs). This is increasingly an issue as the density and speeds of digital circuitry increase. See also [15] for a planar structure constructed with additive techniques. See Section 3.2.5 for more information.

Swengel *et al.* [18] describe a technique for machine wiring coaxial lines on PCB. The technique point-to-point wires 3.14 mil wires coated with a 10 mil thick layer of PTFE. The wire ends are then plated together. The PTFE is plated, creating the coaxial shield. A final selective etch separates the wire ends as required. Because of the solid ground plane, wave-soldering is difficult. The technique is also sensitive to handling and vibration.

A similar technique, described in [21], uses a semirigid cable (PTFE, $D = 9.5$ mil, $d = 3.14$ mil), which is also machine routed point-to-point. The wires are laid down on an adhesive layer during routing and then the shields are plated together. The wire is stripped at its ends, coated with epoxy and the center conductors are plated to plated-through holes. This structure is easily used as part of a multilayer PCB.

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3.2.2 Partially Filled Round Coaxial Cable

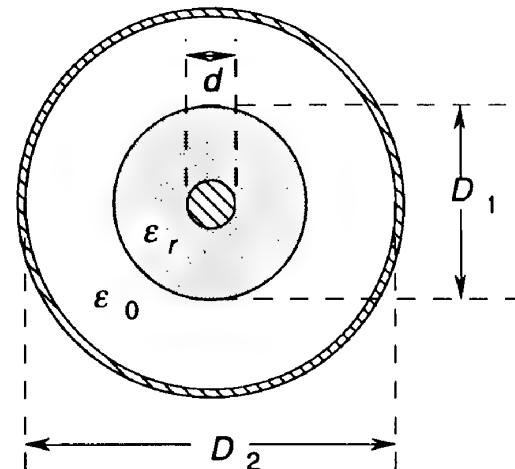


Figure 3.2.2.1: Partially Filled Round Coax

This case is encountered where a low-loss line is being constructed and the dielectric losses need to be reduced by making part of the dielectric air. It also occurs when the dielectric diameter is reduced over a portion of its length for assembly.

$$Z_0 = \frac{\eta_0}{2\pi} \ln\left(\frac{D_2}{d}\right) \sqrt{\frac{\epsilon_r \ln\left(\frac{D_2}{d}\right) + \ln\left(\frac{D_1}{d}\right)}{\epsilon_r \ln\left(\frac{D_2}{d}\right)}} \quad (\Omega) \quad (3.2.2)$$

Ragan [3] also discusses the matching of such structures by proper choice of dimensions. Hatsuda [1] analyzes a round coaxial structure with wedges of dielectric removed.

3.2.2.1 Example: Calculating q for a Partially Filled Coaxial Line

Calculate q for the partially filled round coaxial structure of 3.2.2.1.

Solution: The ratio of the structure's actual capacitance to its capacitance with solid air as the dielectric is defined as q . We can easily get C_{air} from the solid dielectric structure:

$$C_{air} = \frac{2 \pi \epsilon_0}{\ln \left(\frac{D_2}{d} \right)}$$

and using the relationships in Chapter 1 we can write the capacitance in terms of the known Z_0 :

$$C_{actual} = \frac{\sqrt{\mu_0 \epsilon}}{Z_0}$$

$$q = \frac{C_{actual}}{C_{air}} = \frac{\frac{\sqrt{\mu_0 \epsilon}}{Z_0}}{\frac{2 \pi \epsilon_0}{\ln \left(\frac{D_2}{d} \right)}} = \frac{\frac{\eta_0}{2 \pi} \ln \left(\frac{D_2}{d} \right) \sqrt{\frac{\epsilon_r \ln \left(\frac{D_2}{d} \right) + \ln \left(\frac{D_1}{d} \right)}{\epsilon_r \ln \left(\frac{D_2}{d} \right)}}}{\frac{2 \pi \epsilon_0}{\ln \left(\frac{D_2}{d} \right)}}$$

$$q = \frac{\epsilon_r}{\sqrt{\frac{\epsilon_r \ln \left(\frac{D_2}{d} \right) + \ln \left(\frac{D_1}{d} \right)}{\epsilon_r \ln \left(\frac{D_2}{d} \right)}}}$$

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3.2.3 Eccentric Round Coaxial Cable

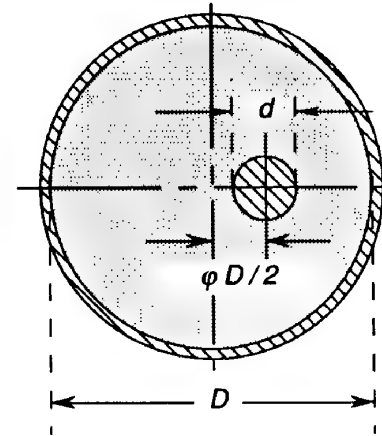


Figure 3.2.3.1: Eccentric Round Coaxial Cable

The eccentric coax equations enable us to analyze the effects of tolerances in the manufacture of the cable. This structure has also been used as a continuously adjustable $\lambda/4$ line. The center conductor is moved within the outer shield by a mechanical probe resulting in smooth variation of the characteristic impedance.

For center conductors off center in one direction, [1] gives:

$$Z_0 = \frac{\eta_0}{2.0 \pi \sqrt{\epsilon_r}} \cosh^{-1} \left[\frac{D}{2.0 d} (1.0 - \phi^2) + \frac{d}{2.0 D} \right] (\Omega) \quad (3.2.3.1)$$

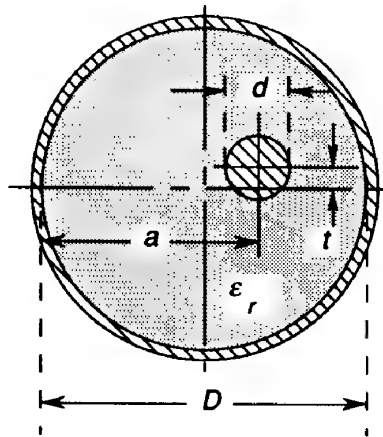


Figure 3.2.3.2: Eccentric Round Coaxial Cable

For center conductors that are eccentric in both directions, as in Figure 3.2.3.2 [2]:

$$Z_0 = \left[\frac{\eta_0}{2\pi\sqrt{\epsilon_r}} \ln \left(\frac{D}{d} \right) \right] \left[1.0 - \frac{e^2 k^2}{(k^2 - 1.0) \log k} \right] (\Omega) \quad (3.2.3.2)$$

where

$$b = \frac{d}{2.0} \quad (3.2.3.3)$$

$$c = t \quad (3.2.3.4)$$

$$e = \frac{c}{a} \quad (3.2.3.5)$$

$$k = \frac{a}{b} \quad (3.2.3.6)$$

The conductor losses of this line are

$$\alpha_c = \alpha_{c,centered} \left[1.0 + \frac{2.0 e^2}{k} - \frac{e^2 k^2}{(k^2 - 1.0) \log k} \right] \text{ (dB / m)} \quad (3.2.3.7)$$

$\alpha_{c,centered}$ = conductor losses of the centered line calculated with (3.2.1.2).

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3.2.4 Square Coaxial Cable

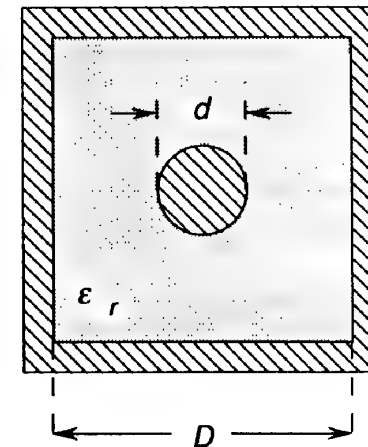


Figure 3.2.4.1: Square Coaxial With Circular Center Conductor

The square coaxial wire configurations are useful at very high frequencies as smaller replacements for waveguide. The square cross section makes them easier to fabricate than a round coax.

For the round center conductor as shown in Figure 3.2.4.1:

$$Z_0 = \frac{\eta_0}{2\pi\sqrt{\epsilon_r}} \ln \left[\frac{1.0787 D}{d} \right] (\Omega) \quad (3.2.4.1)$$

and for $Z_0 \leq 2.0 \Omega$

$$Z_0 = 21.2 \sqrt{D/d - 1.0} (\Omega) \quad (3.2.4.2)$$

which was derived by S. Frankel. Error of the first equation is less than 1.5% above 17 Ω and very small above 30 Ω [2]. The second equation is accurate to 0.5 Ω .

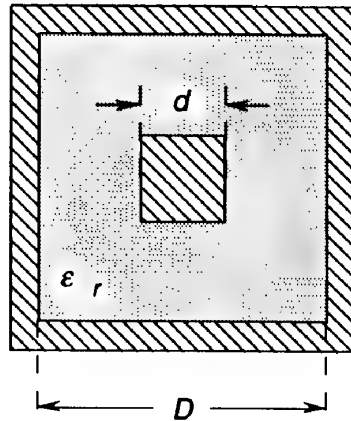


Figure 3.2.4.2: Square Coax with Square Center Conductor

For the square center conductor configuration shown in Figure 3.2.4.2 a conformal mapping solution is available [6]:

$$Z_0 = 5.0 \pi \nu_p \times 10^{-8} \frac{K(k')}{K(k)} = \frac{5.0 \pi c}{\sqrt{\epsilon_r}} \frac{K(k')}{K(k)} \times 10^{-8} \quad (\Omega) \quad (3.2.4.3)$$

$$k = \frac{(\Lambda' - \Lambda)^2}{(\Lambda' + \Lambda)^2} \quad (3.2.4.4)$$

$$\Lambda' = \sqrt{1.0 - \Lambda^2} \quad (3.2.4.5)$$

where Λ is the solution to:

$$\frac{K(\Lambda')}{K(\Lambda)} = \frac{D - d}{D + d} \quad (3.2.4.6)$$

This can be solved iteratively with an equation solver or using the relations of Chapter 12. Observe that the above is always ≥ 0 and ≤ 1 .

An approximation [6] is:

$$Z_0 \cong \frac{1.0}{4.0 \nu_p (C_{pp} + C_c)} = \frac{\eta_0}{\sqrt{\epsilon_r}} \left[\frac{1.0}{4.0 \left(\frac{2.0 d}{D - d} + 0.558 \right)} \right] (\Omega) \quad (3.2.4.7)$$

which can be seen to be the equation for four parallel plate capacitors (C_{pp}) added to the four corner capacitances (C_c). The accuracy of the above equation is better than 1% for

$$D/d \leq 4.0$$

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3.2.5 Rectangular Coaxial Line

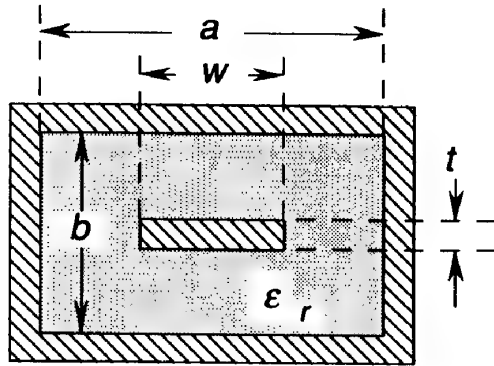


Figure 3.2.5.1: Rectangular Coaxial Line

This structure is a transitional structure between round coax and stripline, microstrip line or other planar lines. Many attempts have been made to utilize this structure for photolithographic construction of fully shielded controlled impedance lines. Rotating the center conductor relative to the shield also allows this structure to have a continuously variable characteristic impedance [4].

As reported in [9] rectangular coax can be used in PCBs constructed with additive techniques to obtain a completely shielded structure. The process by Augat Microtec builds the structure using photolithographic techniques from polyimide and copper.

A related structure [14], Figure 3.2.5.2, was constructed with thick-film techniques. The main limitation appears to have been the thickness, b , achievable without excessive layered printings. Without a sufficiently thick dielectric, center conductor to shield shorts were common. Low dielectric constant pastes were recommended.

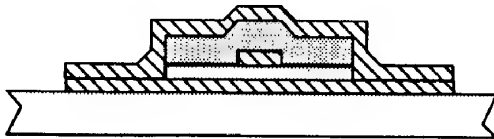


Figure 3.2.5.2: Thick-Film Coaxial Line

$$Z_0 = \frac{\eta_0}{4.0 \sqrt{\epsilon_r}} \left[\frac{1.0}{\frac{w/b}{1.0 - t/b} + \frac{2.0}{\pi} \ln \left(\frac{1.0}{1.0 - t/b} + \coth \frac{\pi a}{2.0 b} \right)} \right] (\Omega) \quad (3.2.5.1)$$

for a square shield and center conductor, $a = b$ and $w = t$:

$$Z_0 = \frac{\eta_0}{4.0 \sqrt{\epsilon_r}} \left[\frac{1.0}{\frac{1.0}{b/t - 1.0} + \frac{2.0}{\pi} \ln \left(\frac{1.0}{1.0 - t/b} + \coth \frac{\pi}{2.0} \right)} \right] (\Omega) \quad (3.2.5.2)$$

As stated earlier, rectangular coaxial structures are also useful at extremely high frequencies as a smaller replacement for waveguide that is more easily fabricated than round coax.

Rotation of the center conductor varies the characteristic impedance of the line. The equations are in Cruz and Brooke, [4]:

$$Z(\theta) = \frac{Z(0^\circ) (1.0 + \cos 2.0\theta) + Z(90^\circ) (1.0 - \cos 2.0\theta)}{2.0} (\Omega) \quad (3.2.5.3)$$

where $Z(0^\circ)$ and $Z(90^\circ)$ are calculated with any of the equations for a rectangular line above. Although the calculations of $Z(0^\circ)$ and $Z(90^\circ)$ were out of range for the [2] and [1] equations, agreement with experimental data to better than 1.4% was achieved.

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3.2.6 Trough Line or Channel Line

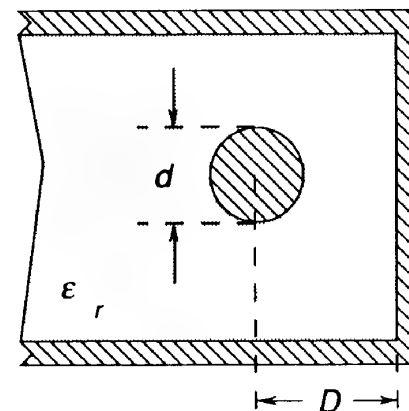


Figure 3.2.6: Trough Line

Wheeler [2] gives:

$$Z_0 = \frac{\eta_0}{4\pi\sqrt{\epsilon_r}} \ln \left[1.0 + \left\{ \frac{1}{2} \left(\frac{4}{\pi} \tanh \frac{\pi}{2} \right)^2 \left[(D/d)^2 - 1.0 \right] \right\} \right. \\ \left. + \sqrt{\left\{ \frac{1}{2} \left(\frac{4}{\pi} \tanh \frac{\pi}{2} \right)^2 \left[(D/d)^2 - 1.0 \right] \right\}^2 + \frac{4}{9} \left[(D/d)^2 - 1.0 \right]} \right] \quad (\Omega) \quad (3.2.6)$$

with a stated accuracy of about 1%.

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3.2.7 Strip-Centered Coaxial Line

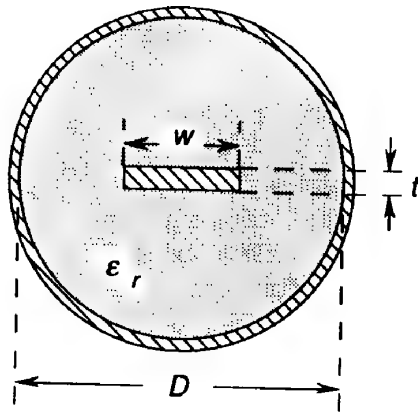


Figure 3.2.7.1: Strip-Centered Coaxial Cable

The configuration above is used to obtain lower loss by building a coaxial cable with a larger center conductor surface area. Bongianni [1] reports cables made with extremely fine dimensions—the center conductor was $150\text{ }\mu\text{m} \times 12.5\text{ }\mu\text{m}$. For $Z_0 \leq 30.0\text{ }\pi / \sqrt{\epsilon_r}$ (Ω):

$$Z_0 = \frac{15.0\text{ }\pi^2}{\sqrt{\epsilon_r}} \frac{1.0}{\ln \left[\frac{2.0(D + w)}{D - w} \right]} \quad (\Omega) \quad (3.2.7.1)$$

For $Z_0 \geq 30.0\text{ }\pi / \sqrt{\epsilon_r}$ (Ω):

$$Z_0 = \frac{60.0}{\sqrt{\epsilon_r}} \ln \left(\frac{2.0 D}{w} \right) \quad (\Omega) \quad (3.2.7.2)$$

There is a bit of a chicken-and-the-egg problem here because it is necessary to know Z_0 in order to choose the equation for Z_0 ; however, the two equations pass quite close and it is okay to use either to make the choice.

REFERENCES

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3.3 PAIRED LINES

3.3.1 Parallel Wires

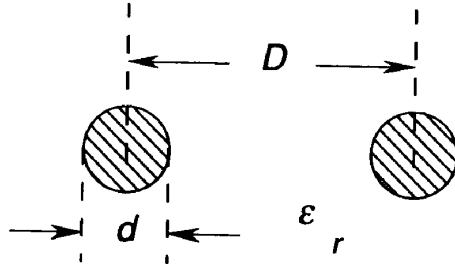


Figure 3.3.1.1: Parallel Wires

$$Z_0 = \frac{\eta_0}{\pi \sqrt{\epsilon_r}} \cosh^{-1} \left(\frac{D}{d} \right) (\Omega) \quad (3.3.1.1)$$

or

$$Z_0 = \frac{\eta_0}{2.0 \pi \sqrt{\epsilon_r}} \cosh^{-1} \left(\frac{2.0 D^2 - d^2}{d^2} \right) (\Omega) \quad (3.3.1.2)$$

valid for

$$d/D \ll 1.0$$

The first reference is found in [4] and [5]. The second was derived with conformal mapping techniques in [2] with a stated accuracy of better than 0.24%. Green, *et al.* [1] gives the conductor losses as:

$$\alpha_c = \frac{P}{2.0 d} \frac{\sqrt{\frac{f \epsilon}{\sigma}}}{\cosh^{-1} v} \quad (\text{nepers/cm}) \quad (3.3.1.3)$$

where

$$P = \frac{v}{\sqrt{v^2 - 1.0}} \quad (3.3.1.4)$$

$$v = \frac{D}{d} \quad (3.3.1.5)$$

and P is a proximity factor; the equation is valid for high frequencies.

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3.3.2 Unequal Size Parallel Wires

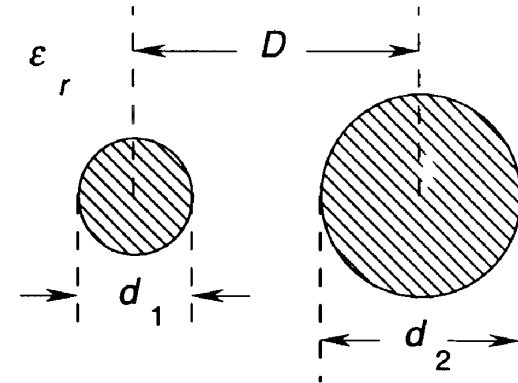


Figure 3.3.2.1: Unequal Size Parallel Wires

$$Z_0 = \frac{\eta_0}{2.0 \pi \sqrt{\epsilon_r}} \cosh^{-1} \left(\frac{4.0 D^2 - d_1^2 - d_2^2}{2.0 d_1 d_2} \right) (\Omega) \quad (3.3.2.1)$$

The stated accuracy [1] is better than 0.24%.

REFERENCES

- [1] Hilberg, Wolfgang, *Electrical Characteristics of Transmission Lines*, Artech House, Norwood, MA, 1979.

3.3.3 Twisted Pair

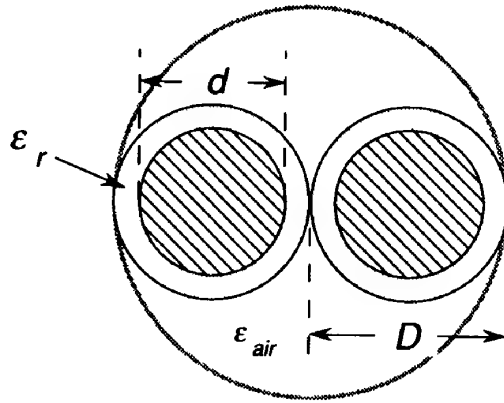


Figure 3.3.3.1: Twisted Pair

The twisted pair provides good low frequency shielding. Undesired signals tend to be coupled equally into each line of the pair. A differential receiver will therefore completely cancel the interference. Lefferson [3] gives design equations for this configuration:

$$Z_0 = \frac{\eta_0}{\pi \sqrt{\epsilon_{eff}}} \cosh^{-1} \left(\frac{D}{d} \right) \quad (\Omega) \quad (3.3.3.1)$$

$$\epsilon_{eff} = 1.0 + q (\epsilon_r - 1.0) \quad (3.3.3.2)$$

$$q = 0.25 + 0.0004 \theta^2 \quad (3.3.3.3)$$

$$T = \frac{\tan \theta}{\pi D} = \text{twists per length} \quad (3.3.3.4)$$

$$\theta = \tan^{-1} (T \pi D) \quad (3.3.3.5)$$

where T and D have the same length units and θ is the pitch angle of the twist; the angle between the twisted pair's center line and the twist. It was found to be optimal for θ to be between 20 and 45°. Smaller angles make the twist loose and create problems in maintaining tolerances. Angles above approximately 50.5° break the wire. The value of q was determined by fitting a line to measurements of the effective dielectric constant.

For the softer insulation PTFE, a different equation should be used for q

$$q = 0.25 + 0.001 \theta^2 \quad (3.3.3.6)$$

An equation for the wire's total length before twisting in terms of number of turns, N , is:

$$l = N \pi D \sqrt{1.0 + \frac{1.0}{\tan^2 \theta}} \quad (3.3.3.7)$$

Lefferson found these equations sufficiently accurate to design transmission lines with $VSWR \leq 1.1:1$.

In [2] a new analysis and experimental data are presented. The derived equations are

$$Z_0 = \sqrt{\frac{L}{C}} \quad (\Omega) \quad (3.3.3.8)$$

$$\epsilon_{eff} = \frac{C}{C_{air}}$$

where

$$L = \left(\frac{\mu_0}{\pi} \right) \cosh^{-1} \left(\frac{D}{d} \right) \quad (3.3.3.9)$$

$$C = C_1 + C_2 - C_3 \quad (3.3.3.10)$$

$$C_1 = \int_a^b \frac{\epsilon_0 dx}{D + (1.0 / \epsilon_r - 1.0) \sqrt{D^2 - x^2} - \frac{\sqrt{d^2 - x^2}}{\epsilon_r}} \quad (3.3.3.11)$$

$$C_2 = \frac{\pi \epsilon_0}{\cosh^{-1} \left(\frac{D}{d} \right)} \quad (3.3.3.12)$$

$$C_3 = \int_a^b \frac{\epsilon_0 dx}{D - \sqrt{d^2 - x^2}} \quad (3.3.3.13)$$

Calculate C_{air} by replacing ϵ_r with ϵ_0 in the equations for C . The integrals are evaluated by the Romberg numerical technique. Accuracy was tested by comparison with measurements of the line impedance. The technique was found to be accurate within the tolerances of the measurement and the film thickness variations.

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3.3.4 Five-Wire Line

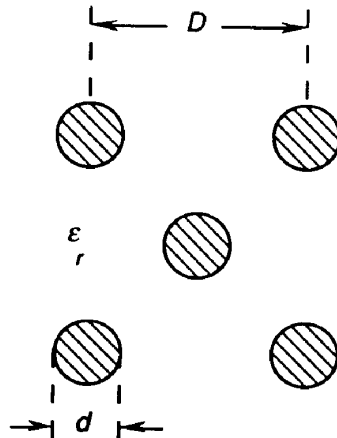


Figure 3.3.4.1: Five-Wire Line

$$Z_0 = \frac{2.5044 \eta_0}{\pi \sqrt{\epsilon_r}} \ln \left[\frac{D}{0.933 d} \right] \quad (\Omega) \quad (3.3.4)$$

where

$$d/D \ll 1.0$$

Of course, five conductors means lots of modes, so although the reference doesn't state it an assumption is that $(D - d)$ is small relative to a wavelength.

REFERENCES

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3.3.5 Paired Strips

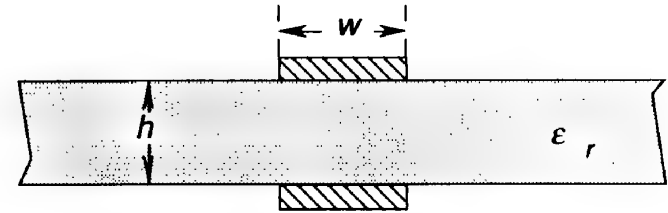


Figure 3.3.5.1: Paired Strips

For wide strips ($a/b > 1$):

$$Z_0 = \frac{\eta_0}{\sqrt{\epsilon_r}} \left\{ \frac{a}{b} + \frac{1.0}{\pi} \ln 4 + \frac{\epsilon_r + 1.0}{2 \pi \epsilon_r} \ln \left[\frac{\pi \epsilon (a/b + 0.94)}{2.0} \right] + \frac{\epsilon_r - 1.0}{2 \pi \epsilon_r^2} \ln \frac{\epsilon \pi^2}{16.0} \right\}^{-1} \quad (\Omega) \quad (3.3.5.1)$$

Stated error is less than 1% for wide strips.

For narrow strips ($a/b < 1$):

$$Z_0 = \frac{\eta_0}{\pi \sqrt{\epsilon_r}} \left[\ln \frac{4.0 b}{a} + \frac{1.0}{8.0} \left(\frac{a}{b} \right)^2 - \frac{\epsilon_r - 1.0}{2.0 (\epsilon_r + 1.0)} \left(\ln \frac{\pi}{2.0} + \frac{\ln \frac{4.0}{\pi}}{\epsilon_r} \right) \right] \quad (\Omega) \quad (3.3.5.2)$$

where

$$a = w/2.0 \quad (3.3.5.3)$$

$$b = h / 2.0 \quad (3.3.5.4)$$

Stated error is less than 1% for narrow strips.

Equations are valid for $2b$ much smaller than half a wavelength in the dielectric (λ_g).

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3.4 COPLANAR STRUCTURES

3.4.1 Coplanar Waveguide

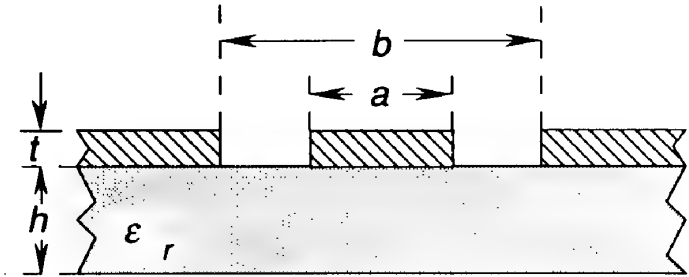


Figure 3.4.1.1: Coplanar Waveguide

The coplanar waveguide (CPW) configuration's advantages stem from its single-sided nature. Grounding components does not require plated through-holes to a plane on the other side of the substrate. This makes it ideal for use with surface mounted components. Another feature of coplanar waveguide is that we can narrow traces to match component lead widths while keeping Z_0 constant. The ground plane should extend greater than $5b$ on each side of the gap or the coplanar strip analysis should be used. The ground planes on either side of the center conductor may need to be connected periodically with wire jumpers depending on the frequency. The enclosure's cover can be used to jumper the two ground planes if it is kept away by at least $(b - a)$. To prevent propagation of higher modes, b should be less than $\lambda / 2$.

The design equations for coplanar waveguide are:

$$Z_0 = \frac{30.0 \pi K(k_1')}{\sqrt{\epsilon_{eff,t}} K(k_1)} \quad (3.4.1.1)$$

$$\epsilon_{eff,t} = \epsilon_{eff} - \frac{\epsilon_{eff} - 1.0}{\frac{(b-a)/2.0}{0.7t} \frac{K(k)}{K'(k)} + 1.0} \quad (3.4.1.2)$$

$$\epsilon_{eff} = 1.0 + \frac{\epsilon_r - 1.0}{2.0} \frac{K(k')K(k_1)}{K(k)K(k_1')} \quad (3.4.1.3)$$

$$k_t = \frac{a_t}{b_t} \quad k = \frac{a}{b} \quad (3.4.1.4)$$

$$k_t' = \sqrt{1.0 - k^2} \quad k' = \sqrt{1.0 - k^2} \quad (3.4.1.5)$$

$$k_1 = \frac{\sinh\left(\frac{\pi a_t}{4.0 h}\right)}{\sinh\left(\frac{\pi b_t}{4.0 h}\right)} \quad (3.4.1.6)$$

$$k_1' = \sqrt{1.0 - k_1^2} \quad (3.4.1.7)$$

$$a_t = a + \frac{1.25 t}{\pi} \left[1.0 + \ln\left(\frac{4.0 \pi a}{t}\right) \right] \quad (3.4.1.8)$$

$$b_t = b - \frac{1.25 t}{\pi} \left[1.0 + \ln\left(\frac{4.0 \pi a}{t}\right) \right] \quad (3.4.1.9)$$

These equations are corrected for thickness which is important for accurate results. The dielectric losses are described in [6, 11] by:

$$\alpha_d = \frac{q \epsilon_r \tan \delta}{\epsilon_{eff} \lambda_g} \quad (\text{Np/m}) \quad (3.4.1.10)$$

where q and λ_g have been previously defined.

The conductor losses are solved numerically and plotted in [6]. Jackson [9] reports that in some cases coplanar waveguide can have lower losses and dispersion than microstrip line for impedances near 50 Ω .

An equation for conductor losses [5] is:

$$\alpha_c = \frac{R_s \sqrt{\epsilon_{eff}} [\Phi(a) + \Phi(b)]}{480.0 \pi K(k) K(k') k'} \quad (\text{Np/m}) \quad (3.4.1.11)$$

where

$$R_s = \text{surface resistivity} = \sqrt{\pi f \mu_0 / \sigma} \quad (\Omega / \text{q}) \quad (3.4.1.12)$$

$$\Phi(x) = \frac{\pi}{x} + \frac{\ln\left[\frac{8.0 \pi x (1.0 - k)}{t (1.0 + k)}\right]}{x} \quad (3.4.1.13)$$

These equations are valid for

$$t \ll a$$

$$t \gg b - a$$

Radiation losses are

$$\alpha_r = \frac{\pi}{Q_{rad} \lambda_g} \quad (3.4.1.14)$$

$$Q_{rad} = \frac{K(k) K(k')}{\Psi(\epsilon_r, h, k_0, b) \Psi_{sc} \Psi_{oc}} \quad (3.4.1.15)$$

$$\begin{aligned} \Psi(\epsilon_r, h, k_0, b) &= (\epsilon_r - 1.0) \left[1.0 + 0.5 (\epsilon_r - 1.0)^2 (k_0 h)^2 \right] \\ &\quad \times (k_0 h) (k_0 b)^2 \left[1.0 + 0.25 (\epsilon_r - 1.0)^2 (k_0 h)^2 \right] \\ &\quad \times \left\{ \frac{1.0}{\sqrt{\epsilon_{eff}} \left[1.0 + (\epsilon_r - 1.0)^2 (k_0 h)^2 \right]} \right\} \end{aligned} \quad (3.4.1.16)$$

$$\begin{aligned} \Psi_{sc} &= \frac{5.0 \pi}{8} \frac{1.0 + \left(1.0 - \frac{\pi^2}{8.0}\right) \left[1.0 + (\epsilon_r - 1.0)^2 (k_0 h)^2 / 4.0 \right]}{3.0 \epsilon_{eff}} \\ &\quad \times \frac{\left[1.0 + (\epsilon_r - 1.0)^2 (k_0 h)^2 / 4.0 \right]}{\epsilon_{eff}} \end{aligned} \quad (3.4.1.17)$$

$$\Psi_{oc} = \frac{\pi}{8} \frac{3.0 + \left(1.0 - \frac{\pi^2}{8.0}\right) \left[1.0 + (\epsilon_r - 1.0)^2 (k_0 h)^2 / 4.0 \right]}{\epsilon_{eff}} \quad (3.4.1.18)$$

These are valid for thin dielectric.

In Houdart [8, Figure 2], the effect of ground plane width (c in Figure 3.4.8.1) on the CPW is analyzed. For $c/b > 5.0$ the impedance is affected by less than about 3%.

3.4.1.1 Example: Hybrid IC Probe in Coplanar Waveguide

What are the required dimensions for a 50 Ω coplanar waveguide probe to a 10 mil \times 10 mil bonding pad on a thick-film hybrid IC? The dielectric is 25.2 mil thick 96% alumina.

Solution: Assuming the probe is the same material as the hybrid IC substrate and using the above equations with

$$h = 0.0252 \text{ in} = 0.0640 \text{ cm}$$

$$a = 0.010 \text{ in} = 0.0254 \text{ cm}$$

$$t = 0.001 \text{ in} = 0.003 \text{ cm}$$

$$\epsilon_r = 10.0$$

The program iteratively adjusts b until the goal impedance 50Ω is achieved. The final dimension of the probe is:

$$b = 0.0234 \text{ in.} = 0.615 \text{ cm}$$

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3.4.2 Micro-Coplanar Stripline

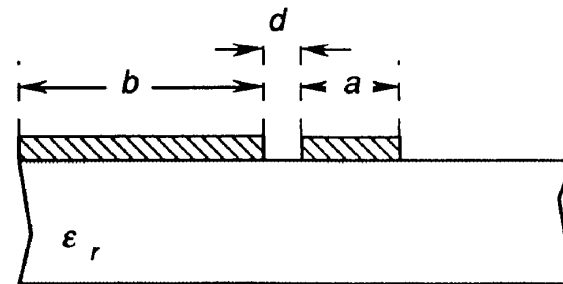


Figure 3.4.2.1: Micro-Coplanar Stripline

This structure has been called nonsymmetrical coplanar waveguide; however, to avoid confusion we prefer to reserve this term for coplanar waveguide having unequal gaps.

Kneppo and Gotzman [2] proposed this structure as an improvement over CPW for the connection of shunt elements. Their equations were derived with conformal transformation:

$$Z_0 = \frac{\eta_0}{(\sqrt{\epsilon_1} + \sqrt{\epsilon_2})} \frac{K(k)}{K'(k)} \quad (\Omega) \quad (3.4.2.1)$$

$$\epsilon_{eff} = \frac{(\sqrt{\epsilon_1} + \sqrt{\epsilon_2})^2}{4.0} \quad (3.4.2.2)$$

where

$$k = \sqrt{\frac{1.0 + a/d + b/d}{(1.0 + b/d)(1.0 + a/d)}} \quad (3.4.2.3)$$

Polynomial equations for the line width, ϵ_{eff} , and conductor losses are given in Quian and Yamashita [3] for dielectrics with $\epsilon_r = 2.22, 9.7, 10.1$, and 12.9 and in Yamashita *et al.* [4] for $\epsilon_r = 12.7$.

Losses may be calculated with the incremental inductance rule.

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3.4.3 Coplanar Waveguide with Ground

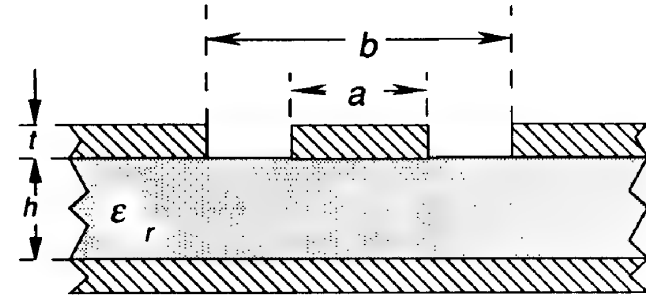


Figure 3.4.3.1: Coplanar Waveguide with Ground

The equations of this section can be used to analyze coplanar waveguide with ground or for microstrip lines with signal side ground plane.

$$Z_0 = \frac{60.0 \pi}{\sqrt{\epsilon_{eff}}} \frac{1.0}{\frac{K(k)}{K'(k)} + \frac{K(k_1)}{K'(k_1')}} \quad (3.4.3.1)$$

$$k = a/b \quad (3.4.3.2)$$

$$k' = \sqrt{1.0 - k^2} \quad (3.4.3.3)$$

$$k_1' = \sqrt{1.0 - k_1^2} \quad (3.4.3.4)$$

$$k_1 = \frac{\tanh\left(\frac{\pi a}{4.0 h}\right)}{\tanh\left(\frac{\pi b}{4.0 h}\right)} \quad (3.4.3.5)$$

$$\epsilon_{eff} = \frac{1.0 + \epsilon_r \frac{K(k')}{K(k)} \frac{K(k_1)}{K(k_1')}}{1.0 + \frac{K(k')}{K(k)} \frac{K(k_1)}{K(k_1')}} \quad (3.4.3.6)$$

The notation varies between references. Some use the a, b notation of Figure 3.4.2.1 or a $2a$ and $2b$ version while others use s (spacing to the adjacent ground) and w (trace width).

These equations "show good agreement" with spectral-domain variational calculation techniques.

The structure may propagate in three different modes: microstrip, coplanar waveguide, and coupled slotlines. To prevent slotline modes, jumpers connecting the two halves of the component side ground plane can be used. To avoid microstrip line modes, it is recommended [2] that $h \gg b$ and that the component side ground extend away from the trace on each side more than b .

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3.4.4 Shielded Coplanar Waveguide

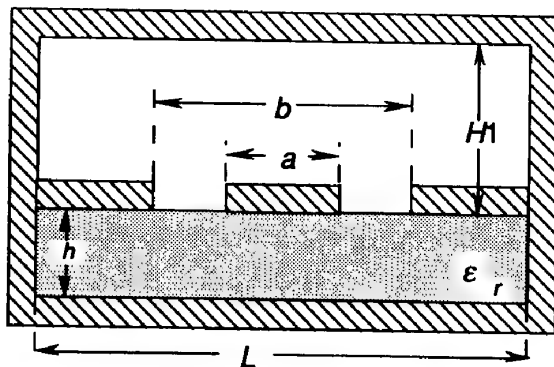


Figure 3.4.4.1: Shielded Coplanar Waveguide

The formulas below allow calculation of coplanar waveguide including the shield's effects [1, 2]. Many times we are simply looking for guidelines on how not to be affected by the shield's presence. For $L/b \geq 1.75$, and $H_1/a \geq 2.50$, Z_0 is affected less than 1.5% by the shield's presence.

The equations of [2] assume H_1 is infinite and are:

$$Z_0 = \frac{1.0}{\frac{5.0 q}{Z_m (1.0 + 5.0 q)} + \frac{1.0}{Z_c (1.0 + q)}} \quad (\Omega) \quad (3.4.4.1)$$

$$q = \frac{a}{h} \left(\frac{b}{a} - 1.0 \right) \left[3.6 - 2 e^{-(\epsilon_r + 1.0) / 4.0} \right] \quad (3.4.4.2)$$

$$Z_c = \text{coplanar waveguide impedance} \quad (3.4.4.3)$$

$$Z_m = \text{microstrip impedance} \quad (3.4.4.4)$$

Use the equations given in the coplanar waveguide and microstrip line sections together with the relevant dimensions to calculate Z_c and Z_m . Accuracy is within 2% of numerical results.

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3.4.5 Asymmetric Coplanar Waveguide

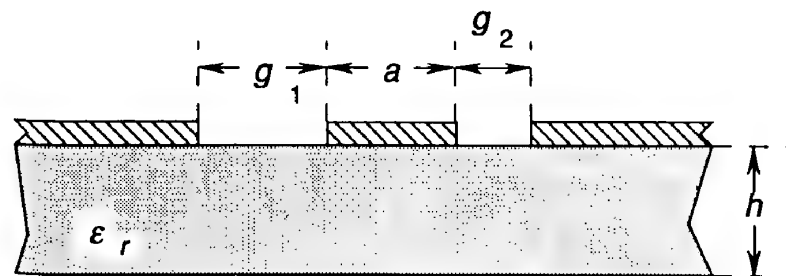


Figure 3.4.5.1: Asymmetric Coplanar Waveguide

The asymmetric coplanar waveguide structure can be used to reduce the line impedance of symmetric coplanar waveguide or to analyze the effects of fabrication tolerances.

The equations of Hanna and Thebault [1, 2] derived with conformal mapping are

$$Z_0 = \frac{30.0 \pi}{\sqrt{\epsilon_{eff}}} \frac{K'(k_1)}{K(k_1)} \quad (\Omega) \quad (3.4.5.1)$$

$$\epsilon_{eff} = 1.0 + \frac{\epsilon_r - 1.0}{2.0} \frac{K(k_2)}{K'(k_2)} \frac{K'(k_1)}{K(k_1)} \quad (3.4.5.2)$$

where

$$k_1 = \frac{0.5 b [1.0 + \alpha (0.5 b + d_1)]}{0.5 b + d_1 \alpha \sqrt{0.5 b}} \quad (3.4.5.3)$$

$$k_2 = \frac{w_A (1.0 + \alpha_1 w_B)}{w_B + \alpha_1 w_A^2} \quad (3.4.5.4)$$

and

$$w_A = \sinh \left(\frac{\pi a}{4.0 h} \right) \quad (3.4.5.5)$$

$$w_B = \sinh \left[\frac{\pi (a / 2.0 + g_1)}{2.0 h} \right] \quad (3.4.5.6)$$

$$w_E = -\sinh \left[\frac{\pi (a / 2.0 + g_2)}{2.0 h} \right] \quad (3.4.5.7)$$

$$\alpha = \frac{d_1 d_2 + 0.5 b (d_1 + d_2) \pm \sqrt{d_1 d_2 (b + d_1) (b + d_2)}}{\sqrt{0.5 b (d_1 - d_2)}} \quad (3.4.5.8)$$

$$\alpha_1 = \left(\frac{1.0}{w_B + w_E} \right) \left[-1.0 - \frac{w_B w_E}{w_A^2} \pm \sqrt{\left(\frac{w_B^2}{w_A^2} - 1.0 \right) \left(\frac{w_E^2}{w_A^2} - 1.0 \right)} \right] \quad (3.4.5.9)$$

The equations assume that the traces have negligible thickness. Comparison of the equations to experimental data showed agreement to within 4% for Z_0 . Choose the '-' sign solution of Equations (3.4.5.8) and (3.4.5.9).

Graphs of the dispersion of this structure are available in [3].

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3.4.6 Coplanar Strips

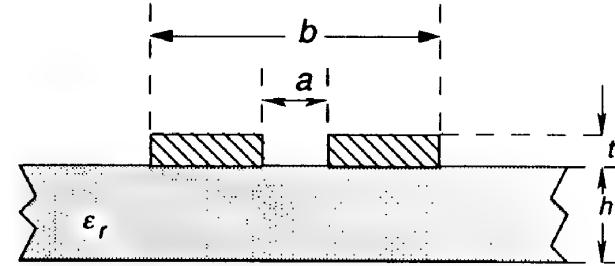


Figure 3.4.6.1: Coplanar Strips

The coplanar strips structure is similar to paired wire transmission line structures. Note also that coplanar strips are the complementary structure to coplanar waveguide.

$$Z_0 = \frac{\eta_0}{\sqrt{\epsilon_{eff}}} \frac{K(k)}{K'(k')} \quad (3.4.6.1)$$

$$\epsilon_{eff} = 1 + \frac{\epsilon_r - 1}{2} \frac{K(k')}{K(k)} \frac{K(k_1)}{K'(k_1')} \quad (3.4.6.2)$$

$$k = \frac{a}{b} \quad (3.4.6.3)$$

$$k' = \sqrt{1.0 - k^2} \quad (3.4.6.4)$$

$$k_1' = \sqrt{1.0 - k_1^2} \quad (3.4.6.5)$$

$$k_1 = \frac{\sinh\left(\frac{\pi a}{4h}\right)}{\sinh\left(\frac{\pi b}{4h}\right)} \quad (3.4.6.6)$$

Losses for this structure are [4]

$$\alpha_c = 17.34 \frac{R_s}{Z_0} \frac{P'}{\pi a} \left(1.0 + \frac{b-a}{2.0a}\right) \times \left(\frac{\frac{1.25}{\pi} \ln \frac{4.0 \pi (b-a)}{2.0 t} + 1.0 + \frac{2.5 t}{\pi (b-a)}}{\left\{1.0 + \frac{b-a}{a} + \frac{1.25 t}{\pi a} \left[1.0 + \ln \frac{2.0 \pi (b-a)}{t}\right]\right\}^2} \right) \quad (\text{dB/m}) \quad (3.4.6.7)$$

$$\alpha_d = \frac{20 \pi}{\ln(10)} \frac{\epsilon_r}{\sqrt{\epsilon_{eff}}} \frac{q \tan \delta}{\lambda_0} \quad (\text{dB/unit length}) \quad (3.4.6.8)$$

for $0 \leq k \leq 0.707$

$$P' = \frac{k}{k'^{3/2} (1.0 - k')} \left[\frac{K(k)}{K'(k)} \right] \quad (3.4.6.9)$$

for $0.707 \leq k \leq 1.0$

$$P' = \frac{1.0}{\sqrt{k} (1.0 - k)} \quad (3.4.6.10)$$

Pintzos [7] gives plots of the dispersion characteristics of this line.

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3.4.7 Asymmetrical Coplanar Strips

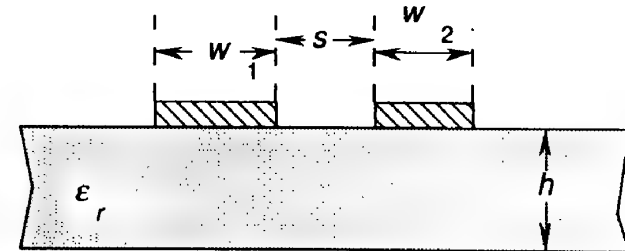


Figure 3.4.7.1: Asymmetrical Coplanar Strips

Hoffman [2] solves the equations for this structure from those of the infinitely thick dielectric structure. His equations are

$$Z_0 = \frac{\eta_0}{2.0 \sqrt{\epsilon_{eff}}} \frac{K(k)}{K'(k')} \quad (\Omega) \quad (3.4.7.1)$$

$$\epsilon_{eff} = 1.0 - \frac{(\epsilon_r - 1.0) K(k_1) K(k)}{2.0 K(k_1') K(k')} \quad (3.4.7.2)$$

where

$$k = \sqrt{\frac{s}{b} \left(1.0 + \frac{b}{a} - \frac{s}{d}\right)} \quad (3.4.7.3)$$

$$k_1 = \sqrt{\frac{(t_1 - t_2)(t_3 - t_2)}{(t_1 + t_2)(t_3 + t_2)}} \quad (3.4.7.4)$$

$$t_n = \frac{e^{\lambda_n} - 1.0}{e^{\lambda_n} + 1.0}, \quad n = 1, 2, 3 \quad (3.4.7.5)$$

$$\lambda_1 = \frac{\pi}{2.0} \left(\frac{2.0 w_2}{h} + \frac{s}{h} \right) \quad (3.4.7.6)$$

$$\lambda_2 = \frac{\pi s}{2.0 h} \quad (3.4.7.7)$$

$$\lambda_3 = \frac{\pi}{2.0} \left(\frac{2.0 w_1}{h} + \frac{s}{h} \right) \quad (3.4.7.8)$$

$$k_1' = \sqrt{1.0 - k_1^2} \quad (3.4.7.9)$$

$$k' = \sqrt{1.0 - k^2} \quad (3.4.7.10)$$

$$b = w_2 + s \quad (3.4.7.11)$$

$$d = w_1 + s \quad (3.4.7.12)$$

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3.4.8 Three Coplanar Strips

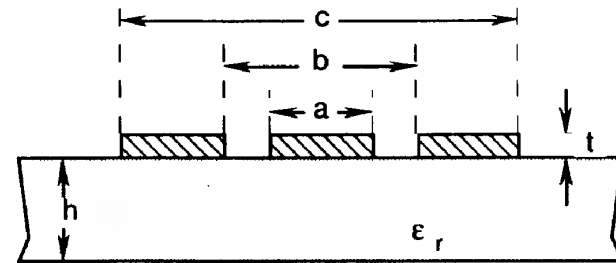


Figure 3.4.8.1: Three Coplanar Strips

This configuration was derived in [4] to analyze and design a tape automated bonding (TAB) IC package with alternate traces grounded.

$$Z_0 = \frac{\eta_0}{4.0 \sqrt{\epsilon_{eff}}} \frac{K(k_1)}{K(k_1')} \quad (\Omega) \quad (3.4.8.1)$$

where

$$\epsilon_{eff} = 1.0 + \frac{\epsilon_r - 1.0}{2.0} \frac{K(k_2')}{K(k_2)} \frac{K(k_1)}{K(k_1')} + \frac{\epsilon_r - 1.0}{2.0} \frac{K(k_2')}{K(k_2)} \left[\frac{K(k_1)}{K(k_1')} \right]^2 \left(\frac{t}{b-a} \right) + \frac{2.0 t}{b-a} \frac{K(k_1)}{K(k_1')} + \left[\frac{t}{b-a} \frac{K(k_1)}{K(k_1')} \right]^2 \quad (3.4.8.2)$$

$$k_1 = \frac{c}{b} \sqrt{\frac{b^2 - a^2}{c^2 - a^2}} \quad (3.4.8.3)$$

$$k_2 = \frac{\sinh\left(\frac{\pi c}{4.0 h}\right)}{\sinh\left(\frac{\pi b}{4.0 h}\right)} \sqrt{\frac{\sinh^2\left(\frac{\pi b}{4.0 h}\right) - \sinh^2\left(\frac{\pi a}{4.0 h}\right)}{\sinh^2\left(\frac{\pi c}{4.0 h}\right) - \sinh^2\left(\frac{\pi a}{4.0 h}\right)}} \quad (3.4.8.4)$$

$$k_n' = \sqrt{1 - k_n^2}, \quad n = 1, 2 \quad (3.4.8.5)$$

In Houdart [3, Figure 2], the effect of ground plane width (c in Figure 3.4.8.1) on the CPW is analyzed. For $c/b > 5.0$ the impedance is affected by less than about 3%.

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3.4.9 Three Coplanar Strips with Ground

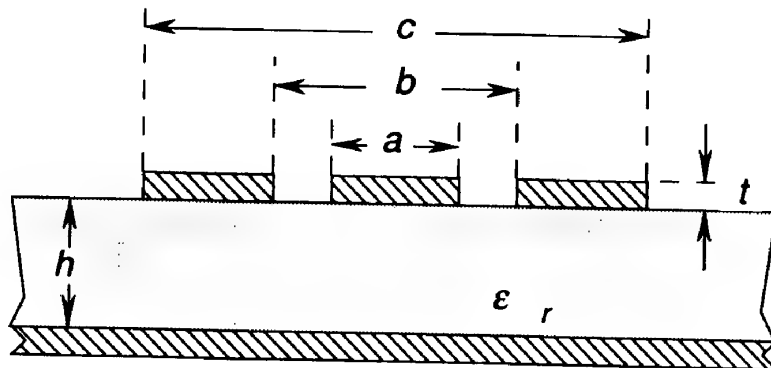


Figure 3.4.9.1: Three Coplanar Strips With Ground

This configuration was derived in [1] to analyze and design a tape automated bonding structure where alternate traces are grounded and are connected to a ground layer with vias.

$$Z_0 = \frac{\eta_0}{2.0 \sqrt{\epsilon_{eff}}} \frac{1.0}{\frac{K(k_1)}{K(k_1')} + \frac{K(k_2)}{K(k_2')} + \frac{2.0 t}{(b-a)}} \quad (3.4.9.1)$$

where

$$\epsilon_{eff} = \frac{\frac{K(k_1)}{K(k_1')} + \frac{\epsilon_r K(k_2)}{K(k_2')} + \frac{2.0 t}{b-a}}{\frac{K(k_1)}{K(k_1')} + \frac{K(k_2)}{K(k_2')} + \frac{2.0 t}{(b-a)}} \quad (3.4.9.2)$$

$$k_1 = \frac{c}{b} \sqrt{\frac{b^2 - a^2}{c^2 - a^2}} \quad (3.4.9.3)$$

$$k_2 = \frac{\tanh\left(\frac{\pi a}{4.0 h}\right)}{\tanh\left(\frac{\pi b}{4.0 h}\right)} \quad (3.4.9.4)$$

$$k_n' = \sqrt{1.0 - k_n^2}, n = 1, 2 \quad (3.4.9.5)$$

REFERENCES

- [1] Wentworth, Stuart M., *et al.*, "The High-Frequency Characteristics of Tape Automated Bonding (TAB) Interconnects," *IEEE Transactions on Components, Hybrids, and Manufacturing Technology*, Vol. CHMT-122, No. 3, September 1989, pp. 340–347.

3.4.10 Covered Coplanar Waveguide with Ground

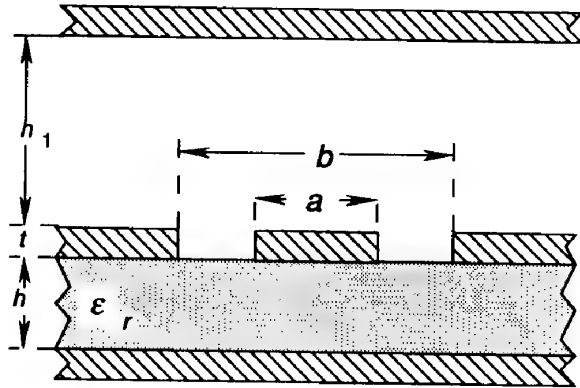


Figure 3.4.10: Covered Coplanar Waveguide with Ground

The equations describing this structure are

$$Z_0 = \frac{\eta_0}{2.0 \sqrt{\epsilon_{eff}}} \frac{1.0}{\frac{K(k_3)}{K(k'_3)} + \frac{K(k_4)}{K(k'_4)}} \quad (3.4.10.1)$$

$$\epsilon_{eff} = 1.0 + \frac{\frac{K(k_3)}{K(k'_3)}}{\frac{K(k_3)}{K(k'_3)} + \frac{K(k_4)}{K(k'_4)}} (\epsilon_r - 1.0) \quad (3.4.10.2)$$

where

$$k_3 = \frac{\tanh\left(\frac{\pi a}{h}\right)}{\tanh\left(\frac{\pi b}{h}\right)} \quad (3.4.10.3)$$

$$k_4 = \frac{\tanh\left(\frac{\pi a}{h_1}\right)}{\tanh\left(\frac{\pi b}{h_1}\right)} \quad (3.4.10.4)$$

which are strictly valid for

$$h = h_1$$

and were found to give good results elsewhere.

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3.4.11 Covered Coplanar Waveguide

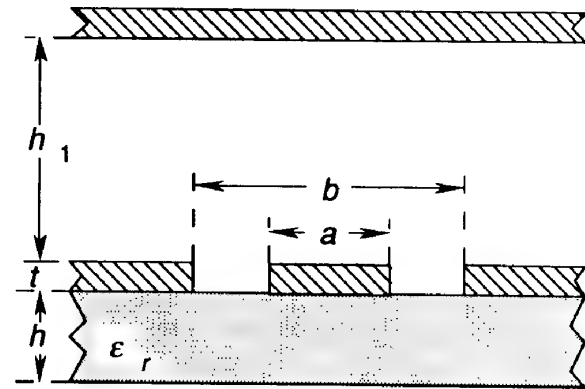


Figure 3.4.11.1: Covered Coplanar Waveguide

The equations of Ghione and Naldi [1] are

$$Z_0 = \frac{\eta_0}{2.0 \sqrt{\epsilon_{eff}}} \frac{1.0}{\frac{K(k_2)}{K(k'_2)} + \frac{K(k)}{K(k')}} \quad (3.4.11.1)$$

$$\epsilon_{eff} = 1.0 + \frac{\frac{K(k_1)}{K(k'_1)}}{\frac{K(k_2)}{K(k'_2)} + \frac{K(k)}{K(k')}} (\epsilon_r - 1.0) \quad (3.4.11.2)$$

where

$$k = a / b \quad (3.4.11.3)$$

$$k_1 = \frac{\sinh\left(\frac{\pi a}{h}\right)}{\sinh\left(\frac{\pi b}{h}\right)} \quad k_1' = \sqrt{1.0 - k_1} \quad (3.4.11.4)$$

$$k_2 = \frac{\tanh\left(\frac{\pi a}{h_1}\right)}{\tanh\left(\frac{\pi b}{h}\right)} \quad k_2' = \sqrt{1.0 - k_2} \quad (3.4.11.5)$$

These equations are valid to X-band for usual dimensions of PCBs, hybrids, and ICs.

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3.5 MICROSTRIP LINE STRUCTURES

Microstrip line may well be the most popular transmission line structure. Ease of fabrication by photolithographic techniques and a good range of impedances and couplings allow it to be used for a wide variety of circuit components. Harold Wheeler developed a planar transmission line (two coplanar strips) which could be rolled up in 1936 and a stripline-like structure in 1942. Flat coaxial transmission line was used by V.H. Rumsey and H.W. Jamieson in WWII. Coaxial cable was first adapted to a flat configuration using printed circuit techniques by Barrett. This then evolved into stripline after WWII. The first use of the microstrip line configuration was reported by engineers at the Federal Telecommunications Research Laboratories (a division of ITT) sometime after 1949 [2].

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3.5.1 Microstrip Line

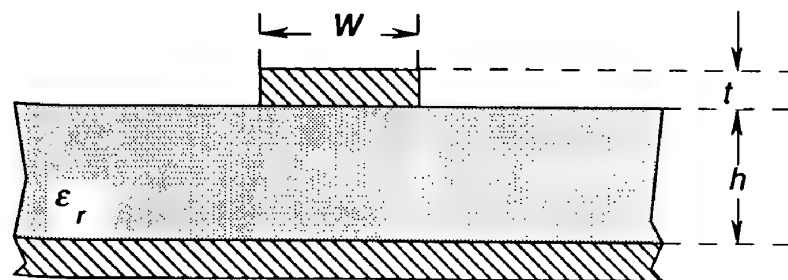


Figure 3.5.1.1: Microstrip Line

This is probably the most common and most analyzed transmission line structure. It is easy to use and has a good range of practical impedances. Equations have been calculated with an incredible variety of techniques both analytic and computational. A FORTRAN

program, MSTRIP2, is commonly encountered in references as a tool for checking the results of new analysis and synthesis equations. It is a numerical analysis program that assumes quasistatic conditions, zero thickness strips, and perfect conductivity. It is also assumed that the dielectric thickness and trace widths are thin relative to a wavelength.

Bogatin [10] experimentally compared various calculation techniques for this structure and recommends using the Wheeler equations with Schneider's ϵ_{eff} .

$$Z_0 = \frac{\eta_0}{2.0 \sqrt{2.0} \pi \sqrt{\epsilon_r + 1.0}} \ln \left\{ 1.0 + \frac{4.0 h}{w'} \left[\frac{14.0 + 8.0 / \epsilon_r}{11.0} \frac{4.0 h}{w'} + \sqrt{\left(\frac{14.0 + 8.0 / \epsilon_r}{11.0} \right)^2 \left(\frac{4.0 h}{w'} \right)^2 + \frac{1.0 + 1.0 / \epsilon_r}{2.0} \pi^2} \right] \right\} (\Omega) \quad (3.5.1.1)$$

Improvements in Schneider's ϵ_{eff} made by Hammerstad and Bekkadal [22] are given here. For $w/h \leq 1.0$:

$$\epsilon_{eff} = \frac{\epsilon_r + 1}{2} + \frac{\epsilon_r - 1}{2} \left[\left(1 + \frac{12 h}{w} \right)^{-0.5} + 0.04 \left(1.0 - \frac{w}{h} \right)^2 \right] \quad (3.5.1.2)$$

and for $w/h \geq 1$:

$$\epsilon_{eff} = \frac{\epsilon_r + 1.0}{2.0} + \frac{\epsilon_r - 1.0}{2.0} \left(1 + \frac{12.0 h}{w} \right)^{-0.5} \quad (3.5.1.3)$$

The equations for ϵ_{eff} are accurate to within 1% for:

$$\epsilon_r \leq 16 \quad (< 2\% \text{ error } \epsilon_r > 16)$$

$$0.05 \leq \frac{w}{h} \leq 20.0 \quad (< 2\% \text{ error } \frac{w}{h} < 0.05)$$

The thickness of the trace can be corrected for by relating it to an equivalent change in the width. Owens and Potok [40] examined a number of formulas for this correction and show that Wheeler's is the most accurate:

$$\frac{\Delta w}{t} = \frac{1.0}{\pi} \ln \left[\frac{4 e}{\sqrt{(t/h)^2 + \left(\frac{1/\pi}{w/t + 1.1} \right)^2}} \right] \quad (3.5.1.4)$$

$$w' = w + \Delta w' \quad (3.5.1.5)$$

$$\Delta w' = \Delta w \left(\frac{1.0 + 1.0 / \epsilon_r}{2.0} \right) \quad (3.5.1.6)$$

Error in Z_0 is less than 2% for any ϵ_r, w .

3.5.1.1 Frequency Dependencies of Microstrip Line

The frequency dependency (dispersion) of the microstrip line's ϵ_{eff} can be calculated with (see also [35]):

$$\epsilon_{eff}(f) = \epsilon_r - \frac{\epsilon_r - \epsilon_{eff}(f=0)}{1.0 + P(f)} \quad (3.5.1.7)$$

$$P(f) = P_1 P_2 [(0.1844 + P_3 P_4) \times (10.0 f h)]$$

$$P_1 = 0.27488 + \left[0.6315 + \frac{0.525}{(1.0 + 0.157 f h)^{20}} \right] \frac{w}{h} - 0.065683 e^{-8.7513 w/h} \quad (3.5.1.8)$$

$$P_2 = 0.33622 [1.0 - e^{-0.03442 \epsilon_r}] \quad (3.5.1.9)$$

$$P_3 = 0.0363 e^{-4.6 w/h} [1.0 - e^{-(f h / 3.87)^{4.97}}] \quad (3.5.1.10)$$

$$P_4 = 1.0 + 2.751 [1.0 - e^{-(\epsilon_r / 0.916)^8}] \quad (3.5.1.11)$$

where f is in GHz and h is in cm. The accuracy of this correction is better than 0.6% for

$$0.1 \leq w/h \leq 100.0$$

$$1.0 \leq \epsilon_r \leq 20.0$$

$$0 \leq h/\lambda \leq 0.13$$

Atwater [2] compares several corrections for $\epsilon_{eff}(f)$ with actual measurement data from eight published papers and found the above correction to be the best (by an admittedly small margin). His version of the equation is a rearranged equivalent.

3.5.1.2 Conductor Losses

Conductor losses are a result of several contributing factors: the actual conductance of the material, the frequency-dependent skin effect losses, and the surface roughness losses

caused by the lengthened path at the surface. Conductor losses due to skin effect and the metal's conductivity are given in [53]. For $w/h \leq 1$:

$$\alpha_c = \frac{10.0 R_s}{\pi \ln 10.0} \left(\frac{8.0 h}{w} - \frac{w}{4 h} \right) \left(1.0 + \frac{h}{w} + \frac{h}{w} \frac{\partial w}{\partial t} \right) \frac{Z_0 / 60}{h Z_0 e} \quad (\text{dB/unit length}) \quad (3.5.1.12)$$

For $w/h \geq 1$:

$$\alpha_c = \frac{Z_0 R_s}{720.0 \pi^2 h \ln 10.0} \left[1.0 + \left(\frac{0.44 h^2}{w^2} \right) + \frac{6.0 h^2}{w^2} \left(1.0 - \frac{h}{w} \right)^5 \right] \times \left(1.0 + \frac{w}{h} + \frac{\partial w}{\partial t} \right) \quad (\text{dB/unit length}) \quad (3.5.1.13)$$

where for $\frac{w}{h} \leq \frac{1}{2\pi}$

$$\frac{\partial w}{\partial t} = \left(\frac{1.0}{\pi} \right) \ln \frac{4.0 \pi w}{t} \quad (3.5.1.14)$$

for $\frac{w}{h} \geq \frac{1}{2\pi}$

$$\frac{\partial w}{\partial t} = \left(\frac{1.0}{\pi} \right) \ln \frac{2.0 h}{t} \quad (3.5.1.15)$$

These equations are valid for ground and strip conductivities the same and

$$t \ll h$$

$$t < h/2.0$$

$$\partial w / \partial t > 1.0$$

Pucel *et al.* [43] also give equations for the conduction losses.

3.5.1.2.1 Effect of Ground Plane on Conductor Losses

The ground plane's resistance can be a significant factor as frequency increases. At low frequencies the current in the ground plane is spread over the full width of the ground plane and the current is evenly distributed in the center conductor. Because of this, the ground plane contributes little resistance relative to the center conductor. As frequency

increases the current in the line begins to concentrate in the sides and bottom of the trace. Simultaneously, the ground current begins to concentrate in the area directly under the strip, increasing the effective resistance. Faraji-Dana and Chow [16] give a closed-form equation suitable for design calculations. The resistance of the ground plane is R_g

$$R_g = 0.55 R_{dc} \sqrt{\frac{\pi}{2.0}} \sqrt{\frac{t}{w}} (1.0 - e^{-w' / 1.2 \pi}) P \quad (3.5.1.16)$$

where

$$w' = w/h \quad (3.5.1.17)$$

$$P = \sqrt{2.0 \mu \sigma f w t} \quad (3.5.1.18)$$

R_{dc} = the dc resistance of the signal conductor

The total ac resistance for use in the microstrip line model is R_t .

$$R_t = R_s + R_g \quad (3.5.1.19)$$

3.5.1.2.2 Effect of Conductor Surface Roughness

The conductor surface analyzed in the formulas above are perfectly smooth. The physical processes which make real world conductors create scratches, bumps, and random bumps. See Tanaka and Okada [62] for some scanning microscope photographs of typical conductor surfaces.

Conductor surface roughness losses are discussed thoroughly in Sanderson [53]. Sanderson also reports on a technique used for calculating the roughness effect on the losses, Z_0 , and v_p . Higher than expected losses were experienced in stainless steel conductors that can be modeled by a lumped resistance. Good agreement between the predicted performance and the experimental measurements was found from 200 MHz to 10 GHz in the other cases. Saad [50] gives a plot from Lending [32] for losses as a function of surface roughness. In [20], an approximate formula is given for the conductor losses due to roughness:

$$\alpha_c = \alpha_{c,0} \left\{ 1.0 + \frac{2.0}{\pi} \tan^{-1} \left[1.40 \left(\frac{\Delta}{\delta} \right)^2 \right] \right\} \quad (3.5.1.20)$$

where

$\alpha_{c,0}$ = the conductor losses calculated for a perfectly smooth conductor; equations (3.5.1.12) and (3.5.1.13)

Δ = the rms surface roughness

δ = the skin depth

The results of a number of references [20, 32, 39, 60, 66] are summarized in Figure 3.5.1.2 below. This figure shows good agreement between Morgan's data and the Goldfarb closed form equation.

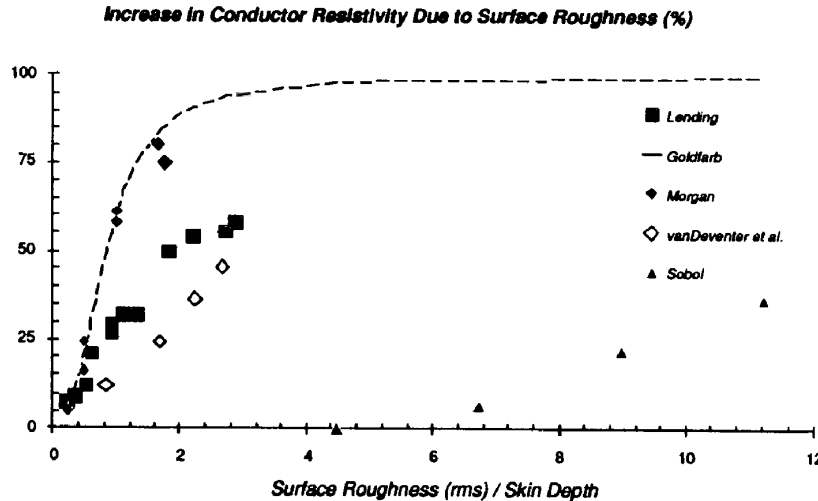


Figure 3.5.1.2.2.1: Microstrip Line Losses vs. Surface Roughness

3.5.1.3 Dielectric Losses

Dielectric losses are given in Schneider [50]:

$$\alpha_d = \frac{20.0 \pi q \tan \delta}{\ln(10.0) \lambda} \quad (\text{dB / unit length}) \quad (3.5.1.21)$$

where q is the filling factor and λ_g is the wavelength in microstrip line both previously defined.

3.5.1.4 Radiation Losses

Radiative losses were discussed in [1]. The equations are:

$$\alpha_r = 60 \left(\frac{2 \pi h}{\lambda_0} \right)^2 F(\epsilon_{eff}) \quad (3.5.1.22)$$

where $F(\epsilon_{eff})$ is:

$$F(\epsilon_{eff}) = \frac{\epsilon_{eff} + 1.0}{\epsilon_{eff}} - \frac{(\epsilon_{eff} - 1.0)^2}{2.0 \epsilon_{eff}^{3/2}} \log \left(\frac{\sqrt{\epsilon_{eff} + 1.0}}{\sqrt{\epsilon_{eff} - 1.0}} \right), \text{ for an open-circuited line} \quad (3.5.1.23)$$

$$F(\epsilon_{eff}) = 1.0 - \frac{(\epsilon_{eff} - 1.0)}{2.0 \sqrt{\epsilon_{eff}}} \log \left(\frac{\sqrt{\epsilon_{eff} + 1.0}}{\sqrt{\epsilon_{eff} - 1.0}} \right), \text{ for a matched line} \quad (3.5.1.24)$$

3.5.1.5 Higher Order Modes in Microstrip Line

The microstrip line structure is not strictly TEM; other modes may propagate. The lowest order TE_1 mode frequency is [69]:

$$f_c = \frac{c}{4.0 h \sqrt{\epsilon_r - 1.0}} \quad (3.5.1.25)$$

These non-TEM modes are not to be confused with waveguide modes which can be propagated when a microstrip line is placed in a metallic enclosure for shielding. These modes are launched or stimulated by radiation or discontinuities in the transmission line. These modes can be suppressed by decreasing the enclosure dimensions, adding shorting posts, lossy films, or by damping materials. See [71, 72] for more information.

3.5.1.6 Magnetic Substrates

Magnetic substrates are useful for building voltage programmable transmission lines. This is useful, for example, in constructing programmable phase shifters. Magnetic substrates are analyzed in [42, 43] which propose the following modifications to the above:

$$Z_0 = Z_0(\epsilon_r = 1, \mu_r = 1, h, w, t) \sqrt{\frac{\mu_{eff}}{\epsilon_{eff}}} \quad (\Omega) \quad (3.5.1.26)$$

$$\mu_{eff}(w/h, \mu) = \frac{1.0}{\epsilon_{eff}(w/h, 1.0/\mu)} \quad (3.5.1.27)$$

where

$$Z_0(\epsilon_r = 1, \mu_r = 1, h, w, t) \text{ is given by (3.5.1.1)} \quad (3.5.1.28)$$

Equation (3.5.1.27) can be rewritten using a simpler version of (3.5.1.2) and (3.5.1.3)

$$\begin{aligned} \mu_{eff}(w/h, \mu) &= \frac{1.0}{\frac{1.0/\mu + 1}{2.0} + \frac{1.0/\mu - 1.0}{2} \left(1.0 + \frac{10.0 h}{w}\right)^{-0.5}} \\ &= \frac{2.0 \mu}{(1.0 + \mu) + (1.0 - \mu) \left(1.0 + \frac{10.0 h}{w}\right)^{-0.5}} \end{aligned} \quad (3.5.1.29)$$

The dielectric is magnetic, so we must use the complete equation for the guide wavelength, (2.3.11):

$$\lambda_g = \frac{\lambda_0}{\sqrt{\epsilon_{eff} \mu_{eff}}} \quad (3.5.1.30)$$

The dielectric losses are modified in the presence of a magnetic substrate to

$$\alpha_d = \frac{27.3}{\lambda_g} (\tan \delta_{d,eff} + \tan \delta_{m,eff}) \quad (\text{dB}) \quad (3.5.1.31)$$

where

$$\tan \delta_{d,eff} = \left(\frac{1.0 - 1.0/\epsilon_{eff}}{1.0 - 1.0/\epsilon_r} \right) \tan \delta = \left(\frac{\epsilon_{eff} \epsilon_r - \epsilon_r}{\epsilon_{eff} \epsilon_r - \epsilon_{eff}} \right) \tan \delta \quad (3.5.1.32)$$

$$\tan \delta_{m,eff} = \left(\frac{1.0 - \mu_{eff}}{1.0 - \mu_r} \right) \tan \delta_m \quad (3.5.1.33)$$

The magnetic loss tangent can be calculated [20]:

$$\tan \delta_m = \frac{\omega_m / T}{\omega^2} \quad (3.5.1.34)$$

where

$$T = \frac{2.0}{\gamma \Delta H} = \text{relaxation time} \quad (3.5.1.35)$$

$$\gamma = \frac{g e}{m c} \quad (3.5.1.36)$$

ΔH = resonance-line width

M_s = saturation magnetization (Wb / m²)

g = spectroscopic splitting factor

e = electron charge

m = electron mass

$\omega_m = \gamma M_s$ = resonance frequency of the ferrite

ω = operating frequency

3.5.1.7 Conductors of Trapezoidal Cross-section

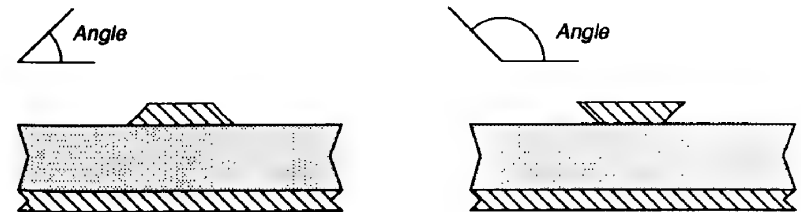


Figure 3.5.1.7.1: Trapezoidal Microstrip Line Traces

Conventionally it is assumed that the microstrip line conductor has a rectangular cross-section, however this is only an approximation to the physical reality. The additive and subtractive processes used to create conductors for PCBs, thin-, and thick-film hybrids produce a conductor of approximately trapezoidal cross-section. For wide strips (large w/h) the thickness and edge effects will be small since most of the capacitance is parallel plate rather than fringing and we can ignore the cross-section. For narrow strips which are common in smaller structures this condition is violated and the thickness and edge shape become important.

Information on this structure is found in references [7, 36, 47, 48, 49, 56]. Barsotti, *et al.* [7] develop a modified incremental inductance rule which can be used for calculating the loss of any shape conductor once a number of parameters are determined. It has been found in Rizzoli [49] that the trapezoidal cross-section results in a characteristic impedance

lying somewhere between that of a rectangular cross-section (having width equal to the widest side of the trapezoid) and a zero thickness conductor.

The program CALIF¹ was used to calculate the effect of trapezoidal cross-section on a microstrip line. The width of the wider side was kept constant and the angle is as shown in Figure 3.5.1.7.1 above. The case examined was:

$$w = 18.02 \text{ mil } (w/h = 1.82)$$

$$t = 1.4 \text{ mil } (t/h = .14, t/w = .078)$$

$$h = 9.9 \text{ mil}$$

$$\epsilon_r = 4.30$$

The limiting cases define the bounds of the impedance effects. The rectangular trace having width w is the lowest impedance we can have (51.58 Ω). The zero-thickness trace is the highest impedance we can achieve with etching of the type shown in the leftmost figure of Figure 3.5.1.7.1 (53.30 Ω). The highest impedance we can achieve with etching similar to the right-hand figure of Figure 3.5.1.7.1 is that of a zero-thickness trace suspended by t over the dielectric substrate (70.20 Ω). The actual impedance variation between these cases is shown in Figure 3.5.1.7.2.

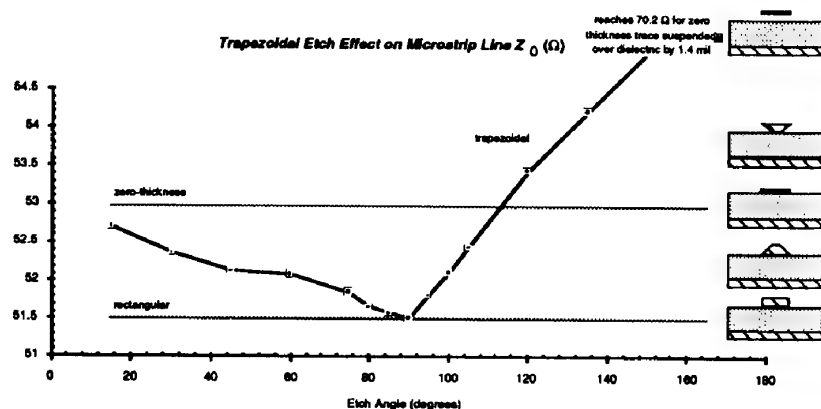


Figure 3.5.1.7.2: Trapezoidal Trace Effect on Z_0 ($t/h = 0.14$)

It can be expected that the increase of t/h (more square trace) will move the two lines labeled *rectangular* and *zero-thickness* relatively further apart on the plot above. In fact, it is seen in the plot below that the asymptotic lines are indeed further apart.

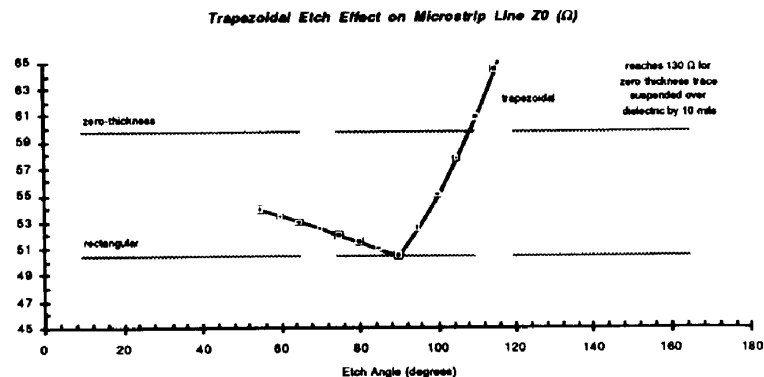


Figure 3.5.1.7.3: Trapezoidal Trace Effect on Z_0 ($t/h = 1.01$)

Based on the two extreme cases shown in Figure 3.5.1.7.2 and 3.5.1.7.3, it appears that calculating the effect of the trapezoidal trace (wide side near dielectric) on Z_0 by linearly interpolating the zero-thickness and rectangular impedances with a line over the 0 to 90° range. This is valid up until the point that the trace becomes an equilateral triangle. At this point, the thickness of the limiting rectangular trace can be decreased and the impedance interpolated in the same fashion. The equation is

$$Z_{0, \text{trapezoid}} = Z_0(t=0) - \frac{\theta}{90^\circ} [Z_0(t=0) - Z_0(t)] \quad (3.5.1.7.1)$$

which is valid for angles, θ

$$\tan^{-1} \left(\frac{2.0 t}{w} \right) \leq \theta \leq 90.0^\circ \quad (3.5.1.7.2)$$

For the trapezoidal trace with the wide side further from the dielectric the relation is not linear. The impedance reaches the zero-thickness impedance by approximately 110°. Again, a linear interpolation will provide reasonable results over the 90 to 110° range.

The effective dielectric constant is plotted for the above two cases in figures 3.5.1.7.4 and 3.5.1.7.5 below.

¹CALIF is a registered trademark of a product of Integrity Engineering. The program (version 2.0) uses the boundary element method to solve a number of transmission line structures.

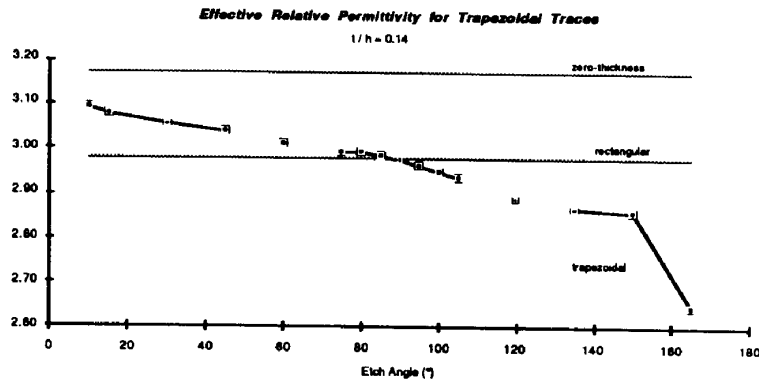


Figure 3.5.1.7.4: ϵ_r for Trapezoidal Trace ($t/h = 0.14$)

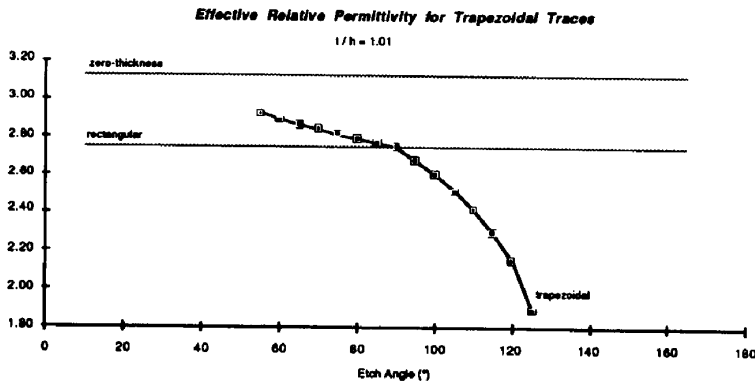


Figure 3.5.1.7.5: ϵ_r for Trapezoidal Trace ($t/h = 1.01$)

3.5.1.8 Superconducting Microstrip Line Losses

A relation for conductor losses in microstrip line can be found in Baiocchi, *et al.* [5].

$$\alpha_c = \frac{\theta}{(1.0 - \theta)^{3/2}} \frac{G \sigma_n \mu^2 \lambda^3}{4.0 Z_0}$$

$$\cdot \left(\coth \frac{A G \sqrt{1.0 - \theta}}{\lambda} + \frac{A G \sqrt{1.0 - \theta}}{\lambda} \operatorname{csech}^2 \frac{A G \sqrt{1.0 - \theta}}{\lambda} \right) \omega^2 \quad (3.5.1.8.1)$$

where

σ_n = conductivity near critical temperature

λ = zero temperature penetration depth

T_c = critical temperature of the superconductor

$$\theta = (T/T_c)^4 \quad (3.5.1.8.2)$$

and the value of G can be found with the relations in [29]:

$$G = \frac{1.0}{\mu} \sum_j \frac{\partial L}{\partial n} \quad (\text{m}^{-1}) \quad (3.5.1.8.3)$$

where

$\frac{\partial L}{\partial n}$ = derivative of external inductance for a small recession of the j th conductor wall (see Chapter 1).

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3.5.2 Microstrip Line with Truncated Ground Plane and Dielectric

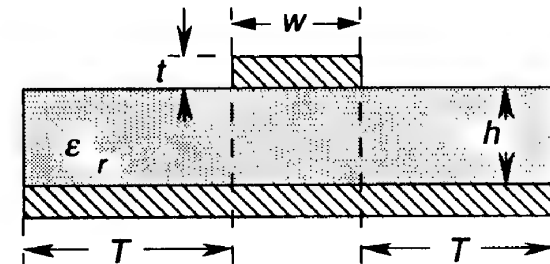


Figure 3.5.2.1: Microstrip Line with Truncated Ground Plane and Dielectric

It is desirable to know the applicable range of the standard microstrip line formulas for finite width grounds and dielectrics. Figure 2 of [1] indicates that $T/w \geq 2$ has < 3% effect on Z_0 . Figure 5 of [2] plots ϵ_{eff} as a function of T/w . As expected ϵ_{eff} decreases with T/w reaching a final value approximately 18% less than $T = \infty$ (for $\epsilon_r = 10.2$, $w/h = 0.73$, 2 GHz). A value of $T/w > 2$ has little effect on ϵ_{eff} .

This effect is a function of frequency. At dc the ground current will be spread equally through the ground cross-section; however, as frequency increases it will concentrate into a strip directly below the signal conductor.

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3.5.3 Microstrip Line with Truncated Dielectric

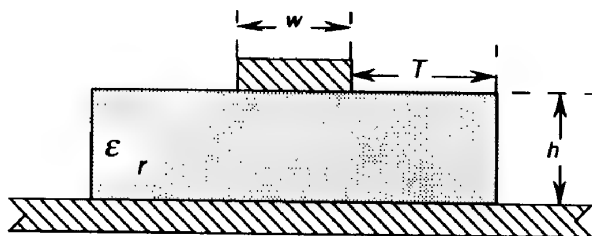


Figure 3.5.3.1: Microstrip Line with Truncated Dielectric

It is desirable to know the range of standard microstrip line equations in the presence of finite width dielectrics. Surprisingly, Figure 2 of [2] indicates for $T/w > 0.5$, Z_0 is affected $< 0.5\%$. This is because most of the field lines are concentrated in the dielectric directly under the microstrip line.

It should be pointed out that this effect is a function of frequency. At dc the ground current will be spread equally through the ground cross-section; however, as frequency increases it will concentrate into a strip directly below the signal conductor.

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3.5.4 Embedded or Buried Microstrip Line

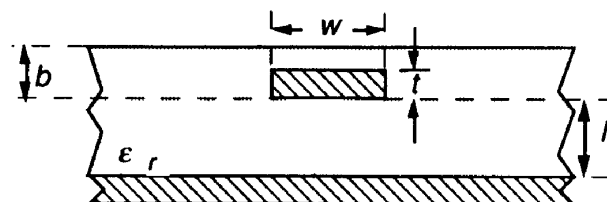


Figure 3.5.4.1: Embedded or Buried Microstrip Line

This structure is encountered in multilayer PCBs, compensated couplers, and even in single-layer PCBs. Some vendors apply a thin coat of epoxy to the PCB to make the dielectric flush with the upper microstrip line surface. This buries surface imperfections (loose glass fibers) remaining after processing.

Many simple approximations are available in the literature. These are included here due to their frequent occurrence in the PCB industry literature. This is not meant to imply their usefulness. The following approximation equation is given by [5] and others:

$$Z_0 = Z_0' - b\% \quad (3.5.4.1)$$

b (mils)

The impedance Z_0' is calculated as for normal microstrip line. The formula then implies that the impedance is reduced 1% per mil of dielectric over the line. In Coombs [4], it is stated that the presence of upper dielectric material reduces Z_0 by 22%. This rule and equation (3.5.4.1) are often quoted by PCB vendors and are in fact not very good estimators.

An equation which appears to have the proper form is found in [6]. The equations are:

$$Z_0 = \frac{Z_{0,\text{microstrip}} \sqrt{\epsilon_{\text{eff},\text{microstrip}}}}{\sqrt{\epsilon_{\text{eff},\text{buried}}}} \quad (3.5.4.2)$$

where,

Z_0 = characteristic impedance of a (not buried, $b = 0$) microstrip line

$\epsilon_{\text{eff},\text{microstrip}}$ = effective relative dielectric constant of a (not buried, $b = 0$) microstrip line

$$\epsilon_{eff,buried} = \epsilon_{eff} e^{(-2.0 b / h)} + \epsilon_r [1.0 - e^{(-2.0 b / h)}] \quad (3.5.4.3)$$

The figures below plot an example of the resulting ϵ_{eff} and $\epsilon_{eff,buried}$ for the microstrip line and embedded microstrip line configurations respectively. Note that the extra layer of dielectric material has more effect on narrow strips. In this example, for $b/h < .1$ there is $< 1\%$ effect on the line's ϵ_{eff} and, for $b/h > 1.0$, ϵ_{eff} is within 5% of ϵ_r . Static results of CALIF are also shown for $w/h = 1.0$. Narrow traces have more fringing capacitance and are effected more by the presence of an overlay than wider traces.

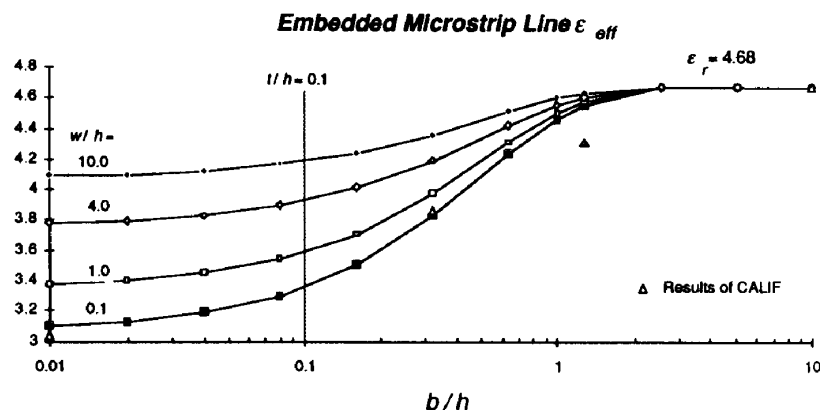


Figure 3.5.4.2.1: Comparison of Microstrip Line and Embedded Microstrip Line ϵ_{eff} 's

Embedded Microstrip Line Z_0 (Ω)

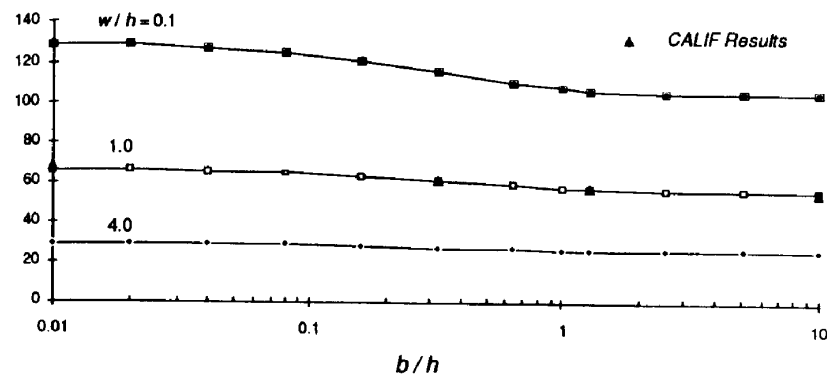


Figure 3.5.4.2.2: Embedded Microstrip Line Z_0

In Chang and Klein [3], data is presented for a line covered with a high permittivity line. The dependency on h_2 appears to be of the form $(1.0 - e^{-x})$, similar to the results of Buntschuh [2]. With equal dielectric constants the value of ϵ_{eff} was nearly ϵ_r for $h_2/h_1 = 2.0$.

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3.5.5 Nonhomogeneous Dielectric Embedded Microstrip Line

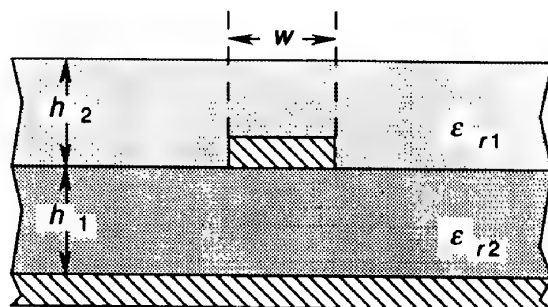


Figure 3.5.5.1: Nonhomogeneous Dielectric Embedded Microstrip Line or Microstrip Line with Overlay

This structure results when we cover a microstrip line with solder mask. Inner layer microstrip lines are sometimes used in multilayer PCBs where the dielectric constants are not equal. If $\epsilon_{r1} = \epsilon_{r2}$ use the equations of the previous section.

In lieu of design equations we can bound the problem. We know that

$$\epsilon_{\text{eff, microstrip}} < \epsilon_{\text{eff, embedded}} < \epsilon_{\text{eff, 2}}$$

where $\epsilon_{\text{eff, 2}}$ is calculated with the technique described in the Introduction.

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3.5.6 Shielded Microstrip Line

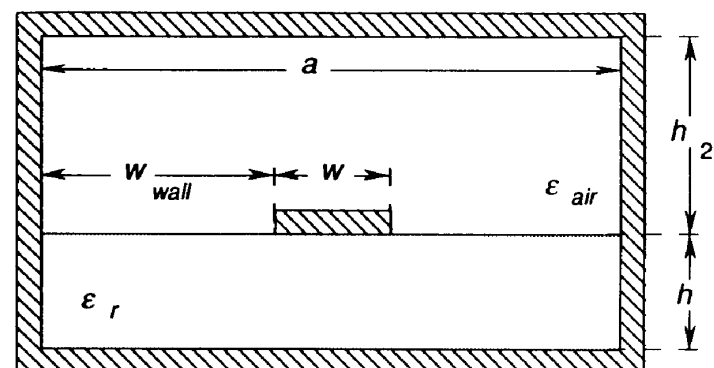


Figure 3.5.6.1: Shielded Microstrip Line

Shielding lowers $\epsilon_{\text{effective}}$ and Z_0 and creates resonances. Neglect the effects of shield for $W_{\text{wall}} > 5.0 \times w$ and $h_2 > 5.0 \times h$.

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3.5.7 Edge-Compensated Microstrip Line (ECM Line)

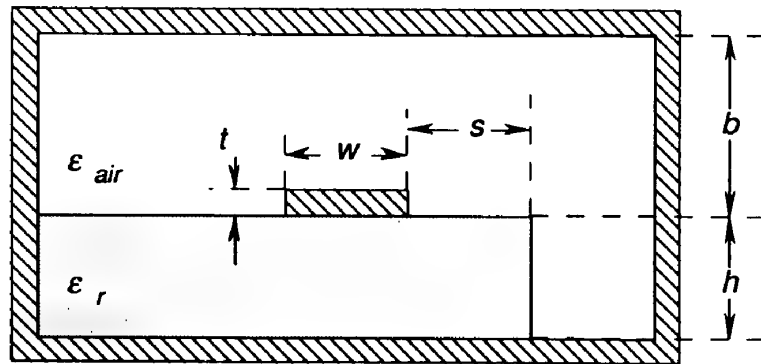


Figure 3.5.7.1: Edge-Compensated Microstrip Line

Yamashita *et al.* [3] show a technique for solving and correcting the effects of a shielded microstrip line near a dielectric truncation. The line is corrected by increasing the width calculated for the usual microstrip line, w_0 , by a factor w/w_0 and using a modified ϵ_{eff} . They also give simplified approximation formulas for GaAs and alumina microstrip lines. For GaAs ($\epsilon_r = 12.9$)

$$\frac{w}{w_0} = 1.0367 - 0.18228 u + 0.17147 u^2 + 0.13289 u^3 - 0.080565 u^4 - 0.081078 u^5 - 0.015946 u^6 \quad (3.5.7.1)$$

$$\epsilon_{eff} = (0.35346 - 0.025010 u + 0.022373 u^2 + 0.018016 u^3 - 0.010253 u^4 - 0.010717 u^5 - 0.00211593 u^6)^{-2} \quad (3.5.7.2)$$

For alumina ($\epsilon_r = 9.7$)

$$\frac{w}{w_0} = 1.0162 - 0.098654 u + 0.13236 u^2 + 0.033574 u^3 - 0.070952 u^4 - 0.024336 u^5 + 0.00022125 u^6 \quad (3.5.7.3)$$

$$\epsilon_{eff} = (0.39554 - 0.017792 u + 0.022627 u^2 + 0.0068338 u^3 - 0.012018 u^4 - 0.0046957 u^5 + 0.0019738 u^6)^{-2} \quad (3.5.7.4)$$

where

$$u = \log \frac{s}{w_0} \quad (3.5.7.5)$$

Equations are curve-fit polynomials that are within 1% of the theoretical results for

$$0.1 \leq s/w_0 \leq 5.0$$

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3.5.8 Covered Microstrip Line

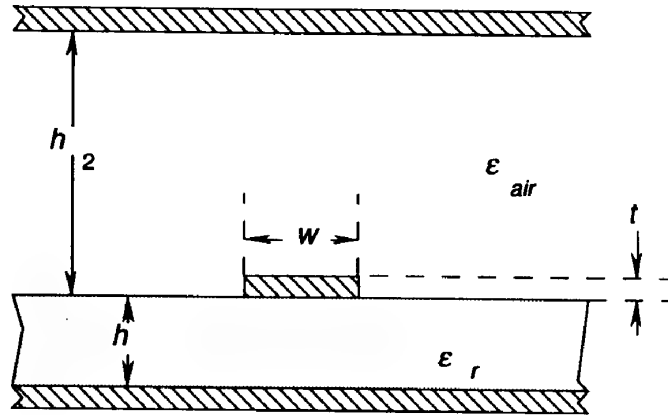


Figure 3.5.8.1: Covered Microstrip Line

The impedance of a microstrip line of a given set of dimensions is lowered by the presence of a cover at a finite distance h_2 . A large number of sources [1–8, 10] all describe this configuration. March [7] corrects a number of errors and compares his results with the “exact” MSTRIP analysis.

The technique is to calculate Z_0 for a completely air dielectric and a shield located at $h_2 = \infty$ and then correct this value for the presence of the cover with ΔZ_0 and ϵ_{eff} .

$$Z_0 = \frac{Z_0(\epsilon_r = 1, h_2 = \infty) - \Delta Z_0(\epsilon_r = 1)}{\sqrt{\epsilon_{eff}}} \quad (3.5.8.1)$$

The equation for Z_0 is after Hammerstad and Jensen [6]:

$$Z_0(\epsilon_r = 1, h_2 = \infty) = \frac{\eta_0}{2.0 \pi} \ln \left\{ \frac{f(w'/h)}{w'/h} + \sqrt{\left[1.0 + \left(\frac{2.0 h}{w'} \right)^2 \right]} \right\} \quad (3.5.8.2)$$

where

$$f(w'/h) = 6.0 + (2.0 \pi - 6.0) e^{-\left(\frac{30.666}{w'/h} \right)^{0.7528}} \quad (3.5.8.3)$$

These equations agree with MSTRIP to within $\pm 0.1\%$ for $0.1 \leq w/h \leq 6.0$ and to within $\pm 0.2\%$ for $6.0 \leq w/h \leq 10.0$.

$$w' = w + \frac{t \left\{ 1.0 + \ln 4.0 - 0.5 \ln \left[(t/h)^2 + \left(\frac{t}{\pi w} \right)^2 \right] \right\}}{\pi} \quad (3.5.8.4)$$

which is Wheeler's width correction [10] for finite strip thickness presented earlier. The value of $\Delta Z(\epsilon_r = 1)$ is calculated with corrected versions of Bahl's equations [1]:

$$\Delta Z(\epsilon_r = 1) = P Q \quad (3.5.8.5)$$

where

$$P = 270.0 \left[1.0 - \tanh \left(1.192 + 0.706 \sqrt{1.0 - h_2/h} - \frac{1.389}{1.0 + h_2/h} \right) \right] \quad (3.5.8.6)$$

$$Q = 1.0109 - \tanh^{-1} \left[\frac{0.012 w/h + 0.177 (w/h)^2 - 0.027 (w/h)^3}{(1.0 + h_2/h)^2} \right] \quad (3.5.8.7)$$

The effective dielectric constant corrects for the presence of the dielectric.

$$\epsilon_{eff} = \frac{\epsilon_r + 1}{2} + q \frac{\epsilon_r - 1}{2} \quad (3.5.8.8)$$

An adjusted value of q is used to calculate ϵ_{eff} . March [7] modifies equations of Bahl [1] by defining a q made up of three pieces:

$$q = q_c (q_\infty - q_T) \quad (3.5.8.9)$$

where

q_∞ = filling factor with no cover and zero-thickness strip

q_T = correction for finite strip thickness

q_c = correction for effect of the shield

March recommends using q_∞ of [7] for $w/h \geq 5$

$$q_\infty = \left(1.0 + \frac{10.0 h}{w} \right)^j \quad (3.5.8.10)$$

$$j = a(w/h) b(\epsilon_r) \quad (3.5.8.11)$$

$$a(w/h) = 1.0 + \frac{1.0}{49.0} \ln \left\{ \frac{(w/h)^2 [(w/h)^2 + (1.0/52.0)^2]}{(w/h)^4 + 0.432} \right. \\ \left. + \frac{1.0}{18.7} \ln \left[1.0 + \left(\frac{w}{18.1 h} \right)^3 \right] \right\} \quad (3.5.8.12)$$

$$b(\epsilon_r) = -0.564 \left(\frac{\epsilon_r - 0.9}{\epsilon_r + 3.0} \right)^{0.053} \quad (3.5.8.13)$$

These equations have a stated accuracy of $\pm 0.2\%$ for

$$0.01 \leq w/h \leq 100$$

$$1 \leq \epsilon_r \leq 128$$

Stated accuracy is better than 0.2% agreement with Bryant and Weiss [3] for

$$1 \leq \epsilon_r \leq 30$$

$$0.05 \leq w/h \leq 20$$

The strip thickness correction is

$$qT = \frac{2.0 t \ln 2}{\pi h \sqrt{w/h}} \quad (3.5.8.14)$$

and the correction for the shield is ($h_2/h \geq 1$)

$$q_c = \tanh \left(1.043 + 0.121 h_2/h - \frac{1.164 h}{h_2} \right) \quad (3.5.8.15)$$

The equations were verified with MSTRIP calculations and the error was found to be less than 1% for the conditions above and $1 < h_2/h < \infty$. March also recommends the dispersion model of [5] for this structure based on exact calculations. To correct for frequency dependencies, we modify the calculated $\epsilon_{eff}(f=0)$

$$\epsilon_{eff}(f) = \epsilon_r - \frac{\epsilon_r - \epsilon_{eff}(0)}{1.0 + G(f/f_p)^2} \quad (3.5.8.16)$$

where

$$f_p = \frac{Z_0}{2.0 \mu_0 h} \quad (3.5.8.17)$$

$$G = \frac{\pi^2 (\epsilon_r - 1.0)}{12.0 \epsilon_{eff}(0)} \sqrt{\frac{2.0 \pi Z_0}{\eta_0}} \quad (3.5.8.18)$$

To calculate the frequency-dependent characteristic impedance, $Z_0(f)$, we convert the calculated zero-frequency impedance to an equivalent parallel-plate transmission line.

$$Z_0(f) = \frac{120.0 \pi h}{w_{eff}(f) \sqrt{\epsilon_{eff}(f)}} \quad (3.5.8.19)$$

$$w_{eff}(f) = w + \frac{w_{eff}(0) - w}{1.0 + (f/f_g)^2} \quad (3.5.8.20)$$

$$w_{eff}(0) = \frac{120.0 \pi h}{Z_0(f=0) \sqrt{\epsilon_{eff}(0)}} \quad (3.5.8.21)$$

where $Z_0(f=0)$ is the previously calculated zero-frequency value of the characteristic impedance, and the frequency where dispersion due to higher mode propagation begins is

$$f_g = \frac{c}{2.0 w_{eff}(f) \sqrt{\epsilon_{eff}(f) - 1.0}} \quad (3.5.8.22)$$

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3.6 STRIPLINE STRUCTURES

The term stripline was originally used to describe any planar transmission line based on strip conductors, including microstrip line. Its meaning in recent years has come to be more specific and now generally refers to a structure having a strip sandwiched between ground planes. Harold Wheeler developed a planar transmission line (two coplanar strips) which could be rolled up in 1936 and a stripline-like structure in 1942. The first use of "flat strip coaxial transmission line" (later to be known as stripline) was by V.H. Rumsey and H.W. Jamieson for an antenna power divider in WWII [2]. In 1949 Robert M. Barrett applied this new structure to printed-circuit boards launching what is now one of the most popular transmission line structures. Stripline was the name coined for a ground-signal-ground transmission line with an air dielectric at AIL around 1952. Tri-Plate was simultaneously developed at Sanders Associates, Inc. (stripline with solid dielectric) and a handbook of design techniques was published in 1956 [1].

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3.6.1 Zero-Thickness Centered Stripline

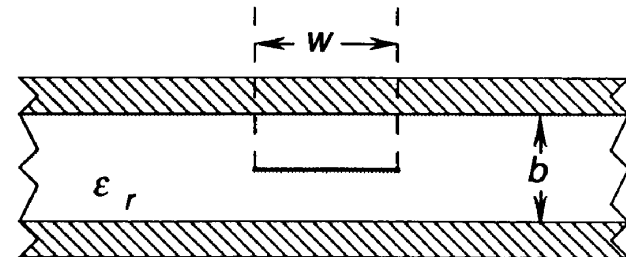


Figure 3.6.1.1: Zero-Thickness Stripline

Equations are exact for the zero-thickness case; however see 3.6.2 for a more realistic finite-thickness trace. The zero-thickness calculation is often used to check the limiting case ($t = 0$) of finite-thickness equations and as a term in finite-thickness calculations.

$$Z_0 = \frac{\eta_0 K(k)}{4.0 K(k')} \quad (3.6.1.1)$$

$$k = \operatorname{sech} \left(\frac{\pi w}{2.0 b} \right) \quad (3.6.1.2)$$

$$k' = \tanh \left(\frac{\pi w}{2.0 b} \right) \quad (3.6.1.3)$$

REFERENCE

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3.6.2 Centered Stripline, Triplate®² or Sandwich Line

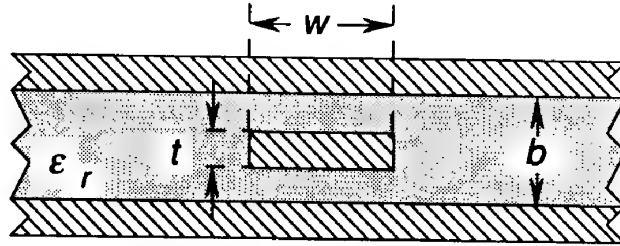


Figure 3.6.2.1: Centered Stripline

Note in Figure 3.6.2.1 and this equation, that the dielectric thickness is b ; some references refer to this as h or $2.0 h$. The equations are

$$Z_0 = \frac{\eta_0}{2.0 \pi \sqrt{\epsilon_r}} \ln \left\{ 1.0 + 0.5 \frac{8.0 b}{\pi w'} \left[\frac{8.0 b}{\pi w'} + \sqrt{\left(\frac{8.0 b}{\pi w'} \right)^2 + 6.27} \right] \right\} \quad (3.6.2.1)$$

where

$$2.0 h = b \quad (3.6.2.2)$$

$$w' = w + \frac{\Delta w}{t} t \quad (3.6.2.3)$$

$$\frac{\Delta w}{t} = \frac{1.0}{\pi} \ln \left[\frac{e}{\sqrt{\left(\frac{1}{2.0 b / t + 1} \right)^2 + \left(\frac{1/4\pi}{w/t + 1.1} \right)^m}} \right] \quad (3.6.2.4)$$

The above is sometimes written:

$$\frac{\Delta w}{t} = \frac{1.0}{\pi} \left\{ 1.0 - 0.5 \ln \left[\left(\frac{1.0}{2.0 b / t + 1.0} \right)^2 + \left(\frac{1.0 / 4\pi}{w / t + 1.1} \right)^m \right] \right\} \quad (3.6.2.5)$$

and

$$m = \frac{6.0}{3.0 + \frac{2.0 t}{b}} \quad (3.6.2.6)$$

The stated error is less than 1.5%. Nauwelaers and Van de Capelle [6] propose a thickness correction that is not dependent on the line width.

$$\frac{\Delta w}{t} = \frac{\ln \left(\frac{5.0 b}{t} \right)}{3.2} \quad (3.6.2.7)$$

The approximation is valid for

$$t / b < 0.1$$

$$w / t > 2.5$$

$$w / b > 0.1$$

Accuracy of the resulting Z_0 is claimed to be better than 0.5%.

The effects of sidewalls can be ignored when they are $>1.5 b$ away. If this condition is not met, use the equations of Section 3.6.4. The errors of these equations are $<0.5\%$ for

$$w' / b < 10.0.$$

²Triplate® is a registered trademark of Sanders Associates.

Losses are calculated with the incremental inductance rule. First the skin depth, δ is calculated:

$$\delta = \sqrt{\frac{2.0}{\omega \mu \sigma}} \quad (\text{m}) \quad (3.6.2.8)$$

where ω is the frequency in radians / second, μ is the permeability in H / m, and σ is the conductivity in mho / m. The power factor ($1.0 / Q$) is then found from the characteristic impedance of the structure when conductor dimensions are increased by δ :

$$p = 1.0 - \frac{Z_0(w + \delta, h + \delta, t + \delta)}{Z_0(w, b, t)} \quad (3.6.2.9)$$

The accuracy of p is better than ± 0.05 for $t/b < 0.5$, and ± 0.015 for $t/b < 0.25$. Recently, an exact analysis of conductor losses was presented in [8]. This solution is not closed-form and shows that the error in the above technique is less than 6%.

Dielectric losses are calculated with

$$\alpha_d = k_0 \frac{\sqrt{\epsilon_r} \tan \delta}{2.0} \quad (3.6.2.10)$$

where k_0 is the wavenumber in free-space.

The next higher mode will propagate in stripline when

$$\lambda_c = \sqrt{\epsilon_r} \left(\frac{2.0 w}{b} + \frac{4.0 d}{b} \right) b \quad (\text{units of } b) \quad (3.6.2.11)$$

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3.6.3 Off-Center Stripline

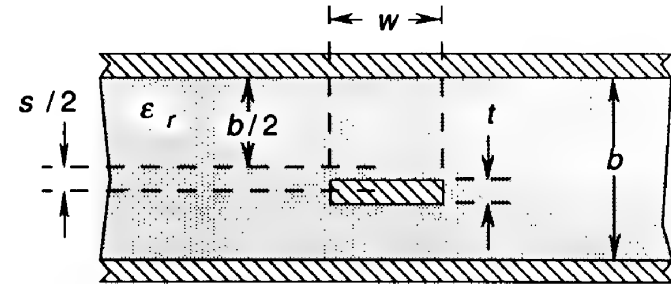


Figure 3.6.3.1: Off-Center Stripline

The off-center stripline structure is found mainly in multilayer boards. Two situations can make its use desirable. In order to make the board symmetric relative to the center-line (to avoid warpage) it is sometimes necessary to place the stripline off-center relative to its surrounding ground planes. It is also common to run two layers of striplines (*dual orthogonal stripline*) between a single pair of ground planes to increase density. To minimize coupling, the two trace layers are run perpendicular to each other. Each of these layers is off-center relative to the ground planes. See Kiang *et al.* [8] for more information on the analysis of dual orthogonal striplines.

Rao and Das [9] and Chang [1] analyze the off-center stripline using conformal mapping techniques. Chang's solution is limited to wide lines ($w \leq b$) and the solution of Rao and Das' equations is somewhat involved. Equations of Robrish [10] extend Cohn [2] to a wider range of w/b .

For $\frac{w}{b-t} < 0.35$ (narrow strips) these equations are

$$Z_0 = \frac{\eta_0}{2.0 \pi \sqrt{\epsilon_r}} \cosh^{-1} A \quad (3.6.3.1)$$

where

$$A = \sin\left(\frac{\pi c_l}{b}\right) \coth\left(\frac{\pi d_0}{2.0 b}\right) \quad (3.6.3.2)$$

$$c_l = \frac{b - s}{2.0} \quad (3.6.3.3)$$

The line width w is replaced with an equivalent round wire diameter, d_0 . A curve for this purpose is available in [3] and others³. A simple curve has been fit to the data for CAE use; the fit error is less than 0.35%:

$$\frac{d_0}{w} = 0.5008 + 1.0235 x - 1.0230 x^2 + 1.1564 x^3 - 0.4749 x^4 \quad (3.6.3.4)$$

$$x = \frac{t}{w} \quad (3.6.3.5)$$

where $w \geq t$ (reverse their positions if this condition is not met).

For $\frac{w}{b-t} > 0.35$ (wide strips)

$$Z_0 = \frac{\eta_0}{\sqrt{\epsilon_r}} \frac{1.0}{\frac{w/b}{\gamma} + \frac{w/b}{\beta - \gamma} + \frac{2.0 C_f}{\epsilon_0 \epsilon_r}} \quad (\Omega) \quad (3.6.3.6)$$

where

$$C_f = \frac{\epsilon_r \epsilon_0}{\pi} \left\{ 2.0 \ln \left[\frac{1.0}{\gamma(\beta - \gamma)} \right] + \frac{1.0}{\gamma(\beta - \gamma)} [F(t/b) - F(c_l/b)] \right\} \quad (3.6.3.7)$$

$$F(x) = (1.0 - 2.0 x) [(1.0 - x) \ln(1.0 - x) - x \ln x] \quad (3.6.3.8)$$

$$\gamma = \frac{c_l}{b} - \frac{t}{2.0 b} \quad (3.6.3.9)$$

$$\beta = 1.0 - \frac{t}{b} \quad (3.6.3.10)$$

for $\frac{w}{b-t} < \frac{t}{b}$ replace w/b with

$$\left(\frac{w}{b}\right)_{eff} = \frac{w}{b} + \left(1.0 - \frac{t}{b}\right)^{8.0} \left[\frac{K(k')}{K(k)} - \frac{2.0}{\pi} \ln 2.0 - \frac{w}{b} \right] \quad (3.6.3.11)$$

$$k = \operatorname{sech} \left(\frac{\pi w}{2.0 b} \right) \quad (3.6.3.12)$$

$$k' = \tanh \left(\frac{\pi w}{2.0 b} \right) \quad (3.6.3.13)$$

Agreement with exact analytic solutions is better than 0.5%. The accuracy of these relations was compared to numerical solutions and found to be better than $\pm 2\%$ for

$$0.2 < \frac{c_l}{b} < 0.8$$

$$w > t$$

$$t/b < 0.2$$

Two other approaches to this problem are also found in the literature. Howe uses a parallel plate and fringing capacitance calculation while Hogan calculates a correction based on equivalent widths.

These equations from Howe [6] have been rearranged to show the underlying parallel plate capacitance calculation.

$$Z_0 = \frac{\eta_0}{\sqrt{\epsilon_r} \frac{C'}{\epsilon}} \quad (3.6.3.14)$$

where

$$\frac{C'}{\epsilon} = \frac{C_{p1}}{\epsilon} + \frac{C_{p2}}{\epsilon} + \frac{2.0 C_{f1}}{\epsilon} + \frac{2.0 C_{f2}}{\epsilon} \quad (3.6.3.15)$$

$$\frac{C_{p1}}{\epsilon} = \frac{w}{\frac{b}{2} - \frac{s}{2} - \frac{t}{2}} = \frac{2.0 \left(\frac{w}{b}\right)}{1.0 - \frac{s}{b} - \frac{t}{b}} \quad (3.6.3.16)$$

³The source of this equivalence relation is C. Flammer, "Equivalent Radii of Thin Cylindrical Antennas with Arbitrary Cross Sections," *Stanford Research Institute Technical Report*, March 15, 1950.

$$\frac{C_{p2}}{\epsilon} = \frac{w}{\frac{b}{2} + \frac{s}{2} - \frac{t}{2}} = \frac{2 \left(\frac{w}{b} \right)}{1.0 + \frac{s}{b} - \frac{t}{b}} \quad (3.6.3.17)$$

$$\begin{aligned} \frac{C_{f1}}{\epsilon} = \frac{1.0}{\pi} & \left(\frac{2.0}{1.0 - t/(b-s)} \ln \left[\frac{1.0}{1.0 - t/(b-s)} + 1.0 \right] \right. \\ & \left. - \left[\frac{1.0}{1.0 - t/(b-s)} - 1.0 \right] \ln \left\{ \frac{1.0}{[1.0 - t/(b-s)]^2} - 1.0 \right\} \right) \end{aligned} \quad (3.6.3.18)$$

$$\begin{aligned} \frac{C_{f2}}{\epsilon} = \frac{1.0}{\pi} & \left(\frac{2.0}{1.0 - t/(b+s)} \ln \left[\frac{1.0}{1.0 - t/(b+s)} + 1.0 \right] \right. \\ & \left. - \left[\frac{1.0}{1.0 - t/(b+s)} - 1.0 \right] \ln \left\{ \frac{1.0}{[1.0 - t/(b+s)]^2} - 1.0 \right\} \right) \end{aligned} \quad (3.6.3.19)$$

For $\frac{w}{b-t} < 0.35$, substitute

$$\frac{w_n}{b} = \frac{0.07 (1 - t/b) + w/b}{1.2} \quad (3.6.3.20)$$

for $\frac{w}{b}$ in the equations for C_{p1}/ϵ and C_{p2}/ϵ .

The term C' is the incremental capacitance of the line and C_{p1} and C_{p2} represent the parallel-plate capacitances of each side and C_{f1} and C_{f2} represent the fringing field capacitance of the edges. A limitation of these equations is that the results of this technique differ substantially from centered stripline equations when the strip is centered. Hogan's [7] equations offer a solution which remedies this problem.

Hogan's solution modifies Wheeler's centered stripline equations and suggests additional corrections for the trapezoidal cross-section caused by the etching process of PCBs. A different figure is required because of substantially different dimensioning.

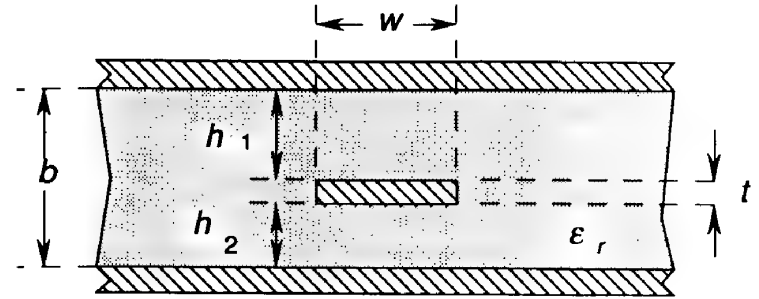


Figure 3.6.3.2: Alternate Off-Center Stripline Labeling

The equations are:

$$Z_0 = \frac{1.0}{\sqrt{\epsilon_r}} [Z_{0,ctr}(1.0, w, b, t) - \Delta Z_{0,air}] \quad (3.6.3.21)$$

where

$Z_{0,ctr}(\epsilon_r, w, b, t)$ = the characteristic impedance of a centered stripline (3.6.2.1)

$$Z_{0,air} = 2.0 \left[\frac{Z_{0,ctr}(1.0, w, h_1, t) Z_{0,ctr}(1.0, w, h_2, t)}{Z_{0,ctr}(1.0, w, h_1, t) + Z_{0,ctr}(1.0, w, h_2, t)} \right] \quad (3.6.3.22)$$

$$\Delta Z_{0,air} = \frac{0.26 \pi}{8.0} Z_{0,air}^2 \left(0.5 - \frac{h_1 + t/2}{h_1 + h_2 + t} \right)^{2.2} \left(\frac{t + w}{h_1 + h_2 + t} \right)^{2.9} \quad (3.6.3.23)$$

For trapezoidal cross-section traces with the widest dimension closest to a ground plane (Figure 3.6.3.3), replace equations (3.6.3.20) and (3.6.3.21) with:

$$Z_{0,air} = 2.0 \left[\frac{Z_{0,ctr}(1.0, w_{top}, h_1, t) Z_{0,ctr}(1.0, w_{bottom}, h_2, t)}{Z_{0,ctr}(1.0, w_{top}, h_1, t) + Z_{0,ctr}(1.0, w_{bottom}, h_2, t)} \right] \quad (3.6.3.24)$$

$$\Delta Z_{0,air} = \frac{0.26 \pi}{8.0} Z_{0,air}^2 \left(0.5 - \frac{h_1 + t/2}{h_1 + h_2 + t} \right)^{2.2} \left(\frac{t + w_{eff}}{h_1 + h_2 + t} \right)^{2.9}$$

where

$$w_{eff} = \frac{w_{top} + w_{bottom}}{2.0} \quad (3.6.3.25)$$

(3.6.3.26)

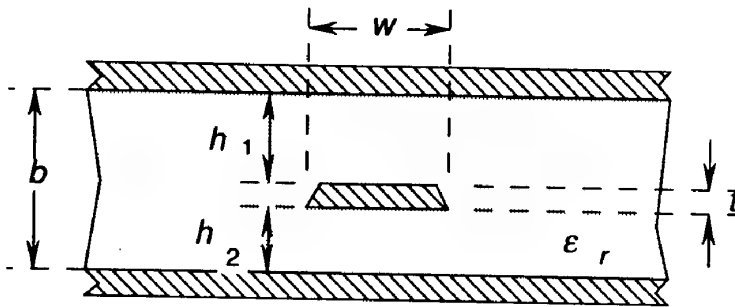


Figure 3.6.3.3: Trapezoidal Off-Center Stripline

For trapezoidal cross-section traces with the narrowest dimension closest to a ground plane (Figure 3.6.3.4) replace w in the equations above with w_{eff} giving:

$$Z_{0,air} = 2.0 \left[\frac{Z_{0,ctr}(1.0, w_{eff}, h_1, t) Z_{0,ctr}(1.0, w_{eff}, h_2, t)}{Z_{0,ctr}(1.0, w_{eff}, h_1, t) + Z_{0,ctr}(1.0, w_{eff}, h_2, t)} \right] \quad (3.6.3.27)$$

$$\Delta Z_{0,air} = \frac{0.26 \pi}{8.0} Z_{0,air}^2 \left(0.5 - \frac{h_1 + t/2}{h_1 + h_2 + t} \right)^{2.2} \left(\frac{t + w_{eff}}{h_1 + h_2 + t} \right)^{2.9} \quad (3.6.3.28)$$

where

$$w_{eff} = \frac{\frac{w_{top} + w_{bottom}}{2.0} + \frac{2.0 w_{top} w_{bottom}}{w_{top} + w_{bottom}}}{2.0} \quad (3.6.3.29)$$

or

$$w_{eff} = \frac{(w_{top} + w_{bottom})^2 + 4.0 w_{top} w_{bottom}}{4.0 (w_{top} + w_{bottom})} \quad (3.6.3.30)$$

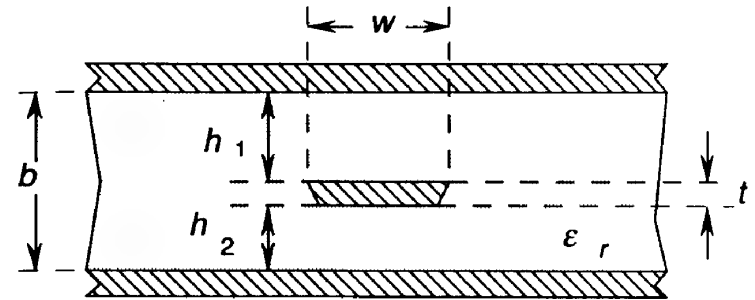


Figure 3.6.3.4: Trapezoidal Off-Center Stripline

Note that these equations are symmetric and the choice of w_{top} or w_{bottom} is relative to h_1 and h_2 and therefore arbitrary.

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3.6.4 Shielded Stripline

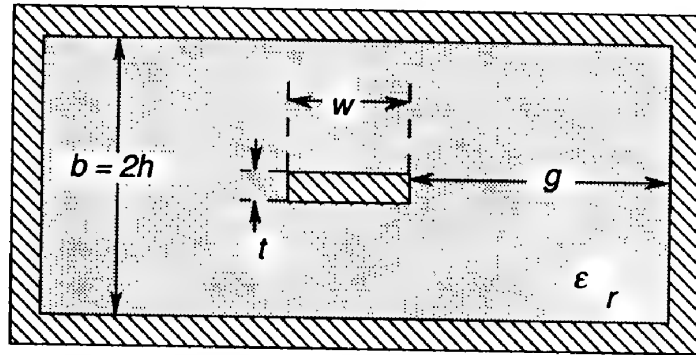


Figure 3.6.4.1: Shielded Stripline

Generally, the stripline configuration is considered to have ground planes extending to infinity. This configuration corrects for the effect of side walls at a finite distance. The equations of [1] are:

$$Z_0 = \frac{30.0 \pi}{\sqrt{\epsilon_r} \left[\frac{w/b}{1.0 - t/b} + \frac{2.0 C_f(t/b)}{\pi C_f(0)} \ln \left(1.0 + \coth \frac{\pi g}{b} \right) \right]} \quad (\Omega) \quad (3.6.4.1)$$

$$C_f(t/b) = \frac{\epsilon}{\pi} \left[\frac{b}{(b-t)} \ln \frac{(2.0 b - t)}{t} + \ln \frac{t(2.0 b - t)}{(b-t)^2} \right] \quad (3.6.4.2)$$

$$C_f(0) = \frac{2.0 \epsilon \ln 2.0}{\pi} \quad (3.6.4.3)$$

Souza [2] describes a conformal mapping procedure that yields the equations for a shielded stripline that is off center.

REFERENCES

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3.7 INVERTED MICROSTRIP LINE STRUCTURES

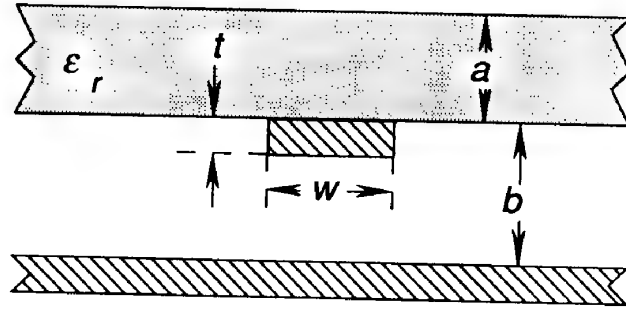


Figure 3.7.1.1: Inverted Microstrip Line

Inverted microstrip is used for precision low-loss lines. It allows a wider range of achievable impedances. In addition, the air dielectric is essentially lossless and dominates the dispersion characteristics.

The equations of [1] are valid for $\epsilon_r \leq 3.8$. For dielectrics with $\epsilon_r \leq 20$, use the more involved equations of [2]. These closed-form equations are the results of curve-fitting data generated by an exact variational analysis.

$$Z_0 = \frac{\eta_0}{2.0 \pi \sqrt{\epsilon_{eff}}} \ln \left[\frac{f(u)}{u} + \sqrt{1.0 + \frac{4.0}{u^2}} \right] (\Omega) \quad (3.7.1.1)$$

$$\epsilon_{eff} = (1.0 + f_1 f_2)^2 \quad (3.7.1.2)$$

where

$$u = \frac{w}{b} \quad (3.7.1.3)$$

$$f(u) = 6.0 + (2\pi - 6.0) e^{-(30.666/u)^{0.7528}} \quad (3.7.1.4)$$

$$f = \ln(\epsilon_r) \quad (3.7.1.5)$$

$$f_1 = \sqrt{\epsilon_r} - 1 \quad (3.7.1.6)$$

$$f_2 = \frac{1}{\sum_{j=0}^3 c_j \left(\frac{w}{b}\right)^j} \quad (3.7.1.7)$$

$$c_i = \frac{1}{\sum_{j=0}^3 d_{ij} \left(\frac{b}{a}\right)^j} \quad (3.7.1.8)$$

$$d_{00} = (2359.401 - 97.1644f - 5.706 \times f^2 + 11.4112 \times f^3) \times 10^{-3} \quad (3.7.1.9)$$

$$d_{01} = (4855.9472 - 3408.5207f + 15296.73f^2 - 2418.1785f^3) \times 10^{-5} \quad (3.7.1.10)$$

$$d_{02} = (1763.34 + 961.0481f - 2089.28f^2 + 375.8805f^3) \times 10^{-5} \quad (3.7.1.11)$$

$$d_{03} = (-556.0909 - 268.6165f + 623.7094 \times f^2 + -119.1402 \times f^3) \times 10^{-6} \quad (3.7.1.12)$$

$$d_{10} = (219.066 - 253.0864f + 208.7469f^2 - 27.3285f^3) \times 10^{-3} \quad (3.7.1.13)$$

$$d_{11} = (915.5589 + 338.4033f - 253.2933f^2 + 40.4745f^3) \times 10^{-3} \quad (3.7.1.14)$$

$$d_{12} = (-1957.379 - 1170.936f + 1480.857f^2 - 347.6403f^3) \times 10^{-5} \quad (3.7.1.15)$$

$$d_{13} = (486.7425 + 279.8323f - 431.3625f^2 + 108.824f^3) \times 10^{-6} \quad (3.7.1.16)$$

$$d_{20} = (5602.767 + 4403.356f - 4517.034f^2 + 743.2717f^3) \times 10^{-5} \quad (3.7.1.17)$$

$$d_{21} = (-2823.481 - 1562.782f + 3646.15f^2 - 823.4223f^3) \times 10^{-5} \quad (3.7.1.18)$$

$$d_{22} = (255.893 + 158.5529f - 3235.485f^2 + 919.3661f^3) \times 10^{-6} \quad (3.7.1.19)$$

$$d_{23} = (-147.0235 + 62.4342f + 887.5211f^2 - 270.7555f^3) \times 10^{-7} \quad (3.7.1.20)$$

$$d_{30} = (-3170.21 - 1931.852f + 2715.327f^2 - 519.342f^3) \times 10^{-6} \quad (3.7.1.21)$$

$$d_{31} = (596.3251 + 188.1409f - 1741.477f^2 + 465.6756f^3) \times 10^{-6} \quad (3.7.1.22)$$

$$d_{32} = (124.9655 + 577.5381 f + 1366.453 f^2 - 481.13 f^3) \times 10^{-7} \quad (3.7.1.23)$$

$$d_{33} = (-530.2099 - 2666.352 f - 3220.096 f^2 + 1324.499 f^3) \times 10^{-9} \quad (3.7.1.24)$$

The stated error of the fit to the exact theoretical calculation is $< 0.6\%$ (Z_0 and $\sqrt{\epsilon_{eff}}$) for:

$$1.0 \leq \epsilon_r \leq 20.0$$

$$0.5 \leq w/b \leq 10.0$$

$$0.06 \leq a/b \leq 1.5$$

REFERENCES

- [1] Tomar, R.S., and Prakash Bhartia, "Suspended and Inverted Microstrip Design," *Microwave Journal*, March 1986, pp. 173–177.
- [2] Tomar, R.S., and Prakash Bhartia, "New Quasi-Static Models for the Computer-Aided Design of Suspended and Inverted Microstrip Lines," *IEEE Transactions on Microwave Theory and Techniques*, Vol. MTT-35, No. 4, April 1987, pp. 453–457, and corrections No. 11, November 1987, p. 1076.

3.8 SUSPENDED MICROSTRIP LINE STRUCTURES

3.8.1 Suspended Microstrip Line

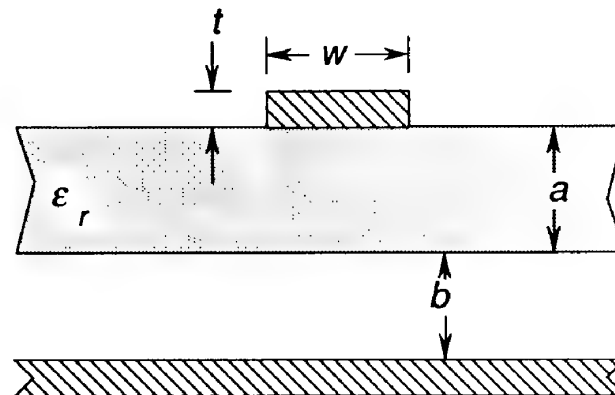


Figure 3.8.1.1: Suspended Microstrip Line

Suspended microstrip is used for precision low-loss lines. This configuration also allows a wider range of realizable Z 's and reduced dispersion. Again the simple equations of [4] are valid for $\epsilon_r \leq 3.8$, otherwise use [6].

$$Z_0 = \frac{\eta_0}{2.0 \pi \sqrt{\epsilon_{eff}}} \ln \left[\frac{f(u)}{u} + \sqrt{1.0 + \frac{4.0}{u^2}} \right] \quad (\Omega) \quad (3.8.1.1)$$

$$\epsilon_{eff} = \frac{1.0}{(1.0 - f_1 f_2)^2} \quad (\text{unitless}) \quad (3.8.1.2)$$

where

$$u = \frac{w/b}{1.0 + a/b} \quad (3.8.1.3)$$

$$f(u) = 6.0 + (2\pi - 6.0) e^{-(30.666/u)^{0.7528}} \quad (3.8.1.4)$$

$$f = \ln(\epsilon_r) \quad (3.8.1.5)$$

$$f_1 = 1.0 - \frac{1.0}{\sqrt{\epsilon_r}} \quad (3.8.1.6)$$

$$f_2 = \frac{1}{\sum_{i=0}^3 c_i \left(\frac{w}{b}\right)^i} \quad (3.8.1.7)$$

$$c_i = \frac{1}{\sum_{j=0}^3 d_{ij} \left(\frac{b}{a}\right)^j} \quad (3.8.1.8)$$

$$d_{00} = (176.2576 - 43.1240f + 13.4094f^2 - 1.7010f^3) \times 10^{-2} \quad (3.8.1.9)$$

$$d_{01} = (4665.2320 - 1790.4000f + 291.5858f^2 - 8.0888f^3) \times 10^{-4} \quad (3.8.1.10)$$

$$d_{02} = (-3025.5070 - 141.9368f - 3099.4700f^2 + 777.6151f^3) \times 10^{-6} \quad (3.8.1.11)$$

$$d_{03} = (2481.5690 + 1430.3860f + 10095.5500f^2 - 2599.1320f^3) \times 10^{-8} \quad (3.8.1.12)$$

$$d_{10} = (-1410.2050 + 149.9293f + 198.2892f^2 - 32.1679f^3) \times 10^{-4} \quad (3.8.1.13)$$

$$d_{11} = (2548.7910 + 1531.9310f - 1027.500f^2 + 138.4192f^3) \times 10^{-4} \quad (3.8.1.14)$$

$$d_{12} = (999.3135 - 4036.7910f + 1762.4120f^2 - 298.0241f^3) \times 10^{-6} \quad (3.8.1.15)$$

$$d_{13} = (-1983.7890 + 8523.9290f - 5235.4600f^2 + 1145.7880f^3) \times 10^{-8} \quad (3.8.1.16)$$

$$d_{20} = (1954.0720 + 333.3873f - 700.7473f^2 + 121.3212f^3) \times 10^{-5} \quad (3.8.1.17)$$

$$d_{21} = (-3931.0900 - 1890.7190f + 1912.2660f^2 - 319.6794f^3) \times 10^{-5} \quad (3.8.1.18)$$

$$d_{22} = (-532.1326 + 7274.7210f - 4955.7380f^2 + 941.4134f^3) \times 10^{-7} \quad (3.8.1.19)$$

$$d_{23} = (138.2037 - 1412.4270f + 1184.2700f^2 - 270.0047f^3) \times 10^{-8} \quad (3.8.1.20)$$

$$d_{30} = (-983.4028 - 255.1229f + 455.8729f^2 - 83.9468f^3) \times 10^{-6} \quad (3.8.1.21)$$

$$d_{31} = (1956.3170 + 779.9975f - 995.9494f^2 + 183.1957f^3) \times 10^{-6} \quad (3.8.1.22)$$

$$d_{32} = (62.8550 - 3462.5000f + 2909.9230f^2 - 614.7068f^3) \times 10^{-8} \quad (3.8.1.23)$$

$$d_{33} = (-35.2531 + 601.0291f - 643.0814f^2 + 161.2689f^3) \times 10^{-9} \quad (3.8.1.24)$$

The stated error of the fit to the exact theoretical calculation is less than 0.6% (Z_0 and $\sqrt{\epsilon_{eff}}$) for:

$$1.0 \leq \epsilon_r \leq 20.0$$

$$0.5 \leq w/b \leq 10.0$$

$$0.06 \leq a/b \leq 1.5$$

An analysis of conductor and dielectric loss is given in Yamashita [7].

Tomar and Bhartia [5] give polynomials for the dispersion of this structure. To use these, replace ϵ_{eff} in the equations above with $\epsilon_{eff}(f)$ calculated below. The variable $\epsilon_{eff}(0)$ is the static value ϵ_{eff} calculated above.

For $\epsilon_r = 12.9$ the equations are

$$\epsilon_r(f) = \epsilon_r - \frac{\epsilon_r - \epsilon_{eff}(0)}{1.0 + G \left(\frac{f}{f_p}\right)^2} \quad (3.8.1.25)$$

where

$$f_p = \frac{Z}{2.0 \pi \mu_0 (a + b)} \frac{a}{0.064} \quad (3.8.1.26)$$

$$a, b = (\text{cm})$$

$$G = C_0 + C_1 Z + C_2 Z^2 \quad (3.8.1.27)$$

and

$$C_0 = 0.0194 - 0.2398 \left(\frac{a}{b}\right) + 0.8977 \left(\frac{a}{b}\right)^2 - 0.9924 \left(\frac{a}{b}\right)^3 + 0.3468 \left(\frac{a}{b}\right)^4 \quad (3.8.1.28)$$

$$C_1 = -0.0008 + 0.0096 \left(\frac{a}{b}\right) - 0.0346 \left(\frac{a}{b}\right)^2 + 0.0384 \left(\frac{a}{b}\right)^3 - 0.0135 \left(\frac{a}{b}\right)^4 \quad (3.8.1.29)$$

$$C_2 = 0.0004 \left(\frac{a}{b}\right)^2 - 0.0004 \left(\frac{a}{b}\right)^3 + 0.0001 \left(\frac{a}{b}\right)^4 \quad (3.8.1.30)$$

The dispersion relations are valid for

$$1.0 \leq w/b \leq 10.0$$

$$20 \text{ GHz} < f < 100 \text{ GHz}$$

The accuracy is within $\pm 2\%$ when compared with exact theoretical calculations.

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- [1] Bedair, S.S., "Closed Form Expressions for the Computer-Aided Design of Suspended Microstrip Circuits for Millimeter Wave Applications," *IEEE Proceedings*, Vol. 134, Pt. H, No. 4, August 1987, pp. 386–388.
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3.8.2 Shielded Suspended Substrate Stripline

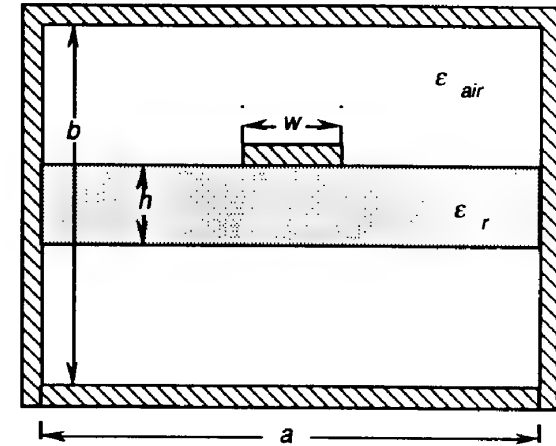


Figure 3.8.2.1: Shielded Suspended Substrate Stripline

The equations for this structure are broken into two ranges.
For narrow strips ($0 < w < a/2.0$):

$$Z_0 = \frac{\eta_0}{2.0 \pi \sqrt{\epsilon_{eff}}} \left\{ V + R \ln \left[\frac{6.0}{w/b} + \sqrt{1.0 + \frac{4.0}{(w/b)^2}} \right] \right\} (\Omega) \quad (3.8.2.1)$$

$$\epsilon_{eff} = \frac{1.0}{\left[1.0 + \left(E - F \ln \frac{w}{b} \right) \ln \left(\frac{1.0}{\sqrt{\epsilon_r}} \right) \right]^2} \quad (3.8.2.2)$$

where

$$V = 1.7866 - 0.2035 \frac{h}{b} + 0.4750 \frac{a}{b} \quad (3.8.2.3)$$

$$R = 1.0835 + 0.1007 \frac{h}{b} - 0.09457 \frac{a}{b} \quad (3.8.2.4)$$

$$E = 0.2077 + 1.2177 \frac{h}{b} - 0.08364 \frac{a}{b} \quad (3.8.2.5)$$

$$F = 0.03451 - 0.1031 \frac{h}{b} + 0.01742 \frac{a}{b} \quad (3.8.2.6)$$

and for wide strips ($a/2 < w < a$)

$$Z_0 = \frac{\eta_0}{\sqrt{\epsilon_{eff}}} \left[V + \frac{R}{\frac{w}{b} + 1.3930 + 0.6670 \ln \left(\frac{w}{b} + 1.444 \right)} \right] (\Omega) \quad (3.8.2.7)$$

$$\epsilon_{eff} = \frac{1.0}{\left[1.0 + \left(E - F \ln \frac{w}{b} \right) \ln \left(\frac{1.0}{\sqrt{\epsilon_r}} \right) \right]^2} \quad (3.8.2.8)$$

where

$$V = -0.6301 - 0.07082 \frac{h}{b} + 0.2470 \frac{a}{b} \quad (3.8.2.9)$$

$$R = 1.9492 + 0.1553 \frac{h}{b} - 0.5123 \frac{a}{b} \quad (3.8.2.10)$$

$$E = 0.4640 + 0.9647 \frac{h}{b} - 0.2063 \frac{a}{b} \quad (3.8.2.11)$$

$$F = -0.1424 + 0.3017 \frac{h}{b} - 0.02411 \frac{a}{b} \quad (3.8.2.12)$$

The above relations are valid for:

$$1.0 \leq a/b \leq 2.5$$

$$1.0 \leq \epsilon_r \leq 4.0$$

$$0.1 < h/b < 0.5$$

The wide strip equations agree with finite-differential techniques to $\pm 3\%$. Narrow strip agreement is $\pm 2\%$.

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[3] Yamashita, Eikichi, and Kuzuhiko Atsuki, "Strip Line with Rectangular Outer Conductor and Three Dielectric Layers," *IEEE Transactions on Microwave Theory and Techniques*, Vol. MTT-18, No. 5, May 1970, p. 238-244.

3.9 OTHER STRUCTURES

3.9.1 Cylindrical Conductor in a Corner

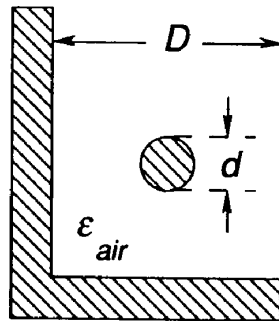


Figure 3.9.1.1: Cylindrical Conductor in a Corner

In [1], Wheeler gives solutions for a number of "wire-in-a-shield" problems. For this configuration:

$$Z_0 = \frac{\eta_0}{4.0 \pi} \ln \left\{ \left(\frac{D}{d} \right)^2 + \sqrt{\left[\left(\frac{D}{d} \right)^2 - 1.0 \right]^2 + \left(\frac{D}{d} \right)^2 - 1.0} \right\} \quad (3.9.1)$$

The stated accuracy is about 1%.

REFERENCES

- [1] Wheeler, Harold A., "Transmission-Line Properties of a Round Wire in a Polygon Shield," *IEEE Transactions on Microwave Theory and Techniques*, Vol. MTT-27, No. 8, August 1979, pp. 717-721.

3.9.2 Slabline

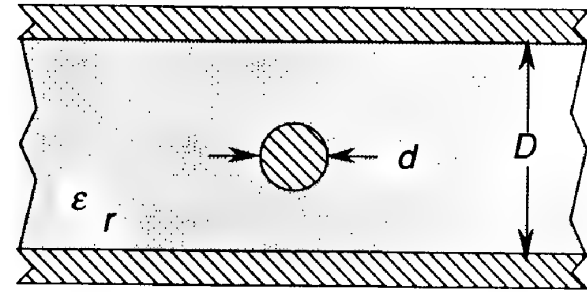


Figure 3.9.2.1: Slabline

The slabline is a transmission line where a thin wire replaces a printed trace, e.g., for an inductor. A simple, frequently encountered relation is

$$Z_0 = \frac{\eta_0}{2.0 \pi \sqrt{\epsilon_r}} \ln \left(\frac{4.0 D}{\pi d} \right) \quad (\Omega) \quad (3.9.2.1)$$

The stated error is less than 1.0% for $\frac{D}{d} \geq 2.0$ with better accuracy as Z_0 increases ($d/D \rightarrow 0$).

White [5, p. 520] has a more accurate equation to be used for synthesis. For d approaching D , Wheeler's [4] equations are more accurate than the above:

$$Z_0 = \frac{15.0}{\sqrt{\epsilon_r}} \ln \left[1.0 + \frac{k_0 g}{2.0} + \sqrt{\left(\frac{k_0 g}{2} \right)^2 + 2.0 g} \right] \quad (\Omega) \quad (3.9.2.2)$$

where

$$w/h = \frac{2.0}{\left[1.0 + \frac{(e^{8.0 \pi r} - 1.0)^2}{k_0(e^{8.0 \pi r} - 1.0) + 2.0} \right]^{0.25}} - 1.0 \quad (3.9.2.3)$$

$$k_0 = \left(\frac{4}{\pi} \right)^4 \quad (3.9.2.4)$$

$$g = (1.0 + 2.0 h / 2.0)^4 - 1.0 \quad (3.9.2.5)$$

Stated accuracy is better than 0.5% for Z_0 .

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3.9.3 Wire over Dielectric over Ground Plane

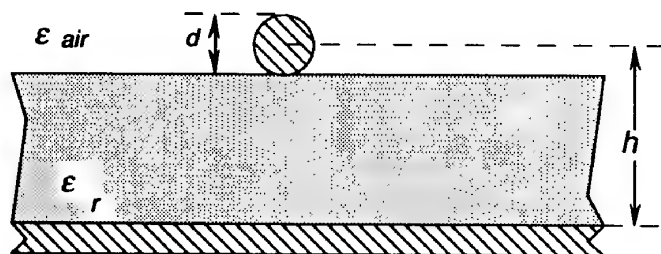


Figure 3.9.3.1: Wire over Dielectric over Ground Plane

This configuration is useful when a line is too fine for standard PCB etching. See Chapter 11 for available wire gauges.

$$Z_0 = \frac{\eta_0}{2.0 \pi \sqrt{\epsilon_{eff}}} \cosh^{-1} \left(\frac{2.0 h}{d} \right) \approx \frac{\eta_0}{2.0 \pi \sqrt{\epsilon_{eff}}} \ln \left(\frac{4.0 h}{d} \right) \quad (3.9.3.1)$$

The approximate equation above is also commonly encountered and is valid for $h/d \gg d$. The accuracy of Z_0 is stated to be better than 0.024% in [2] and 1% in [3].

To calculate ϵ_{eff} , use the equation for microstrip line (3.5.1.2) replacing the microstrip line width, w , with an equivalent width calculated from the wire diameter, d . A curve for this purpose is available in [1] and others. We have fit a simple curve to the data for CAE use; the fit error is less than 0.35%:

$$\frac{d_0}{w} = 0.5008 + 1.0235 x - 1.0230 x^2 + 1.1564 x^3 - 0.4749 x^4 \quad (3.9.3.2)$$

$$x = \frac{t}{w} \quad (3.9.3.3)$$

where $w \geq t$ (reverse their positions if this condition is not met).

REFERENCES

- [1] Cohn, Seymour B., "Characteristic Impedance of the Shielded-Strip Transmission Line," *IRE Transactions on Microwave Theory and Techniques*, Vol. MTT-2, No. 3, July 1954, pp. 52-57.
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3.9.4 Bond Wire

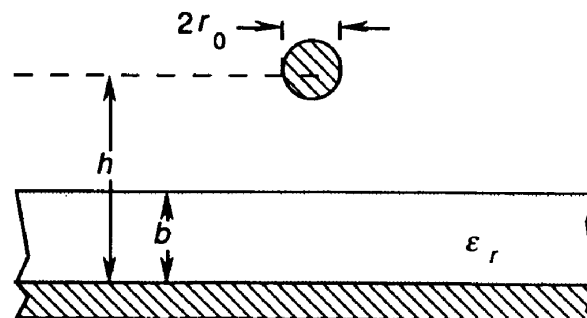


Figure 3.9.4.1: Wire above Grounded Dielectric

The wire over grounded dielectric is encountered where we have a bond wire connecting a circuit substrate to the package leads, the case, or another substrate. Caverly [1], uses conformal mapping to derive

$$Z_0 = \frac{\eta_0}{2.0 \pi \sqrt{\epsilon_{eff}}} \cosh^{-1} \left(\frac{1.0 - u_1^2}{2.0 R_1} + \frac{R_1}{2.0} \right) \quad (\Omega) \quad (3.9.4.1)$$

where

$$R_1 = \frac{2.0}{(4.0 h / r_0 - r_0 / h)} \quad (3.9.4.2)$$

$$u_1 = \frac{1.0}{(2.0 h / r_0)^2 - 1.0} \quad (3.9.4.3)$$

Caverly calculates ϵ_{eff} using an equivalent resistive network and SPICE. For calculations we will use the closed-form ϵ_{eff} of Kuester and Chang [2]:

$$\epsilon_{eff} = \frac{\ln \left(\frac{2.0 h}{r_0} \right)}{\ln \left[\frac{2.0 (h - b)}{r_0} + \frac{2.0 b}{r_0 \epsilon_r} \right]} \quad (\text{unitless}) \quad (3.9.4.4)$$

The equation for ϵ_{eff} is valid for quasi-TEM propagation or

$$\frac{\epsilon_r (h - b)}{\lambda_0} \ln \left[\frac{2.0 \lambda_0 (h - b)}{r_0} \right] < 1.0 \quad (3.9.4.5)$$

where

λ_0 = free-space wavelength in units of h, b, r_0

Mondal [3] derives equations for a bond wire arc as shown in Figure 3.9.4.2.

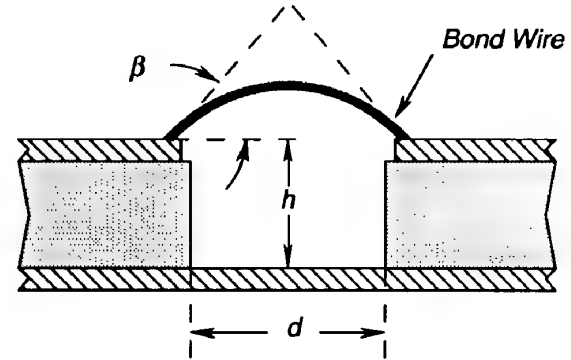


Figure 3.9.4.2: Bond Wire

The equations for this structure are

$$L = 10.0 \int_{0.0}^{d/2.0} \ln [p(x)] dx \quad (\text{nH}) \quad (3.9.4.6)$$

where

d, h, r are all in inches

$$p(x) = \frac{2.0}{r} \left[\sqrt{\left(\frac{d \operatorname{cosec} \beta}{2.0} \right)^{2.0} - x^2} + \left(h - \frac{d}{2.0} \cot \beta \right) \right] \quad (3.9.4.7)$$

And the capacitance is calculated with

$$C = \frac{1.4337}{\ln p(x)} \quad (\text{pF / in}) \quad (3.9.4.8)$$

where

l = the length of the bond wire

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3.9.5 Multiwire®⁴

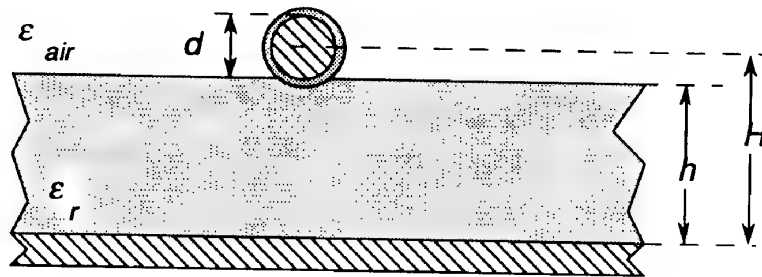


Figure 3.9.5.1: Multiwire®

Multiwire® is a process that routes wires over an adhesive layer on a ground plane using a numerically controlled machine. The grid used for routing is usually 0.050 in. It can provide very quick prototypes. Multiwire allows line lengths to be monitored and controlled and easily rerouted to avoid coupling. Lines may be criss-crossed to form a pseudo-twisted-pair structure. After the wires are routed, the board processing continues similar to a standard PCB, including plated through holes.

The dielectric is a mixture of G-10 and adhesives that have a combined dielectric constant of about 5.0. Available wire sizes are 34, 38, and 41 awg.

Shibata and Terakado [4] examine the effect of different adhesive layer dielectric constants ($\epsilon_r = 1.0, 2.65, \text{ and } 5.0$) and the positioning of the wire in the adhesive on Z_0 (Figure 6a), v_p (Figure 6b), and ϵ_{eff} (Figure 7). The reference analyzes three cases: (1) wire riding on top of dielectric, (2) wire halfway embedded in dielectric, and (3) wire fully embedded in dielectric. An approximation is given [4]:

$$Z_0 = 17.08 + 34.83 \ln \left(\frac{2.0 h}{d} \right) (\Omega) \quad (3.9.5.1)$$

Accuracy is stated to be better than 0.35% for

$$\epsilon_r = 5.0$$

$$0.01 \leq \frac{d}{2.0 h} \leq 0.3$$

A controlled impedance environment may also be simulated by creating a pseudo-ground plane from a grid of grounded wires beneath the conductors. Burr [2] gives the following empirical relations for the characteristic impedance:

$$Z_0 = \frac{41.2}{\epsilon_{eff}} \ln (1.88 H) (\Omega) \quad (3.9.5.2)$$

$$\epsilon_{eff} = 0.475 \epsilon_r + 0.67 \quad (3.9.5.3)$$

Sensitivity analysis shows that for 50 Ω lines the adhesive thickness needs to be monitored in order to achieve $\pm 10\%$. Higher impedances do not require tight tolerancing of h , and at 90 Ω the line dimensions are not critical.

The dielectric constant of the wire's insulation is approximately 3.5. Bends are of constant diameter and do not require compensation. Where wires cross, they are flattened which introduces a capacitance between the lines of approximately 0.026 pF (parallel plate equation).

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⁴Multiwire® is a registered trademark of the Multiwire Division of the Kollmorgen Corporation. Similar products are produced by other manufacturers.

3.9.6 Slotline

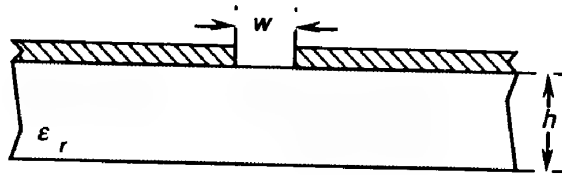


Figure 3.9.6.1: Slotline

Slotline is not a TEM structure; however, its single-sided, planar nature make it of interest for modern fabrication techniques. This structure requires a high dielectric constant material ($\epsilon_r \geq 10$) to confine the fields and minimize radiation. The equations for design are from [3] for $0.02 \leq w/h \leq 0.2$

$$Z_0 = 72.62 - 81.028 \ln \epsilon_r + \frac{50.0 (w/h - 0.02)(w/h - 0.1)}{w/h} \\ + 2.3026 \ln (w/h \times 10^2) (44.28 - 45.0849 \ln \epsilon_r) \\ - [0.7368 \ln \epsilon_r - 0.11 + w/h (2.4638 \ln \epsilon_r + 1.44)] \\ \times (11.4 - 13.9768 \ln \epsilon_r - h/\lambda \times 10^2)^2 \quad (3.9.6.1)$$

$$\epsilon_{eff} = \left(\frac{\lambda_g}{\lambda_0} \right)^{-2} \\ = [0.923 - 1.0316 \ln \epsilon_r + 0.2 w/h \\ - (0.6678 w/h + 0.1082) \ln (h/\lambda_0 \times 10^2)]^{-2} \quad (3.9.6.2)$$

and for $0.2 \leq w/h \leq 1.0$

$$Z_0 = 113.19 - 123.3042 \ln \epsilon_r + 1.25 w/h (114.59 - 119.4588 \ln \epsilon_r) \\ + 20.0 (w/h - 0.2) (1.0 - w/h) \\ - [0.15 + 0.5296 \ln \epsilon_r + w/h (-0.79 + 4.7664 \ln \epsilon_r)]$$

$$\times \left\{ [10.25 - 11.513 \ln \epsilon_r + w/h (2.1 - 3.2697 \ln \epsilon_r) - h/\lambda \times 10^2]^2 \right\} \quad (3.9.6.3)$$

$$\epsilon_{eff} = \left(\frac{\lambda_g}{\lambda_0} \right)^{-2} \\ = [0.987 - 1.1122 \ln \epsilon_r + w/h (0.111 - 0.0022 \epsilon_r) \\ - (0.2786 + 0.2164 w/h - 0.007368 \epsilon_r) \ln (h/\lambda \times 10^2)]^{-2} \quad (3.9.6.4)$$

Equations are curve fits to data calculated by the technique of [2] and are accurate to 2.0% for

$$9.7 \leq \epsilon_r \leq 20.0$$

$$0.02 \leq w/h \leq 1.0$$

$$0.01 \leq h/\lambda \leq (h/\lambda)_0$$

The TE_{10} mode cutoff frequency is given by

$$(h/\lambda)_0 = \frac{0.25}{\sqrt{\epsilon_r - 1.0}} \quad (3.9.6.5)$$

For lower dielectric constants ($2.2 \leq \epsilon_r \leq 3.8$) Janaswamy and Schaubert [7] give (for $0.0015 \leq w/\lambda_0 \leq 0.075$):

$$Z_0 = 60.0 + 3.69 \sin \left[\frac{(\epsilon_r - 2.22) \pi}{2.36} \right] + 133.5 \ln (10.0 \epsilon_r) \sqrt{w/\lambda_0} \\ + 2.81 [1.0 - 0.011 \epsilon_r (4.48 + \ln \epsilon_r)] (w/h) \ln (100.0 h/\lambda_0) \\ + 131.1 (1.028 - \ln \epsilon_r) \sqrt{h/\lambda_0} \\ + 12.48 (1.0 + 0.18 \ln \epsilon_r) \frac{w/h}{\sqrt{\epsilon_r - 2.06 + 0.85 (w/h)^2}} \quad (3.9.6.6)$$

$$\begin{aligned}\epsilon_{eff} &= \left(\frac{\lambda_g}{\lambda_0}\right)^{-2} \\ &= \left\{ 1.045 - 0.365 \ln \epsilon_r + \frac{6.4 (w/h) \epsilon_r^{0.945}}{238.64 + 100.0 w/h} \right. \\ &\quad \left. - \left[0.148 - \frac{8.81 (\epsilon_r + 0.95)}{100.0 \epsilon_r} \right] \ln \frac{h}{\lambda_0} \right\}^{-2} \quad (3.9.6.7)\end{aligned}$$

For $0.075 \leq w/\lambda_0 \leq 1.0$

$$\begin{aligned}Z_0 &= 133.0 + 10.34 (\epsilon_r - 1.0)^2 + 2.87 [2.96 + (\epsilon_r - 1.582)^2] \\ &\quad \left\{ (w/h + 2.32 \epsilon_r - 0.56) \left[(32.5 - 6.67 \epsilon_r)(100.0 h/\lambda_0)^2 - 1.0 \right] \right\}^{1/2} \\ &\quad - (684.45 h/\lambda_0)(\epsilon_r + 1.35)^2 + 13.23 [(\epsilon_r - 1.722) w/\lambda_0]^2 \quad (3.9.6.8)\end{aligned}$$

$$\begin{aligned}\epsilon_{eff} &= \left(\frac{\lambda_g}{\lambda_0}\right)^{-2} \\ &= \left\{ 1.194 - 0.24 \ln \epsilon_r - \frac{0.621 \epsilon_r^{0.835} (w/\lambda_0)^{0.48}}{1.344 + w/h} \right. \\ &\quad \left. - 0.0617 \left[1.91 - \frac{\epsilon_r + 2.0}{\epsilon_r} \right] \ln \frac{h}{\lambda_0} \right\}^{-2} \quad (3.9.6.9)\end{aligned}$$

The above are accurate (first set: $Z_0 \leq 2.7\%$, $\lambda_g/\lambda_0 \leq 2.2\%$; second set: $Z_0 \leq 5.4\%$, $\lambda_g/\lambda_0 \leq 2.6\%$) for

$$2.2 \leq \epsilon_r \leq 3.8$$

$$0.006 \leq h/\lambda_0 \leq 0.060$$

For intermediate dielectric constants ($3.8 \leq \epsilon_r \leq 9.8$) and $0.0015 \leq w/\lambda_0 \leq 0.075$

$$Z_0 = 73.6 - 2.15 \epsilon_r + (638.9 - 31.37 \epsilon_r) (w/\lambda_0)^{0.6}$$

$$\begin{aligned}&+ \left(36.23 \sqrt{\epsilon_r^2 + 41.0} - 225.0 \right) \frac{w/h}{w/h + 0.876 \epsilon_r - 2.0} \\ &+ 0.51 (\epsilon_r + 2.12) (w/h) \ln \frac{100.0 h}{\lambda_0} \\ &- 0.753 \epsilon_r \frac{h/\lambda_0}{\sqrt{w/\lambda_0}} \quad (3.9.6.10)\end{aligned}$$

$$\begin{aligned}\epsilon_{eff} &= \left(\frac{\lambda_g}{\lambda_0}\right)^{-2} \\ &= \left\{ 0.9217 - 0.277 \ln \epsilon_r + 0.0322 (w/h) \sqrt{\frac{\epsilon_r}{w/h + 0.435}} \right. \\ &\quad \left. - 0.01 \ln (h/\lambda_0) \left[4.6 - \frac{3.65}{\epsilon_r^2 \sqrt{w/\lambda_0} (9.06 - 100 w/\lambda_0)} \right] \right\}^{-2} \quad (3.9.6.11)\end{aligned}$$

For $0.075 \leq w/\lambda_0 \leq 1.0$

$$\begin{aligned}Z_0 &= 120.75 - 3.74 \epsilon_r + 50.0 [\tan^{-1} (2.0 \epsilon_r) - 0.8] \\ &\quad (w/h)^{\{1.11 + [0.132 (\epsilon_r - 27.7) / (100.0 h/\lambda_0 + 5.0)]\}} \\ &\quad \ln [100 h/\lambda_0 + \sqrt{(100 h/\lambda_0)^2 + 1.0}] \\ &\quad + 14.21 (1.0 - 0.458 \epsilon_r) (100.0 h/\lambda_0 + 5.1 \ln \epsilon_r - 13.1) (w/\lambda_0 + 0.33)^2 \quad (3.9.6.12)\end{aligned}$$

$$\begin{aligned}\epsilon_{eff} &= \left(\frac{\lambda_g}{\lambda_0}\right)^{-2} \\ &= \left\{ 1.05 - 0.04 \epsilon_r + 1.411 \times 10^{-2} (\epsilon_r - 1.421) \right\}\end{aligned}$$

$$\begin{aligned}
& \ln [w / h - 2.012 (1.0 - 0.146 \epsilon_r)] \\
& + 0.111 (1.0 - 0.366 \epsilon_r) \sqrt{w / \lambda_0} \\
& + 0.139 [1.0 + 0.52 \epsilon_r \ln (14.7 - \epsilon_r)] (h / \lambda_0) \ln (h / \lambda_0) \}^{-2}
\end{aligned}
\tag{3.9.6.13}$$

The above are accurate (first set: Z_0 accurate to 5.4%, ϵ_{eff} accurate to 3%; second set: Z_0 accurate to 5.8%, ϵ_{eff} accurate to 3.2%) for

$$3.8 \leq \epsilon_r \leq 9.8$$

Cohn [2] indicates that losses may be approximated with the losses of microstrip line. Kirschning *et al.* [8] analyzes slotline between two dielectrics.

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3.9.7 Finline

Again, this is not a TEM structure but because of its planar nature it is very useful for fabricating circuits. It is also easy to interface to waveguide because of its fully shielded structure. Propagation is a hybrid mode that is useful mainly in the 30 to 100 GHz range. There are several different forms of finline. Metallization can be on a single side of the dielectric, both sides, or diagonal. The finline may be captivated in the shield and grounded or isolated.

3.9.7.1 Unilateral Finline

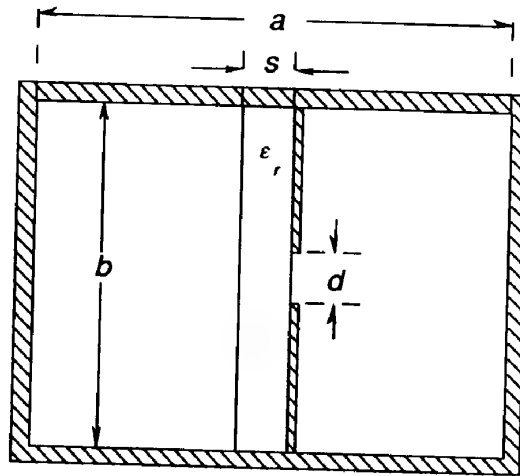


Figure 3.9.7.1: Unilateral Finline

Unilateral finline closed-form equations are in Praminick and Bhartia [10]

$$Z_0 = \frac{240.0 \pi^2 (P X + Q) \frac{b}{a}}{\sqrt{\epsilon_{eff}(f)} (0.385 X + 1.762)^2} \quad (3.9.7.1.1)$$

$$\epsilon_{eff}(f) = k_e - \left(\frac{\lambda_0}{\lambda_{ca}} \right)^2 \quad (3.9.7.1.2)$$

where for $d/b > 0.3$

$$P = -0.763 \left(\frac{b}{\lambda_0} \right)^2 + 0.58 \left(\frac{b}{\lambda_0} \right) + 0.0775 \left[\ln \left(\frac{a}{s} \right) \right]^2 - 0.668 \left(\ln \frac{a}{s} \right) + 1.262 \quad (3.9.7.1.3)$$

$$Q = 0.372 \left(\frac{b}{\lambda} \right) + 0.914 \quad (3.9.7.1.4)$$

and for $d/b \leq 0.3$

$$P = 0.17 \left(\frac{b}{\lambda} \right) + 0.0098 \quad (3.9.7.1.5)$$

$$Q = 0.138 \left(\frac{b}{\lambda} \right) + 0.873 \quad (3.9.7.1.6)$$

where

$$X = \ln \csc \left(\frac{\pi d}{2.0 b} \right) \quad (3.9.7.1.7)$$

For thin dielectric and low ϵ_r [11] gives

$$k_e \approx k_c = \left(\frac{\lambda_{cd}}{\lambda_{ca}} \right)^2 \quad (3.9.7.1.8)$$

$$\lambda_{cd} = 2.0 (a - s) \sqrt{1.0 + N \left[1.0 - \frac{s}{a} (-0.0769 \epsilon_r + 1.231) F \left(\frac{s}{a} \right) X_a \right]} \quad (3.9.7.1.9)$$

$$\lambda_{ca} = 2.0 a \sqrt{1.0 + \frac{4.0}{\pi} \frac{b}{a} \left(1.0 + 0.2 \sqrt{\frac{b}{a}} \right) \ln \csc \left(\frac{\pi d}{2.0 b} \right)} \quad (3.9.7.1.10)$$

$$N = \frac{4.0}{\pi} \frac{b}{a - s} \left(1.0 + 0.2 \sqrt{\frac{b}{a - s}} \right) \quad (3.9.7.1.11)$$

$$X_a = \ln \csc \left(\frac{\pi d}{2.0 b} \right) + \epsilon_r (G_d + G_a) \quad (3.9.7.1.12)$$

$$G_a = \eta_a \tan^{-1} \left(\frac{1.0}{\eta_a} \right) + \ln \sqrt{1.0 + \eta_a^2} \quad (3.9.7.1.13)$$

$$G_d = \eta_d \tan^{-1} \left(\frac{1.0}{\eta_d} \right) + \ln \sqrt{1.0 + \eta_d^2} \quad (3.9.7.1.14)$$

$$\eta_a = \frac{\eta_d d}{b} \quad (3.9.7.1.15)$$

$$\eta_d = \frac{s/a}{b d / a b} \quad (3.9.7.1.16)$$

where for $d/b \leq 0.5$

$$F\left(\frac{s}{a}\right) = -25.1223 + 31.524 \ln \frac{a}{s} - 12.504 \left(\ln \frac{a}{s} \right)^2 + 1.9454 \left(\ln \frac{a}{s} \right)^3 \quad (3.9.7.1.17)$$

and for $0.5 \leq d/b \leq 0.75$

$$F\left(\frac{s}{a}\right) = -33.934 + 42.451 \ln \frac{a}{s} - 17.057 \left(\ln \frac{a}{s} \right)^2 + 2.5885 \left(\ln \frac{a}{s} \right)^3 \quad (3.9.7.1.18)$$

and for $0.75 \leq d/b \leq 1.0$

$$F\left(\frac{s}{a}\right) = -48.0487 + 59.2846 \ln \frac{a}{s} - 23.77 \left(\ln \frac{a}{s} \right)^2 + 3.47 \left(\ln \frac{a}{s} \right)^3 \quad (3.9.7.1.19)$$

The accuracy of the equation for Z_0 is $\pm 2.0\%$ for $s/a \leq 0.05$ and $\pm 3\%$ for $s/a > 0.05$.

3.9.7.2 Finite Thickness Off-center Unilateral Finline

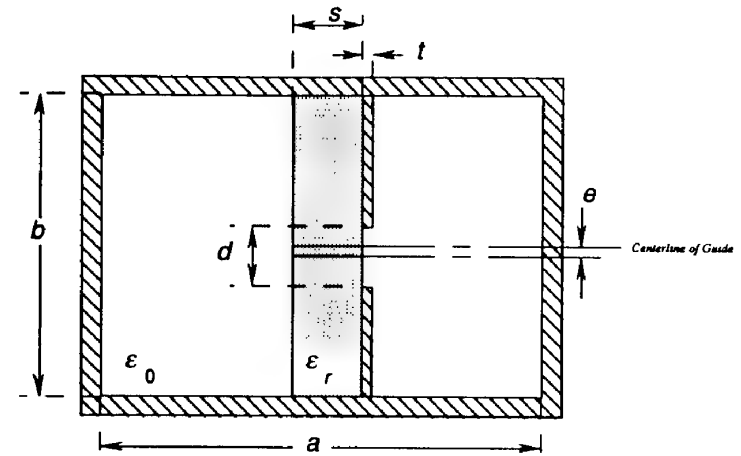


Figure 3.9.7.2: Finite Thickness, Off-Center Finline

Pramanick *et al.* [14] include the effects of finite thickness metallization and an off-center fin as shown in Figure 3.9.7.2. The equations are

$$Z_0 = \frac{120.0 \pi^2 d}{\lambda_{ca} \sqrt{\epsilon_{eff}(f)} (f_1 + q f_2)} \left[\frac{k_e(f) - 1.0}{k_c - 1.0} \right] \quad (3.9.7.2.1)$$

$$\epsilon_{eff}(f) = k_e(f) - \left(\frac{\lambda_0}{\lambda_{ca}} \right)^2 \quad (3.9.7.2.2)$$

where

$$f_1 = \frac{2.0 d}{\lambda_{ca}} \cos^2 \left(\frac{\pi t}{\lambda_{ca}} \right) + \frac{\pi t}{2.0 \lambda_{ca}} + \frac{\sin^2 \left(\frac{\pi t}{\lambda_{ca}} \right)}{4.0} \quad (3.9.7.2.3)$$

$$f_2 = \frac{\pi (a - t)}{2.0 \lambda_{ca}} - \frac{\sin^4 \left[\frac{\pi (a - t)}{\lambda_{ca}} \right]}{4.0} \quad (3.9.7.2.4)$$

$$q = \left[\frac{\frac{d}{b} \cos^2 \left(\frac{\pi t}{\lambda_{ca}} \right)}{\sin^2 \left(\frac{a-t}{\lambda_{ca}} \right)} \right] \quad (3.9.7.2.5)$$

$$k_e(f) = \epsilon_r - \frac{\epsilon_r - k_c}{1.0 + \left[\frac{s}{a} \sqrt{\frac{f - f_c}{f}} \right]^{(1.0 + 0.06 b/d)}} \quad (3.9.7.2.6)$$

$$k_c = 1.0 + \frac{s}{a} (a_1 X + b_1) (\epsilon_r - 1.0) \quad (3.9.7.2.7)$$

$$\lambda_{ca} = 2.0 (a - t) \sqrt{1.0 + N X + \left(2.75 + 0.2 \frac{t}{a} \right) \left(\frac{t}{a-t} \right) b/d} \quad (3.9.7.2.8)$$

$$f_c = \frac{c}{\lambda_{ca}}$$

$$N = \frac{4.0}{\pi} \left(1.0 + 0.2 \sqrt{\frac{b}{a-t}} \right) \left(\frac{b}{a-t} \right) \quad (3.9.7.2.9)$$

$$a_1 = p_0 + p_1 \ln \frac{a}{s} + p_2 \left(\ln \frac{a}{s} \right)^2 \quad (3.9.7.2.10)$$

$$b_1 = q_0 + q_1 \ln \frac{a}{s} + q_2 \left(\ln \frac{a}{s} \right)^2 \quad (3.9.7.2.11)$$

and

$$p_0 = 2.455 - 69.366 \left(\frac{t}{a} \right) + 1826.189 \left(\frac{t}{a} \right)^2 \quad (3.9.7.2.12)$$

$$p_1 = 2.5646 + 54.2675 \left(\frac{t}{a} \right) - 1381.6 \left(\frac{t}{a} \right)^2 \quad (3.9.7.2.13)$$

$$p_2 = 0.7837 - 18.954 \left(\frac{t}{a} \right) + 466.3185 \left(\frac{t}{a} \right)^2 \quad (3.9.7.2.14)$$

$$q_0 = 2.1697 + 3.638 \left(\frac{t}{a} \right) + 240.356 \left(\frac{t}{a} \right)^2 \quad (3.9.7.2.15)$$

$$q_1 = 3.4661 - 7.3661 \left(\frac{t}{a} \right) - 180.6 \left(\frac{t}{a} \right)^2 \quad (3.9.7.2.16)$$

$$q_2 = -0.6149 + 3.11624 \left(\frac{t}{a} \right) + 19.605 \left(\frac{t}{a} \right)^2 \quad (3.9.7.2.17)$$

For a centered fin gap use

$$X = -\ln \left[\sin \left(\frac{0.5 \pi d}{b} \right) \right] \quad (3.9.7.2.18)$$

For an off-center fin gap replace this with

$$X = \ln \left\{ \csc \left(\frac{0.5 \pi d}{b} \right) \csc \left[0.5 \pi (1.0 - 2.0 e/b) \right] \right\} \quad (3.9.7.2.19)$$

The equations are valid for

$$\epsilon_r \leq 4.0$$

$$t/a \leq 0.02094$$

$$1.0/64.0 \leq s/a \leq 1.0/8.0$$

$$1.0/32.0 \leq d/b \leq 1.0/2.0$$

Results are valid to within $\pm 0.8\%$ of the complex power calculations.

3.9.7.3 Bilateral Finline

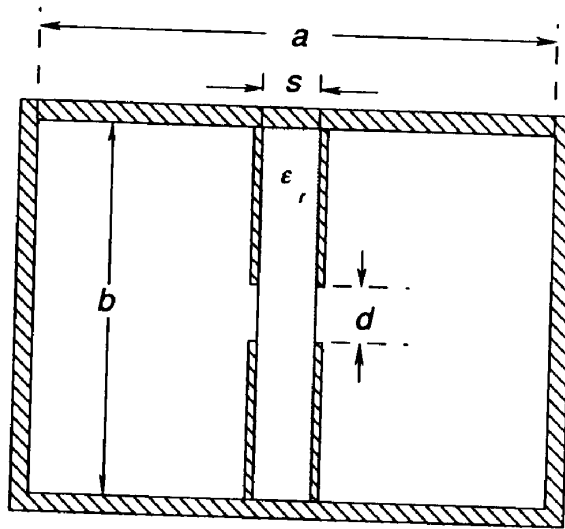


Figure 3.9.7.3: Bilateral Finline

Bilateral finline closed-form equations are in Praminick and Bhartia [10]

$$Z_0 = \frac{240.0 \pi^2 (P X + Q) \frac{b}{a-s}}{\sqrt{\epsilon_{eff}(f)} (0.385 X + 1.762)^2} \quad (3.9.7.3.1)$$

$$\epsilon_{eff}(f) = k_c - \left(\frac{\lambda_0}{\lambda_{ca}} \right)^2 \quad (3.9.7.3.2)$$

where

$$P = 0.01 \left(\frac{b}{\lambda_0} \right)^2 + 0.097 \left(\frac{b}{\lambda_0} \right) + 0.04095 \quad (3.9.7.3.3)$$

$$Q = 0.0031 \left(\frac{b}{\lambda_0} \right) + 0.89 \quad (3.9.7.3.4)$$

$$X = X_b - G_d (\epsilon_r - 1.0) \quad (3.9.7.3.5)$$

$$X_b = \ln \csc \left(\frac{\pi d}{2.0 b} \right) + \epsilon_r G_d \quad (3.9.7.3.6)$$

$$G_d = \eta_d \arctan \left(\frac{1.0}{\eta_d} \right) + \ln (1.0 + \eta_d^2)^{1/2} \quad (3.9.7.3.7)$$

$$\eta_d = \frac{(s/a)}{(b/a)(d/b)} \quad (3.9.7.3.8)$$

For low ϵ_r and thin dielectric

$$k_e = k_c = \left(\frac{\lambda_{cd}}{\lambda_{ca}} \right)^2 \quad (3.9.7.3.9)$$

where

$$\lambda_{cd} = 2.0 (a-s) \sqrt{1.0 + N X_b} \quad (3.9.7.3.10)$$

X_b is defined above and

$$N = \frac{4.0}{\pi} \frac{b}{a-s} \left(1.0 + 0.2 \sqrt{\frac{b}{a-s}} \right) \quad (3.9.7.3.11)$$

To calculate λ_{ca} , set $\epsilon_r = 1.0$ in the equations for λ_{cd} . The accuracy of ϵ_{eff} is stated to be within $\pm 0.8\%$ for $s/a \leq 0.125$ and Z_0 is accurate to $\pm 2\%$ for $d/b \leq 0.2$.

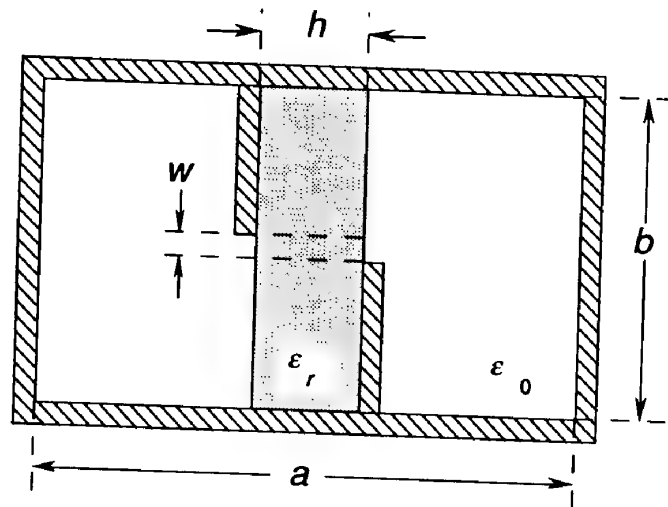
3.9.7.4 *Antipodal Finline*

Figure 3.9.7.4: Antipodal Finline

Antipodal finline provides a wider range of impedance than the previous finline structures. For sufficiently large gap w , and thin dielectric the equations of unilateral finline can be used for design. Also see [1] for design curves.

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3.9.8 Others

Other useful types of transmission lines not covered in this book include image guide, H-guide (100 to 200 GHz), groove guide (100 to 300 GHz), and microguide.

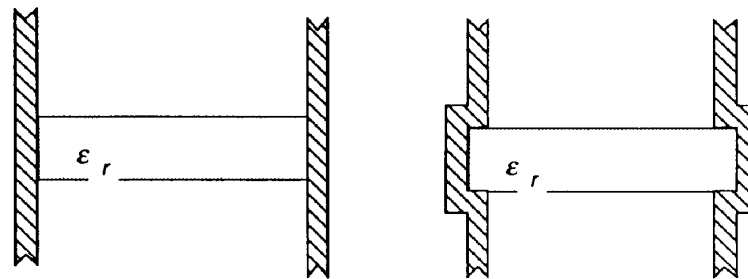


Figure 3.9.8.1: H-Guide and Groove Guide

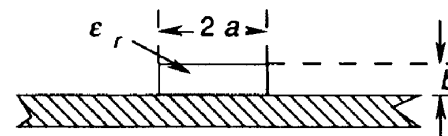


Figure 3.9.8.2: Image Guide

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3.10 WAVEGUIDES

The analysis of waveguides is often used when analyzing planar transmission line discontinuities. An equivalence can be established between a planar line discontinuity and a waveguide discontinuity allowing existing waveguide models to be transformed. The analysis of waveguides is a very large topic in and of itself and will not be covered in depth here. There are a number of excellent references. See Marcuvitz [2], for example.

3.10.1 Rectangular Waveguide

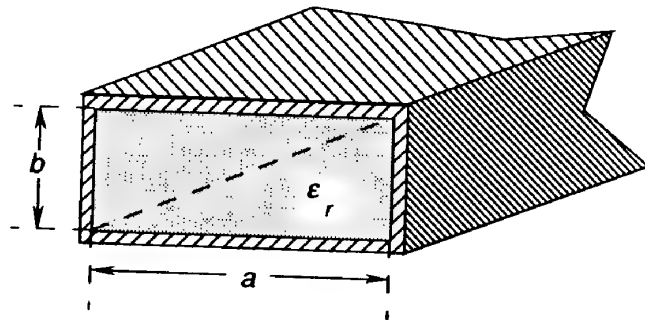


Figure 3.10.1.1: Rectangular Waveguide

Waveguide has a major difference from most of the other transmission lines described so far—it does not pass dc. As such, it behaves somewhat like a high-pass filter. There are many *modes* or solutions of the wave equations in the guide and each has an associated *cutoff wavelength* or *cutoff frequency*. A wave can propagate only when the frequency is greater than the cutoff frequency. Because there is a single conductor, the modes of propagation cannot be *TEM*. The field patterns or modes are named *TE_{m,n}* or *TM_{m,n}*. The subscript *m* represents the number of half wavelengths in the *x* direction and *n* represents the number of half wavelengths in the *y* direction (where *z* is length). *TE* (*transverse electric*) waves, also known as *H* waves have a nonzero longitudinal magnetic field component. *TM* (*transverse magnetic*) waves, also known as *E* waves have a nonzero longitudinal electric field component. *TE* modes may have only one of *m, n* equal to zero. *TM* modes may have neither equal to zero. Conventionally, *a > b* as shown in Figure 3.10.1 above. This is assumed throughout this section.

A *degenerate mode* occurs whenever multiple field patterns have the same cutoff frequency. The lowest frequency that can actually propagate in the waveguide (the *TE_{1,0}* mode) is known as the *dominant mode* and occurs when:

$$\lambda_c = 2.0 a \quad (3.10.1.1)$$

The wavelength in the guide is:

$$\lambda_g = \frac{\lambda_0}{\sqrt{\epsilon_r \mu_r \left[1.0 - \left(\frac{\lambda_0}{\lambda_{c_{m,n}}} \right)^2 \right]}} \quad (3.10.1.2)$$

The cutoff frequencies of both *TE_{m,n}* and *TM_{m,n}* modes are:

$$\omega_{c_{m,n}} = \frac{c_0}{\sqrt{\mu_r \epsilon_r}} \sqrt{\left(\frac{m \pi}{a} \right)^2 + \left(\frac{n \pi}{b} \right)^2} \quad (\text{rad / s}) \quad (3.10.1.3)$$

The guide losses of the *TE* modes (*n* ≠ 0) are:

$$\alpha_{c_{TE_{m,n}}} = \frac{2.0 R_s}{b \eta \sqrt{1.0 - (f_c / f)^2}} \left\{ \left(1.0 + \frac{b}{a} \right) \left(\frac{f_c}{f} \right)^2 + \left[1.0 - \left(\frac{f_c}{f} \right)^2 \right] \left[\frac{(b/a) [(b/a) m^2 + n^2]}{(b^2 m^2 / a^2) + n^2} \right] \right\} \quad (3.10.1.4)$$

and the losses for the *TE_{m,0}* mode are

$$\alpha_{c_{TE_{m,0}}} = \frac{2.0 R_s}{b \eta \sqrt{1.0 - (f_c / f)^2}} \left[1.0 + \frac{2.0 a}{b} \left(\frac{f_c}{f} \right)^2 \right] \quad (3.10.1.5)$$

For the *TM* modes the attenuation is calculated with

$$\alpha_{c_{TM_{m,n}}} = \frac{2.0 R_s}{b \eta \sqrt{1.0 - (f_c / f)^2}} \frac{(b/a)^3 m^2 + n^2}{(b^2 m^2 / a^2) + n^2} \quad (3.10.1.6)$$

The dielectric losses are

$$\alpha_d = \frac{k_0^2 \epsilon_r \tan \delta}{2.0 \beta} \quad (3.10.1.7)$$

where

$$k_0^2 = k_c^2 + \beta^2 \quad (3.10.1.8)$$

$$\beta = \frac{2.0 \pi}{\lambda_g} \quad (3.10.1.9)$$

$$k_c = \frac{2.0 \pi \sqrt{\mu_r \epsilon_r}}{\lambda_c}$$

(3.10.1.10)

For the commonly used $TE_{1,0}$ mode the impedance is

$$Z_{TE_{1,0}} = \frac{\eta_0}{\sqrt{1.0 - \left(\frac{\lambda}{2.0 a}\right)^2}}$$

(3.10.1.11)

The phase and group velocities are given by

$$v_p v_g = v_c^2$$

(3.10.1.12)

$$v_g = \frac{\lambda_0}{\lambda_g} v_c = \text{velocity of energy propagation}$$

(3.10.1.13)

$$v_c = \sqrt{\mu_0 \epsilon_0} = \text{velocity of the TEM mode in free space}$$

(3.10.1.14)

$$v_p = \frac{\omega}{\beta} = \frac{k_0}{\beta} v_c = \text{speed of a constant phase point on the wave}$$

(3.10.1.15)

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3.10.2 Circular Waveguide

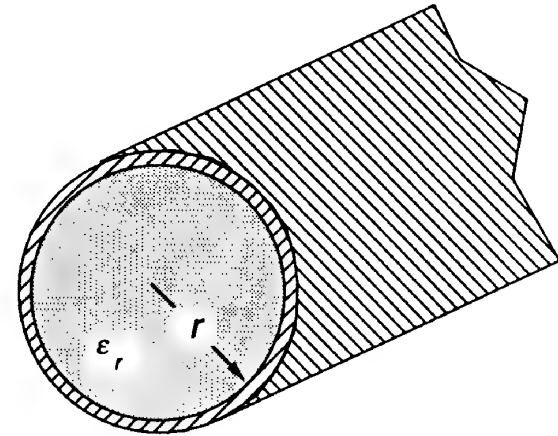


Figure 3.10.2.1: Circular Waveguide

The $TE_{m,n}$ cutoff wavelengths for this structure are

$$\lambda_c = \frac{2 \pi a \sqrt{\epsilon_r}}{u'_{m,n}} \quad (3.10.2.1)$$

The subscript m indicates angular variation in the field and n indicates radial field variations. The roots of the Bessel function are given in Marcuvitz [2]:

$u'_{m,n}$	m				
n	0	1	2	3	4
1	3.832	1.841	3.054	4.201	5.317
2	7.016	5.331	6.706	8.015	9.282
3	10.173	8.536	9.969	11.346	12.682
4	13.324	11.706	13.170		

Table 3.10.2.1: Values of $u'_{m,n}$ for $TE_{m,n}$ Mode Calculations

$TM_{m,n}$ modes are calculated in the same fashion with the following values for $u_{m,n}$.

$u_{m,n}$	m				
n	0	1	2	3	4
1	2.405	3.832	5.136	6.380	7.588
2	5.520	7.016	8.417	9.761	11.065
3	8.654	10.173	11.620	13.015	14.372
4	11.792	13.323	14.796		

Table 3.10.2.2: Values of $u_{m,n}$ for $TM_{m,n}$ Mode Calculations

The wavelength inside the waveguide is λ_g

$$\lambda_g = \frac{\lambda}{\sqrt{\epsilon_r - \left(\frac{\lambda}{\lambda_c}\right)^2}} \quad (3.10.2.2)$$

The conductor losses for TM modes the conductor loss is

$$\alpha_c = \frac{R_s}{\eta a} \frac{1.0}{\sqrt{1.0 - \left(\frac{\lambda}{\lambda_c}\right)^2}} \quad (3.10.2.3)$$

for TE modes the conductor loss is

$$\alpha_c = \frac{R_s}{\eta a} \frac{1.0}{\sqrt{1.0 - \left(\frac{\lambda}{\lambda_c}\right)^2}} \left[u'_{m,n} - m^2 + \left(\frac{\lambda}{\lambda_c}\right)^2 \right] \quad (3.10.2.4)$$

Dielectric losses are

$$\alpha_d = \frac{k_0^2 \epsilon_r \lambda_g \tan \delta}{4.0 \pi} \quad (3.10.2.5)$$

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Chapter 4

Coupled Lines

When two transmission lines are in close enough proximity for their field patterns to be disrupted (relative to their isolated patterns) it is discovered that a portion of the signal in one line is now present in the second line. This property can be desired—as when we design a directional coupler, or undesired—as in a multi-conductor cable carrying several signals. In the former case we use term the second signal *coupling*, in the latter we refer to *crosstalk*. In either case we wish to understand the source of the phenomena and the equations we will need to control it.

This chapter covers the equations for simple side-by-side transmission lines. Side-by-side lines can be used alone, in series, or in combination with other lines and components to form broadband components.

For design of symmetric, directional couplers we make use of two equations. The first determines the amount of coupling and the second designs the coupler to be matched at its ports to the impedance Z_0 .

$$k = \frac{Z_{0,e} - Z_{0,o}}{Z_{0,e} + Z_{0,o}} \quad (4.0.1)$$

$$Z_0 = \sqrt{Z_{0,e} Z_{0,o}} \quad (4.0.2)$$

or equivalently

$$Z_{0,e} = Z_0 \sqrt{\frac{1+k}{1-k}} \quad (4.0.3)$$

$$Z_{0,o} = Z_0 \sqrt{\frac{1-k}{1+k}} \quad (4.0.4)$$

The value of k is the maximum coupling, which is achieved when the electrical length of the line is $\lambda/4$. Equation 4.0.5 is the equation as a function of electrical length θ , and Figure 4.0 shows this repeating function.

$$\frac{V_2}{V_1} = \frac{j k \sin \theta}{\sqrt{1.0 - k^2 \cos \theta + j \sin \theta}} \quad (4.0.5)$$

$$k \equiv \left| \frac{V_2}{V_1} \right| \quad \theta = 90^\circ = \pi/2 \text{ radians} \quad (4.0.6)$$

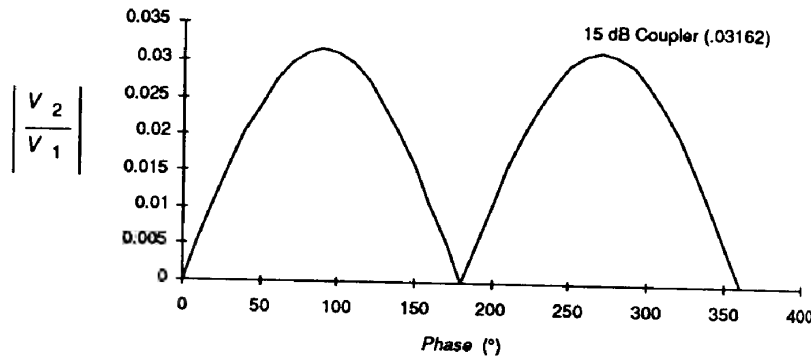


Figure 4.0: Directional Coupler Coupled Port Response

4.1 COUPLED LINE ABCD MATRIX

The physical lines described in the previous chapter are two-port networks. Coupled lines are four-port networks. Coupled lines have two ABCD matrices corresponding to the even and odd modes:

$$\begin{bmatrix} \cos \theta_e & j Z_{0,e} \sin \theta_e \\ \frac{j \sin \theta_e}{Z_{0,e}} & \cos \theta_e \end{bmatrix}$$

$$\begin{bmatrix} \cos \theta_o & j Z_{0,o} \sin \theta_o \\ \frac{j \sin \theta_o}{Z_{0,o}} & \cos \theta_o \end{bmatrix}$$

where the subscript e denotes the even mode parameters and o denotes the odd mode.

4.2 COAXIAL LINES

4.2.1 Shielded Pair

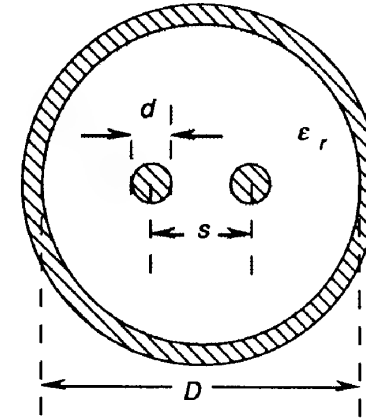


Figure 4.2.1.1: Shielded Pair

For circuits requiring differential signal transmission or low frequency shielding, the addition of an extra conductor in the standard coaxial line provides a good solution. The shielded pair may also be used for coupler, balun, or transformer designs. This structure has higher losses than coax for equivalent dimensions and dielectric.

Ramo and Whinnery [6] give the following equations:

$$Z_{0,o} = \frac{\eta_0}{\pi \sqrt{\epsilon_r}} \left\{ \ln \left[\frac{2s}{d} \frac{1.0 - (s/D)^2}{1.0 + (s/D)^2} \right] - \frac{[1.0 + 4(s/d)^2][1.0 - 4(s/D)^2]}{16(s/d)^4} \right\} \quad (4.2.1.1)$$

$$Z_{0,e} = \frac{\eta_0}{2\pi \sqrt{\epsilon_r}} \ln \left\{ \frac{s/d}{2(s/D)^2} [1.0 - (s/D)^4] \right\} \quad (4.2.1.2)$$

The above equations are valid for

$$d \ll D$$

$$d \ll s$$

Attenuation is calculated with:

$$\alpha_d = \frac{\pi \sqrt{\epsilon_r \mu_r} \epsilon_1''}{\lambda_0 \epsilon_r} \quad (4.2.1.3)$$

$$\alpha_c = \frac{R}{2.0 Z_0} \quad (4.2.1.4)$$

where

λ_0 = wavelength in free space (2.3.13)

ϵ_1'' = loss factor; see (2.4.1.1.)

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4.2.2 Shielded Twisted Pair

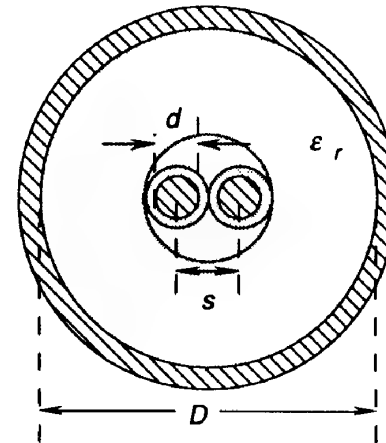


Figure 4.2.2.1: Shielded Twisted Pair

Most analyses use the equations given for the shielded pair of Section 4.2.1; however, those equations are not valid for d approaching s or d approaching D , which are common situations for this type of cable. Until exact equations are available, a reasonable starting point might be to use twisted-pair equations for the odd mode and coax equations (Chapter 3) for the even mode.

Marx [1] analyzes and gives the equations for proper termination of a shielded twisted pair. It is important to terminate both propagation modes to prevent reflections from being converted from one mode to the other. This helps minimize overshoot.

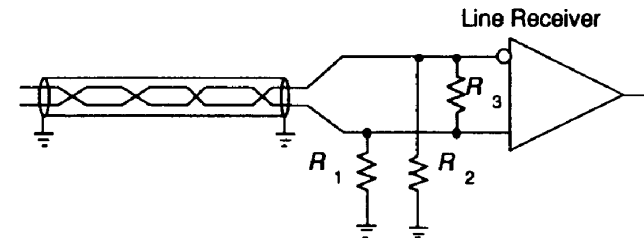


Figure 4.2.2.2: Terminating Shielded Twisted Pair

For the circuit of Figure 4.2.2.2 the termination equations are (it is assumed that the source is similarly matched to the line):

$$R_1 = R_2 = Z_{0,e} \quad (4.2.2.2.1)$$

$$R_3 = \frac{2 Z_{0,e} Z_{0,o}}{Z_{0,e} - Z_{0,o}} \quad (4.2.2.2.2)$$

and to terminate as in the circuit of Figure 4.2.2.3:

$$R_1 = R_2 = Z_{0,o} \quad (4.2.2.2.3)$$

$$R_3 = \frac{Z_{0,e} - Z_{0,o}}{2} \quad (4.2.2.2.4)$$

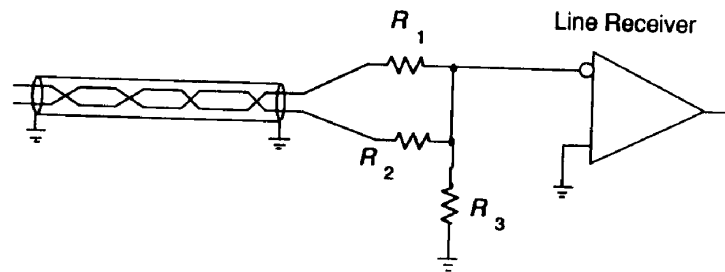


Figure 4.2.2.3: Terminating Shielded Twisted Pair

REFERENCES

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4.3 PAIRED LINES

4.3.1 Wireline™ Coupler

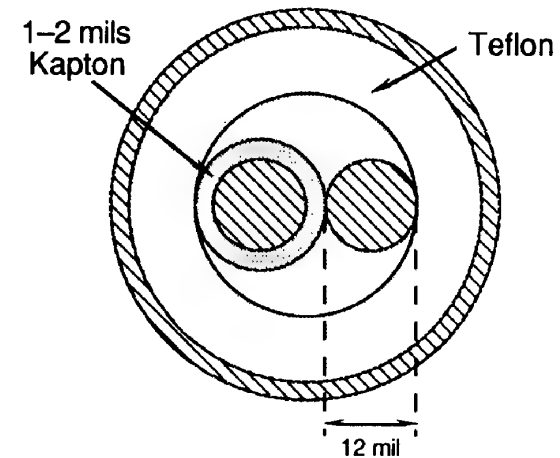


Figure 4.3.1.1: Wireline™ Coupler

This structure is a shielded twisted pair used as a coupler. It can be used as a 3 dB coupler at its center frequency or narrowband (≈ 1 octave) as any other amount of coupling by operating off the center frequency. The shield may be foil, braid, or solid tubing. The maximum frequency is limited to about 5 GHz due to strays caused by the transition from the Wireline's leads to the PCB.

A Wireline™ coupler may be designed in two ways:

- a 3 dB coupler at a desired frequency
- an X dB coupler at a desired frequency

For a 3 dB coupler, calculate the $\lambda/4$ length at the desired center frequency using $\epsilon_r = 2.55$.

For other couplers, because the spacing of the lines is fixed, we operate off of the center frequency of the coupling curve to get the desired coupling. If the desired operating frequency is f_{op} and the desired coupling (dB) is X, we calculate:

$$K_2 = \sin^{-1} \left(\sqrt{\frac{10^{-X/10}}{1.0 - 10^{-X/10}}} \right) \quad (4.3.1.1.1)$$

and

$$f_q = \text{quarter wave frequency} = f_{op} \frac{90^\circ}{K_2} \quad (4.3.1.1.2)$$

Rearranging we get an equation for the coupling over frequency:

$$X = -10 \log \left(\frac{\sin^2 (f / f_q)}{\sin^2 (f / f_q) - 1} \right) \quad (\text{dB}) \quad (4.3.1.1.3)$$

4.3.1.1 Example: 10 dB Wireline Coupler

Design a 10 dB Wireline coupler at 372.8 MHz.

Solution:

$$K_2 = \sin^{-1} \left(\sqrt{-\frac{10^{-10/10}}{10^{-10/10} - 1}} \right) = \sin^{-1} (\sqrt{1/9}) = \sin^{-1} (0.3333) = 19.47^\circ$$

$$f_q = 372.8 \text{ MHz} \times 90^\circ / 19.47^\circ = 1.7233 \text{ GHz (quarter-wave frequency)}$$

$$l = c / f_q / 4 / \sqrt{\epsilon_r} = 3 \times 10^8 \text{ m/s} / 1.7233 \text{ GHz} / 4 / \sqrt{2.55}$$

$$l = 0.0273 \text{ m} = 1.073 \text{ inches}$$

4.3.1.2 Example: 3 dB Wireline Coupler

Design a 3 dB Wireline coupler at 402 MHz.

Solution:

$$f_q = 402 \text{ MHz}$$

$$l = c / f_q / 4 / \sqrt{\epsilon_r} = 3 \times 10^8 \text{ m/s} / 402 \text{ MHz} / 4 / \sqrt{2.55}$$

$$l = 0.1168 \text{ m} = 4.6 \text{ inches}$$

REFERENCES

- [1] Saad, Theodore S., "Wireline™," 1987 SBMO International Microwave Symposium Proceedings, Rio de Janeiro, pp. 563–569.
- [2] Sage Laboratories, "Wireline® Quadrature Hybrids and Couplers."

4.3.2 Coupled Rectangular Coaxial Lines

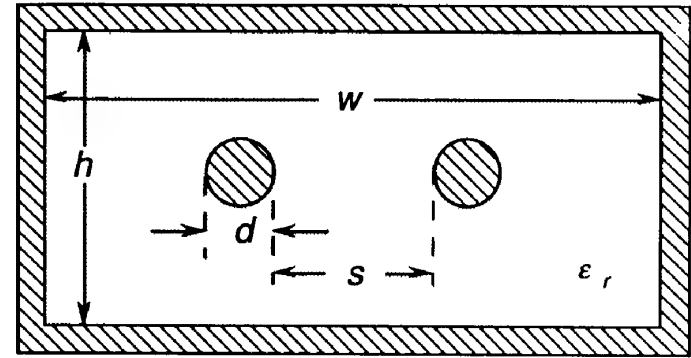


Figure 4.3.2.1: Coupled Lines in Rectangular Coaxial

This configuration is found in interdigital filters and when we have close side walls in slab line designs. Agarwal [1] modifies the (no side wall) equations of Stracca *et al.* [3] for slab line and gives

$$Z_{0,e} = \frac{\eta_0}{2.0 \pi \sqrt{\epsilon_r}} \ln \left(\frac{4.0}{\pi \frac{d}{h} F F_e} \right)$$

$$- \frac{138}{\sqrt{\epsilon_r}} \sum_{m=1}^{\infty} (-1)^m \log \frac{\tanh \left\{ \frac{\pi}{h} \left[m \left(D + C / 2.0 \right) - C / 2.0 \right] \right\}}{\tanh \left\{ \frac{\pi}{h} \left[m \left(D + C / 2.0 \right) + C / 2.0 \right] \right\}}$$

(4.3.2.1.1)

$$Z_{0,o} = \frac{\eta_0}{2.0 \pi \sqrt{\epsilon_r}} \ln \left(\frac{4.0}{\pi \frac{d}{h} F F_o} \right)$$

$$\frac{138}{\sqrt{\epsilon_r}} \sum_{m=1}^{\infty} \log \frac{1.0 + \left\{ \sinh \left(\frac{\pi c}{2.0 h} \right) / \left[(D + C / 2.0) \cosh \left(\frac{\pi m}{h} \right) \right] \right\}}{1.0 - \left\{ \sinh \left(\frac{\pi c}{2.0 h} \right) / \sinh \left[\frac{\pi m (D + C / 2.0)}{h} \right] \right\}}$$

(4.3.2.1.2)

where

$$F_e = M_e \tanh \left(\frac{\pi c}{2.0 h} \right) \quad (4.3.2.1.3)$$

$$F_o = M_o \coth \left(\frac{\pi c}{2.0 h} \right) \quad (4.3.2.1.4)$$

$$M_e = 1.0 + \exp [(-1.8861 x^2 + 1.4177 x - 5.5142)y] \\ - \exp [(-2.7030 x^2 + 4.6772 x - 8.5900)y] \quad (4.3.2.1.5)$$

$$M_o = 1.0 + \exp [(3.0989 x^2 - 8.2997 x - 3.0843)y] \\ - \exp [(4.0797 x^2 - 2.1808 x - 8.5534)y] \quad (4.3.2.1.6)$$

$$F = \frac{[1.0 + e^{16.0 (d/h) - 1.142}]}{\sqrt{1.0 - \frac{(d/h)^4}{5.905}}} \quad (4.3.2.1.7)$$

and

$$x = d/h \quad (4.3.2.1.8)$$

$$y = s/d \quad (4.3.2.1.9)$$

The accuracy of these equations was checked by designing a 10 dB coupler. The experimental results showed good agreement with the calculation. The original equations of Stracca *et al.* are curve fits to numerical results. These equations are accurate to 1.4% for $Z_{0,e}$ and 3.2% for $Z_{0,o}$ for

$$d/h \leq 0.9$$

$$s/d \geq 0.3$$

REFERENCES

- [1] Agarwal, A., *et al.*, "Coupled Bars in Rectangular Coaxial," *Electronics Letters*, Vol. 25, No. 1, January 5, 1989, pp. 66-67. (Bars of circular cross-section.)
- [2] Perlow, Stewart M., "Analysis of Edge-Coupled Shielded Strip and Slabline Structures," *IEEE Transactions on Microwave Theory and Techniques*, Vol. MTT-35, No. 5, May 1987, pp. 522-529. (Bars of rectangular and unequal cross-section.)
- [3] Stracca, Giovanni B., *et al.*, "Numerical Analysis of Various Configurations of Slab Lines," *IEEE Transactions on Microwave Theory and Techniques*, Vol. MTT-34, No. 3, March 1986, pp. 359-363.

4.4 COUPLED COPLANAR LINES

4.4.1 Broadside Coupled Coplanar Waveguide

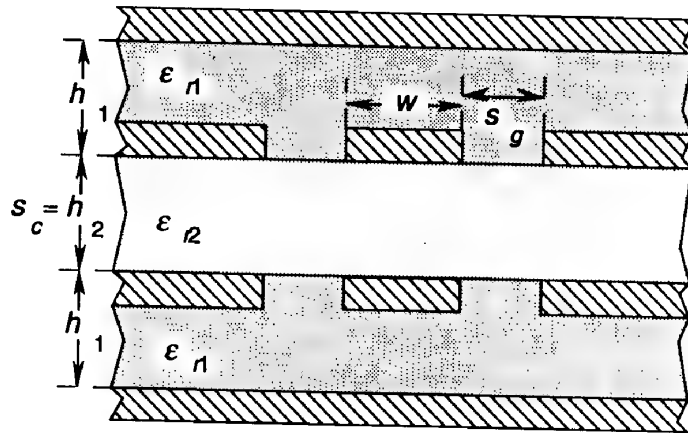


Figure 4.4.1.1: Broadside Coupled Coplanar Waveguide

This multiple dielectric configuration is useful in integrated circuits. Bedair and Wolff [1] gives the even and mode impedances in terms of the the even and odd mode capacitances:

$$Z_{0,o} = \frac{1.0}{c \sqrt{C_o C_o^a}} \quad (4.4.1.1)$$

$$Z_{0,e} = \frac{1.0}{c \sqrt{C_e C_e^a}} \quad (4.4.1.2)$$

$$\epsilon_{eff,o} = \frac{C_o}{C_o^a} \quad (4.4.1.3)$$

$$\epsilon_{eff,e} = \frac{C_e}{C_e^a} \quad (4.4.1.4)$$

$$C_o = C_{o,1} + C_{o,2} \quad (4.4.1.5)$$

$$C_{o,1} = 2.0 \epsilon_0 \epsilon_{r1} \frac{K(k_{o,1})}{K(k'_{o,1})} \quad (4.4.1.6)$$

$$C_{o,2} = 2.0 \epsilon_0 \epsilon_{r2} \frac{K(k_{o,2})}{K(k'_{o,2})} \quad (4.4.1.7)$$

$$C_e = 2.0 \epsilon_0 \epsilon_{r1} \frac{K(k_{e,1})}{K(k'_{e,1})} + 2.0 \epsilon_0 \epsilon_{r2} \frac{K(k_{e,2})}{K(k'_{e,2})} \quad (4.4.1.8)$$

$$k_{o,1} = \frac{\tanh \frac{\pi w}{4 h_1}}{\tanh \frac{\pi (w + 2 s_g)}{4 h_1}} \quad (4.4.1.9)$$

$$k_{o,2} = \frac{\tanh \frac{\pi w}{4 h_2}}{\tanh \frac{\pi (w + 2 s_g)}{4 h_2}} \quad (4.4.1.10)$$

$$k_{e,2} = \frac{\sinh \frac{\pi w}{4 h_2}}{\sinh \frac{\pi (w + 2 s_g)}{4 h_2}} \quad (4.4.1.11)$$

and C_i^a ($i = e, o$) are the capacitances when the dielectric is replaced with air ($\epsilon_r = 1$). As before,

$$k_i' = \sqrt{1 - k_i^2} \quad (4.4.1.12)$$

These equations are quasi-static and are exact at zero frequency and for coplanar propagation. The authors state that these equations are useful to 40 GHz for small s_g . A graph of dispersion for two cases is given in Bedair and Wolff [2]; however, no general conclusions or formulas are available. For coplanar waveguide-like propagation to within 1%, $s_g / h \geq 0.56$ is recommended.

REFERENCES

- [1] Bedair, Said S., and Ingo Wolff, "Computer Oriented Analytic Formulas for the Quasi-TEM Parameters of Broadside-Coupled Coplanar Waveguides with a General Dielectric for MMICs," *18th European Microwave Conference*

Proceedings, Stockholm, Sweden, 1988, pp.1139–1144. (Omits the factor of 2 in front of $C_{0,1}$, $C_{0,2}$, and in C_e .)

- [2] Bedair, Said S., and Ingo Wolff, "Fast and Accurate Analytic Formulas for Calculating the Parameters of a General Broadside-Coupled Coplanar Waveguide for (M)MIC Applications," *IEEE Transactions on Microwave Theory and Techniques*, Vol. MTT-37, No. 5, May 1989, pp. 843–850.

4.4.2 Edge-Coupled CPW

Edge-coupled coplanar waveguide (CPW) has all of the advantages of the planarity of the coplanar structure together with an increased possible isolation over microstrip line structures.

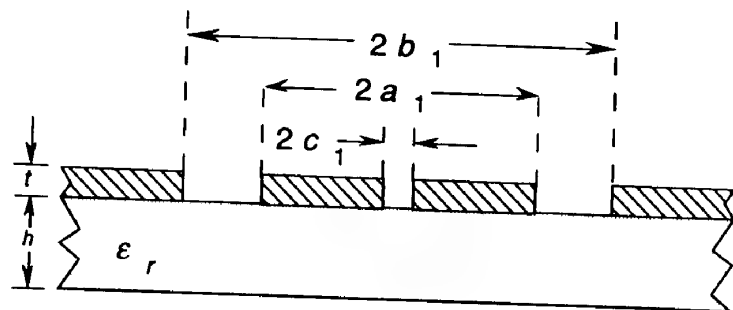


Figure 4.4.2.1: Edge-Coupled CPW

$$Z_{0,o} = \frac{\eta_0}{2} \frac{K'(k_5) c}{\sqrt{\epsilon_{eff}} K(k_5) 754} \quad (\Omega) \quad (4.4.2.1)$$

$$Z_{0,e} = 2 Z_{0,k_1} \frac{K(k_7) K'(k_1)}{K'(k_7) K(k_1)} \quad (\Omega) \quad (4.4.2.2)$$

$$\epsilon_{eff} = \sqrt{\frac{\epsilon_r + 1}{2}} \quad (4.4.2.3)$$

where

$$k = \frac{Z_{0,e} - Z_{0,o}}{Z_{0,e} + Z_{0,o}} \quad (4.4.2.4)$$

$$k_1 = \frac{a_1}{b_1} \quad (4.4.2.5)$$

$$k_5 = \frac{c_5}{a_5} \text{ which is found by solving } \frac{F(k_1, k_5)}{K(k_1)} = 1.0 - \frac{F(k_1, c_1/a_1)}{K(k_1)} \quad (4.4.2.6)$$

$$k_7 = \frac{c_7}{b_7} \text{ which is found by solving } F(k_1, c_7/b_7) = 1.0 - \frac{F(k_1, c_1/a_1)}{K(k_1)} \quad (4.4.2.7)$$

where $F(k_1, m)$ is the elliptic integral of the first kind.

In Bandrand *et al.* [1] a coupled CPW is described that is fabricated as an active device (FET) to allow DC programmability of the coupling.

Another coupling configuration, Figure 4.4.2.2, with increased isolation between strips is described in [2]. Chang models this as two isolated strips coupled by a lumped capacitance.

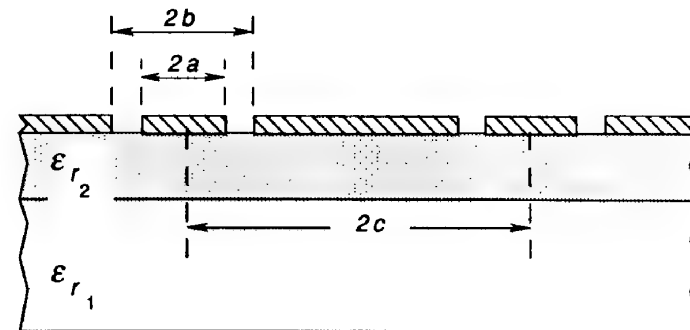


Figure 4.4.2.2: Edge-Coupled CPW

The coupling of the configuration shown in Figure 4.4.2.2 can be calculated approximately with equations of Ghione and Naldi [3].

$$Z_{0,o} = Z_{\infty} \frac{K(k_8') K(k_9)}{K(k_8) K(k_9')} \quad (4.4.2.8)$$

$$Z_{0,e} = Z_{\infty} + \Delta Z_e \quad (4.4.2.9)$$

where

Z_{∞} = impedance of an isolated CPW with dimensions $2a$, $2b$ (3.4.1.1)

$$\Delta Z_e = Z_{\infty} - Z_{0,o} \quad (4.4.2.10)$$

$$k_9 = \frac{2.0 \sqrt{a b}}{a + b} \quad (4.4.2.11)$$

$$k_8 = \frac{k_9}{\sqrt{1.0 - \frac{(b-a)^2}{(2.0 c)^2}}} \quad (4.4.2.12)$$

These equations assume loose coupling and that the substrate is infinitely thick.

REFERENCES

- [1] Bandrand, H., *et al.*, "Bias-Variable Characteristics of Coupled Coplanar Waveguide on GaAs Substrate," *Electronics Letters*, Vol. 23, No. 4, February 13, 1987, pp. 171-172.
- [2] Chang, Ching Ten, and Graham A. Garcia, "Crosstalk Between Two Coplanar Waveguides," *Archiv Für Elektronik und Übertragungstechnik*, Band 43, Heft 1, 1989, pp. 55-58.
- [3] Ghione, Giovanni, and Carlo U. Naldi, "Coplanar Waveguides for MMIC Applications: Effect of Upper Shielding Conductor Backing, Finite-Extent Ground Planes, and Line-to-Line Coupling," *IEEE Transactions on Microwave Theory and Techniques*, Vol. MTT-35, No. 3, March 1987, pp. 260-267.
- [4] Wen, Cheng P., "Coplanar-Waveguide Directional Couplers," *IEEE Transactions on Microwave Theory and Techniques*, Vol. MTT-18, No. 6, June 1970, pp. 318-322.

4.4.3 Edge-Coupled CPWG

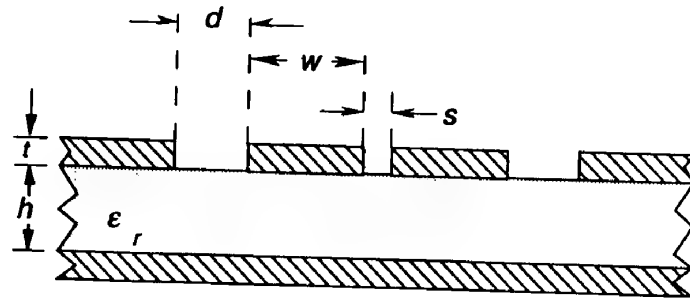


Figure 4.4.3.1: Edge-Coupled CPWG

$$Z_{0,o} = \frac{\eta_0}{\sqrt{\epsilon_{eff,o}}} \left[\frac{1.0}{2.0 \frac{K(k_o)}{K'(k_o)} + \frac{K(\beta_1)}{K'(\beta_1)}} \right] (\Omega) \quad (4.4.3.1)$$

$$Z_{0,e} = \frac{\eta_0}{\sqrt{\epsilon_{eff,e}}} \left[\frac{1.0}{2.0 \frac{K(k_e)}{K'(k_e)} + \frac{K(\beta_1 k_1)}{K'(\beta_1 k_1)}} \right] (\Omega) \quad (4.4.3.2)$$

$$\epsilon_{eff,o} = \frac{2.0 \epsilon_r \frac{K(k_o)}{K'(k_o)} + \frac{K(\beta_1)}{K'(\beta_1)}}{2.0 \frac{K(k_o)}{K'(k_o)} + \frac{K(\beta_1)}{K'(\beta_1)}} \quad (4.4.3.3)$$

$$\epsilon_{eff,e} = \frac{2.0 \epsilon_r \frac{K(k_e)}{K'(k_e)} + \frac{K(\beta_1 k_1)}{K'(\beta_1 k_1)}}{2.0 \frac{K(k_e)}{K'(k_e)} + \frac{K(\beta_1 k_1)}{K'(\beta_1 k_1)}} \quad (4.4.3.4)$$

where

$$k_o = \Lambda \frac{-\sqrt{\Lambda^2 - t_c^2} + \sqrt{\Lambda^2 - t_B^2}}{t_B \sqrt{\Lambda^2 - t_c^2} + t_c \sqrt{\Lambda^2 - t_B^2}} \quad (4.4.3.5)$$

$$k_e = \Lambda' \frac{-\sqrt{\Lambda'^2 - t_c'^2} + \sqrt{\Lambda'^2 - t_B'^2}}{t_B' \sqrt{\Lambda'^2 - t_c'^2} + t_c' \sqrt{\Lambda'^2 - t_B'^2}} \quad (4.4.3.6)$$

$$\Lambda = \frac{\sinh^2 \left[\frac{\pi (s / 2.0 + w + d)}{2.0 h} \right]}{2} \quad (4.4.3.7)$$

$$t_c = \sinh^2 \left[\frac{\pi (s / 2.0 + w)}{2.0 h} \right] - \Lambda \quad (4.4.3.8)$$

$$t_B = \sinh^2 \left(\frac{\pi s}{4.0 h} \right) - \Lambda \quad (4.4.3.9)$$

$$\Lambda' = \frac{\cosh^2 \left[\frac{\pi (s / 2.0 + w + d)}{2.0 h} \right]}{2.0} \quad (4.4.3.10)$$

$$\epsilon'_c = \sinh^2 \left[\frac{\pi (s / 2.0 + w)}{2.0 h} \right] - \Lambda' + 1.0 \quad (4.4.3.11)$$

$$\epsilon'_B = \sinh^2 \left[\frac{\pi s}{4.0 h} \right] - \Lambda' + 1.0 \quad (4.4.3.12)$$

To guarantee coplanar propagation,

$$s + 2.0 w + 2.0 d \leq h.$$

The analysis is a conformal mapping technique, so the equations are exact for low frequency calculation.

REFERENCES

- [1] Chang, Ching Ten, and Graham A. Garcia, "Crosstalk Between Two Coplanar Waveguides," *Archiv Für Elektronik und Übertragungstechnik*, Band 43, Heft 1, 1989, pp. 55–58. (Edge-coupled CPW's separated by ground plane.)
- [2] Hanna, Victor Fouad, "Parameters of Coplanar Directional Couplers with Lower Ground Plane," *15th European Microwave Conference Proceedings*, 1985, pp. 820–825. ([2] dropped primes in (8), corrected here.)

4.5 COUPLED MICROSTRIP LINES

4.5.1 Edge-Coupled Microstrip Lines

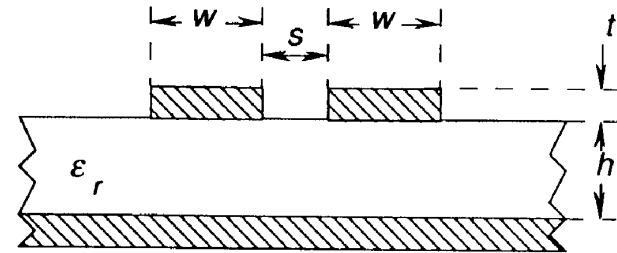


Figure 4.5.1.1: Coupled Microstrip Line

Kirschning and Jansen [8] analyzed this structure with a rigorous spectral-domain hybrid mode calculation. The results of this were then fit numerically. These equations are:

$$Z_{0,e}(0) = \frac{Z_0(0) \sqrt{\frac{\epsilon_{eff}(0)}{\epsilon_{eff,e}(0)}}}{1.0 - \frac{Z_0(0)}{\eta_0} \sqrt{\epsilon_{eff}(0)} Q_4} \quad (4.5.1.1)$$

$$Z_{0,o}(0) = \frac{Z_0(0) \sqrt{\frac{\epsilon_{eff}(0)}{\epsilon_{eff,o}(0)}}}{1.0 - \frac{Z_0(0)}{\eta_0} \sqrt{\epsilon_{eff}(0)} Q_{10}} \quad (4.5.1.2)$$

where $Z_0(0)$ is the zero-frequency impedance of a single, isolated strip (3.5.1.1) and

$$Q_1 = 0.8695 u^{0.194} \quad (4.5.1.3)$$

$$Q_2 = 1.0 + 0.7519 g + 0.189 g^{2.31} \quad (4.5.1.4)$$

$$Q_3 = 0.1975 + [16.6 + (8.4 / g)^{6.0}]^{-0.387} + \frac{\ln \left[\frac{g^{10}}{1.0 + (g / 3.4)^{10}} \right]}{241} \quad (4.5.1.5)$$

$$Q_4 = \frac{2 Q_1}{e^{-8} u^{Q_3} + (2.0 - e^{-8}) u^{-Q_3}} \quad (4.5.1.6)$$

$$Q_5 = 1.794 + 1.14 \ln \left(1.0 + \frac{0.638}{g + 0.517 g^{2.43}} \right) \quad (4.5.1.7)$$

$$Q_6 = 0.2305 + \frac{\ln \left[\frac{g^{10}}{1.0 + (g / 5.8)^{10}} \right]}{281.3} + \frac{\ln (1.0 + 0.598 g^{1.154})}{5.1} \quad (4.5.1.8)$$

$$Q_7 = \frac{10.0 + 190.0 g^2}{1.0 + 82.3 g^3} \quad (4.5.1.9)$$

$$Q_8 = e^{[-6.5 - 0.95 \ln(g) - (g / 0.15)^5]} \quad (4.5.1.10)$$

$$Q_9 = \ln(Q_7) (Q_8 + 1 / 16.5) \quad (4.5.1.11)$$

$$Q_{10} = \frac{Q_2 Q_4 - Q_5 e^{[\ln(u) Q_6 u^{-Q_9}]}}{Q_2} \quad (4.5.1.12)$$

The frequency-dependent even and odd mode impedances are calculated from the above with:

$$Z_{0,e}(f_n) = \frac{Z_e(0) \left\{ 0.9408 [\epsilon_{eff}(f_n)]^{C_e} - 0.9603 \right\} Q_0}{\left\{ (0.9408 - d_e) [\epsilon_{eff}(0)]^{C_e} - 0.9603 \right\} Q_0} \quad (4.5.1.13)$$

$$Z_{0,o}(f_n) = Z_0(f_n) + \frac{Z_{0,o}(0) \left[\frac{\epsilon_{eff,o}(f_n)}{\epsilon_{eff,o}(0)} \right]^{Q_{22}} - Z_0(f_n) Q_{23}}{1.0 + Q_{24} + (0.46 g)^{2.2} Q_{25}} \quad (4.5.1.14)$$

where

$$C_e = 1.0 + 1.275 \left[1.0 - e^{-0.004625 p_e \epsilon_r^{1.674} \left(\frac{f_n}{18.365} \right)^{2.745}} \right] - Q_{12} + Q_{16} - Q_{17} + Q_{18} + Q_{20} \quad (4.5.1.15)$$

$$d_e = 5.086 q_e \frac{r_e}{0.3838 + 0.386 q_e} \frac{e^{-22.2 u^{1.92}}}{1.0 + 1.2992 r_e} \frac{(\epsilon_r - 1.0)^6}{1.0 + 10.0 (\epsilon_r - 1.0)^6} \quad (4.5.1.16)$$

$$p_e = 4.766 e^{-3.228 u^{0.641}} \quad (4.5.1.17)$$

$$q_e = 0.016 + (0.0514 \epsilon_r Q_{21})^{4.524} \quad (4.5.1.18)$$

$$r_e = \left(\frac{f_n}{28.843} \right)^{12} \quad (4.5.1.19)$$

$$Q_{11} = 0.893 \left[1.0 - \frac{0.3}{1.0 + 0.7(\epsilon_r - 1.0)} \right] \quad (4.5.1.20)$$

$$Q_{12} = 2.121 \frac{(f_n / 20)^{4.91}}{[1.0 + Q_{11} (f_n / 20)^{4.91}]} e^{-2.87 g g^{0.902}} \quad (4.5.1.21)$$

$$Q_{13} = 1.0 + 0.038 (\epsilon_r / 8)^{5.1} \quad (4.5.1.22)$$

$$Q_{14} = 1.0 + \frac{1.203 (\epsilon_r / 15)^4}{1.0 + (\epsilon_r / 15)^4} \quad (4.5.1.23)$$

$$Q_{15} = \frac{1.887 e^{-1.5 g^{0.84}} g^{Q_{14}}}{1.0 + \frac{0.41 (f_n / 15)^3 u^{(2 / Q_{13})}}{0.125 + u^{(1.626 / Q_{13})}}} \quad (4.5.1.24)$$

$$Q_{16} = Q_{15} \left[1.0 + \frac{9.0}{1.0 + 0.403 (\epsilon_r - 1.0)^2} \right] \quad (4.5.1.25)$$

$$Q_{17} = 0.394 \left[1.0 - e^{-1.47 (u / 7)^{0.672}} \right] \left[1.0 - e^{-4.25 (f_n / 20)^{1.87}} \right] \quad (4.5.1.26)$$

$$Q_{18} = 0.61 \frac{1.0 - e^{-2.13 (u / 8)^{1.593}}}{1.0 + 6.544 g^{4.17}} \quad (4.5.1.27)$$

$$Q_{19} = \frac{0.21 g^4}{(1.0 + 0.18 g^{4.9}) (1.0 + 0.1 u^2) [1.0 + (f_n / 24.0)^3]} \quad (4.5.1.28)$$

$$Q_{20} = Q_{19} \left[0.09 + \frac{1.0}{1.0 + 0.1 (\epsilon_r - 1.0)^{2.7}} \right] \quad (4.5.1.29)$$

$$Q_{21} = \left| 1.0 - \frac{42.54 g^{0.133} e^{-0.812 g} u^{2.5}}{1.0 + 0.033 u^{2.5}} \right| \quad (4.5.1.30)$$

$$Q_{22} = \frac{0.925 (f_n / Q_{26})^{1.536}}{1.0 + 0.3 (f_n / 30)^{1.536}} \quad (4.5.1.31)$$

$$Q_{23} = 1.0 + \frac{0.005 f_n Q_{27}}{[1.0 + 0.812 (f_n / 15.0)^{1.9}] (1.0 + 0.025 u^2)} \quad (4.5.1.32)$$

$$Q_{24} = 2.506 Q_{28} u^{0.894} \frac{\left[\frac{(1.0 + 1.3 u) f_n}{99.25} \right]^{4.29}}{3.575 + u^{0.894}} \quad (4.5.1.33)$$

$$Q_{25} = \left(\frac{0.3 f_n^2}{10.0 + f_n^2} \right) \left[1.0 + \frac{2.333 (\epsilon_r - 1.0)^2}{5.0 + (\epsilon_r - 1.0)^2} \right] \quad (4.5.1.34)$$

$$Q_{26} = 30.0 - \frac{22.2 \left[\frac{(\epsilon_r - 1.0)}{13} \right]^{12}}{1.0 + 3.0 \left[\frac{(\epsilon_r - 1.0)}{13} \right]^{12}} - Q_{29} \quad (4.5.1.35)$$

$$Q_{27} = 0.4 g^{0.84} \left[1.0 + \frac{2.5 (\epsilon_r - 1.0)^{1.5}}{5.0 + (\epsilon_r - 1.0)^{1.5}} \right] \quad (4.5.1.36)$$

$$Q_{28} = \frac{0.149 (\epsilon_r - 1.0)^3}{94.5 + 0.038 (\epsilon_r - 1.0)^3} \quad (4.5.1.37)$$

$$Q_{29} = \frac{15.16}{1.0 + 0.196 (\epsilon_r - 1.0)^2} \quad (4.5.1.38)$$

$$u = w / h \quad (4.5.1.39)$$

$$g = s / h \quad (4.5.1.40)$$

$$f_n = f h \quad (4.5.1.41)$$

where f is the frequency in GHz and h is the substrate thickness in mm. The effective dielectric constants for the even and odd modes are calculated by beginning with the dc ϵ_{eff} for a single strip derived in [5]:

$$\epsilon_{eff}(0) = 0.5 (\epsilon_r + 1.0) + 0.5 (\epsilon_r - 1.0) \left(1.0 + \frac{10.0}{v} \right)^{-a} e^{v b_e(\epsilon_r)} \quad (4.5.1.42)$$

where

$$v = \frac{u (20.0 + g^2)}{10.0 + g^2} + g e^{-g} \quad (4.5.1.43)$$

$$a_e(v) = 1.0 + \frac{\ln \left[\frac{v^4 + (v / 52.0)^2}{v^4 + 0.432} \right]}{49.0} + \frac{\ln [1.0 + (v / 18.1)^3]}{18.7} \quad (4.5.1.44)$$

$$b_e(\epsilon_r) = 0.564 \left(\frac{\epsilon_r - 0.9}{\epsilon_r + 3.0} \right)^{0.053} \quad (4.5.1.45)$$

The even and odd mode effective relative dielectric constants at dc are then calculated:

$$\epsilon_{eff,e}(0) = [0.5 (\epsilon_r + 1.0) + a_0(u, \epsilon_r) - \epsilon_{eff}(0)] e^{-c_0 g^{d_0}} + \epsilon_{eff}(0) \quad (4.5.1.46)$$

$$\epsilon_{eff,o}(0) = 0.5 (\epsilon_r + 1.0) + 0.5 (\epsilon_r - 1.0) \left[1.0 + \left(\frac{10.0}{v} \right)^{-a_e(v)} b_e(\epsilon_r) \right] \quad (4.5.1.47)$$

where

$$a_0(u, \epsilon_r) = 0.7287 [\epsilon_{eff}(0) - 0.5 (\epsilon_r + 1.0)] (1.0 - e^{-0.179 u}) \quad (4.5.1.48)$$

$$b_0(\epsilon_r) = \frac{0.747 \epsilon_r}{0.15 + \epsilon_r} \quad (4.5.1.49)$$

$$c_0 = b_0(\epsilon_r) - [b_0(\epsilon_r) - 0.207] e^{-0.414 u} \quad (4.5.1.50)$$

$$d_0 = 0.593 + 0.694 e^{-0.562 u} \quad (4.5.1.51)$$

These can be corrected for frequency-dependent dispersion with:

$$\epsilon_{eff,o}(f_n) = \epsilon_r - \frac{\epsilon_r - \epsilon_{eff,o}(0)}{1.0 + F_o(f_n)} \quad (4.5.1.52)$$

$$\epsilon_{eff,e}(f_n) = \epsilon_r - \frac{\epsilon_r - \epsilon_{eff,e}(0)}{1.0 + F_e(f_n)} \quad (4.5.1.53)$$

where

$$F_o(f_n) = P_1 P_2 [(P_3 P_4 + 0.1844) f_n P_{15}]^{1.5763} \quad (4.5.1.54)$$

$$F_e(f_n) = P_1 P_2 [(P_3 P_4 + 0.1844 P_7) f_n]^{1.5763} \quad (4.5.1.55)$$

$$P_1 = 0.27488 + \left[0.6315 + \frac{0.525}{(1.0 + 0.0157 f_n)^{20}} \right] u - 0.065683 e^{-8.7513 u} \quad (4.5.1.56)$$

$$P_2 = 0.33622 (1.0 - e^{-0.03442 \epsilon_r}) \quad (4.5.1.57)$$

$$P_3 = 0.0363 e^{-4.6 u} \left[1 - e^{-(f_n / 38.7)^{4.97}} \right] \quad (4.5.1.58)$$

$$P_4 = 1.0 + 2.751 \left[1.0 - e^{(-\epsilon_r / 15.916)^8} \right] \quad (4.5.1.59)$$

$$P_5 = 0.334 e^{-3.3(\epsilon_r / 15)^3} + 0.746 \quad (4.5.1.60)$$

$$P_6 = P_5 e^{-(f_n / 18)^{0.368}} \quad (4.5.1.61)$$

$$P_7 = 1.0 + 4.069 P_6 g^{0.479} e^{(-1.347 g^{0.595} - 0.17 g^{2.5})} \quad (4.5.1.62)$$

$$P_8 = 0.7168 \left[1.0 + \frac{1.076}{1.0 + 0.0576 (\epsilon_r - 1.0)} \right] \quad (4.5.1.63)$$

$$P_9 = P_8 - 0.7913 \left[1.0 - e^{-(f_n / 20)^{1.424}} \right] \tan^{-1} [2.481 (\epsilon_r / 8)^{0.946}] \quad (4.5.1.64)$$

$$P_{10} = 0.242 (\epsilon_r - 1.0)^{0.55} \quad (4.5.1.65)$$

$$P_{11} = 0.6366 (e^{-3.401 f_n} - 1.0) \tan^{-1} [1.263 (u / 3)^{1.629}] \quad (4.5.1.66)$$

$$P_{12} = P_9 + \frac{1.0 - P_9}{1.0 + 1.183 u^{1.376}} \quad (4.5.1.67)$$

$$P_{13} = \frac{1.695 P_{10}}{0.414 + 1.605 P_{10}} \quad (4.5.1.68)$$

$$P_{14} = 0.8928 + 0.10722 \left[1.0 - e^{-0.42(f_n / 20.0)^{3.215}} \right] \quad (4.5.1.69)$$

$$P_{15} = \left| 1.0 - \frac{0.8928 (1.0 + P_{11}) P_{12} e^{-P_{13} g^{1.092}}}{P_{14}} \right| \quad (4.5.1.70)$$

Equations are valid for:

$$0.1 \leq w / h \leq 10.0$$

$$0.1 \leq s / h \leq 10.0$$

$$1.0 \leq \epsilon_r \leq 18$$

and the stated accuracy for the range $\epsilon_r \leq 12.9$ and $f_n \leq 15$ is better than 1.5%.

The above equations are lengthy; they can be simplified to their dispersionless form in many applications. PCB dimensions are not usually practical for less than about 6 dB of coupling (see the section on Lange couplers for tighter couplings). Other coupling structures (*i.e.*, branch-line hybrids, Lange couplers, Wilkinson dividers, and directional couplers) can be used when tighter or wideband coupling is required.

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4.5.2 Asymmetric Coupled Microstrip Line

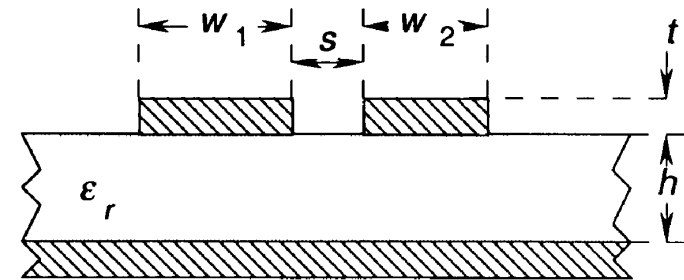


Figure 4.5.2.1: Asymmetric Coupled Microstrip Lines

This configuration has received a great deal of attention for its potential applications in filters, etc. Current designs are based on symmetric couplers and the asymmetry would give us an additional degree of freedom in design. Asymmetric couplers allow the coupling of power at the same time as an impedance transformation takes place.

Unfortunately, closed-form equations for microstrip line do not exist. The reader with sufficient skill may be able to adapt the curves or analyses presented in the references. EM programs may also be used.

The general relationships for coupling and impedances in asymmetric couplers can be defined [1]

$$Z_{0,e}^a = \frac{\sqrt{1.0 - k^2}}{G_a - k \sqrt{G_a G_b}} \quad (4.5.2.1)$$

$$Z_{0,o}^a = \frac{\sqrt{1.0 - k^2}}{G_a - k \sqrt{G_a G_b}} \quad (4.5.2.2)$$

$$Z_{0,e}^b = \frac{\sqrt{1.0 - k^2}}{G_a - k \sqrt{G_a G_b}} \quad (4.5.2.3)$$

$$Z_{0,o}^b = \frac{\sqrt{1.0 - k^2}}{G_a - k \sqrt{G_a G_b}} \quad (4.5.2.4)$$

$$k = \frac{Y_{0,o}^a - Y_{0,e}^a}{\sqrt{(Y_{0,o}^a + Y_{0,e}^a)(Y_{0,o}^b - Y_{0,e}^b)}} \quad (4.5.2.5)$$

where the superscript a corresponds to the impedance of line 2 measured at port 1 and b corresponds to measurements at port 2. For a matched condition to exist the ports are terminated with R_a and R_b such that

$$G_a G_b = \frac{1.0}{4.0} \left[(Y_{0,o}^a + Y_{0,e}^a)(Y_{0,o}^b + Y_{0,e}^b) - (Y_{0,o}^a + Y_{0,e}^a)^2 \right] \quad (4.5.2.6)$$

$$\frac{R_b}{R_a} = \frac{Y_{0,o}^a + Y_{0,e}^a}{Y_{0,o}^b + Y_{0,e}^b} \quad (4.5.2.7)$$

The limitations on impedance transformation for a given coupling are set by the following equations

$$k \leq \sqrt{\frac{R_a}{R_b}} \quad \text{and} \quad k \leq \sqrt{\frac{R_b}{R_a}} \quad (4.5.2.8)$$

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4.5.3 Covered Coupled Microstrip Line

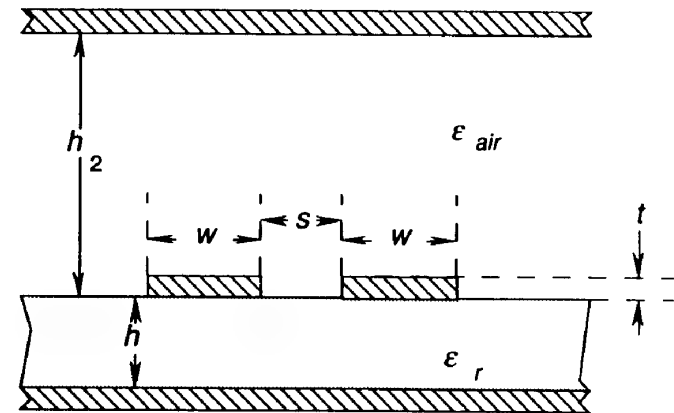


Figure 4.5.3.1: Covered Coupled Microstrip Line

March [3, 4] proposes a technique that corrects and solves a number of problems with previous solutions. The equations of Hammerstad and Jensen [2] are modified for the presence of the cover by a series of corrections. Results were verified by comparison with MSTRIP simulations. The equations of Hammerstad and Jensen are:

$$Z_{0,e}(u, g) = \frac{Z_0(u)}{1.0 - Z_0(u) \Phi_e(u, g) / \eta_0} \quad (\Omega) \quad (4.5.3.1)$$

$$Z_{0,o}(u, g) = \frac{Z_0(u)}{1.0 - Z_0(u) \Phi_o(u, g) / \eta_0} \quad (\Omega) \quad (4.5.3.2)$$

$$g = s / h \quad (4.5.3.3)$$

$$u = w / h \quad (4.5.3.4)$$

$$\epsilon_{eff,e}(u, g, \epsilon_r) = \frac{\epsilon_r + 1.0}{2.0} + F_e(u, g, \epsilon_r) \frac{\epsilon_r - 1.0}{2.0} \quad (4.5.3.5)$$

$$\epsilon_{eff,o}(u, g, \epsilon_r) = \frac{\epsilon_r + 1.0}{2.0} + F_o(u, g, \epsilon_r) \frac{\epsilon_r - 1.0}{2.0} \quad (4.5.3.6)$$

$$F_e(u, g, \epsilon_r) = \left[1.0 + \frac{10.0}{\mu(u, g)} \right]^{-a(u) b(\epsilon_r)} \quad (4.5.3.7)$$

$$F_o(u, g, \epsilon_r) = f_0(u, g, \epsilon_r) \left[1.0 + \frac{10.0}{\mu(u, g)} \right]^{-a(u) b(\epsilon_r)} \quad (4.5.3.8)$$

$$\Phi_e(u, g) = \frac{\varphi(u)}{(\psi(g) \{ \alpha(g) u^{m(g)} + [1.0 - \alpha(g)] u^{-m(g)} \})} \quad (4.5.3.9)$$

$$\varphi(u) = 0.8645 u^{0.172} \quad (4.5.3.10)$$

$$\psi(g) = 1.0 + g / 1.45 + \frac{g^{2.09}}{3.95} \quad (4.5.3.11)$$

$$\alpha(g) = 0.5 e^{-g} \quad (4.5.3.12)$$

$$m(g) = 0.2175 + \left[4.113 + \left(\frac{20.36}{g} \right)^6 \right]^{-0.251} + \left(\frac{1.0}{323.0} \right) \ln \left[\frac{g^{10}}{1.0 + \left(\frac{g}{13.8} \right)^{10}} \right] \quad (4.5.3.13)$$

$$\phi_o(u, g) = \phi_o(u, g) - \frac{\theta(g)}{\psi(g)} e^{\beta(g) u^{-n(g)} \ln u} \quad (4.5.3.14)$$

$$\theta(g) = 1.729 + 1.175 \ln \left(1.0 + \frac{0.627}{g + 0.327 g^{2.17}} \right) \quad (4.5.3.15)$$

$$\beta(g) = 0.2306 + \frac{1.0}{301.8} \ln \left[\frac{g^{10}}{1.0 + \left(\frac{g}{3.73} \right)^{10}} \right] + \frac{1.0}{5.3} \ln (1.0 + 0.646 g^{1.175}) \quad (4.5.3.16)$$

$$n(g) = \left\{ \frac{1.0}{17.7} + e^{\left[-6.424 - 0.76 \ln g - \left(\frac{g}{0.23} \right)^5 \right]} \right\} \ln \left(\frac{10.0 + 68.3 g^2}{1.0 + 32.5 g^{3.093}} \right) \quad (4.5.3.17)$$

$$\mu(u, g) = g e^{-g} + u \frac{20.0 + g^2}{10.0 + g^2} \quad (4.5.3.18)$$

$$f_0(u, g, \epsilon_r) = f_{0,1}(g, \epsilon_r) e^{p(g) \ln u + q(g) \sin \left(\frac{\pi \ln u}{\ln 10.0} \right)} \quad (4.5.3.19)$$

$$p(g) = \frac{e^{-0.745 g^{0.295}}}{\cosh(g^{0.15})} \quad (4.5.3.20)$$

$$q(g) = e^{-1.366 - g} \quad (4.5.3.21)$$

$$f_{0,1}(g, \epsilon_r) = 1.0 - e^{\left\{ -0.179 g^{0.15} - \frac{0.328 g^{r(g, \epsilon_r)}}{\ln [e^{1.0} + (g / 7.0)^{2.8}]} \right\}} \quad (4.5.3.22)$$

$$r(g, \epsilon_r) = 1.0 + 0.15 \left\{ 1.0 - e^{\left[\frac{1.0 - (\epsilon_r - 1.0)^{2.0}}{8.2} \right]} \right\} \frac{1.0 + g^{-6.0}}{1.0 + g^{-6.0}} \quad (4.5.3.23)$$

The values of the modal impedances and effective dielectric constants agree with MSTRIP to better than $\pm 0.5\%$ for

$$0.1 \leq w/h \leq 10.0$$

$$0.01 \leq s/h$$

$$1.0 \leq \epsilon_r \leq 128.0$$

March corrects for the finite trace thickness by replacing w in the above equations with equivalent widths w_e' and w_o' for the even and odd modes, respectively.

$$w_e' = w + \Delta w_e \quad (4.5.3.24)$$

$$s_e' = s - \Delta w_e \quad (4.5.3.25)$$

$$\Delta w_e = \Delta w \left[1.0 - 0.5 e^{(-0.69 \Delta w / \Delta t)^2} \right] \quad (4.5.3.26)$$

$$w_o' = w + \Delta w_o \quad (4.5.3.27)$$

$$s_o' = s - \Delta w_o \quad (4.5.3.28)$$

$$\Delta w_o = \Delta w_e + \Delta t \quad (4.5.3.29)$$

$$\Delta t = \frac{2.0 \ t \ h}{s (\epsilon_r + 1.0)} \quad (4.5.3.30)$$

$$\Delta w = \frac{t}{\pi} \left\{ 1.0 + \ln 4.0 - 0.5 \ln \left[\left(\frac{t}{h} \right)^2 + \left(\frac{t}{\pi w} \right)^2 \right] \right\} \quad (4.5.3.31)$$

Finite trace thickness corrections are valid for $s \geq 2.0 t$. The equations have been verified for $0.1 \leq u \leq 10.0$ and $g \geq 0.01$.

To correct for the presence of the cover when $h_2 / h \leq 39.0$

$$q_{e,c} = \tanh \left(1.626 + 0.107 \ h_2 / h - \frac{1.733}{\sqrt{h_2 / h}} \right) \quad (4.5.3.32)$$

and when $h_2 / h > 39.0$

$$q_{e,c} = 1.0 \quad (4.5.3.33)$$

and the odd mode is corrected for the cover when $h_2 / h \geq 7.0$ with

$$q_{o,c} = \tanh \left[\frac{9.575}{7.0 - h_2 / h} - 2.965 + 1.68 \ h_2 / h - 0.311 \ (h_2 / h)^2 \right] \quad (4.5.3.34)$$

and for $h_2 / h < 7.0$

$$q_{o,c} = 1.0 \quad (4.5.3.35)$$

These equations are accurate to $\pm 0.4\%$ and typically to $\pm 0.2\%$. Even mode corrections for the cover are

$$\Delta Z_{0,e}(\epsilon_r = 1.0) = f_e(w / h, h_2 / h) g_e(s / h, h_2 / h) \quad (4.5.3.36)$$

$$f_e(w / h, h_2 / h) = 1.0 - \tanh^{-1} \left[A + \left(B + \frac{C}{h} \right) \frac{w}{h} \right] \quad (4.5.3.37)$$

$$A = \frac{-4.351}{(1.0 + h_2 / h)^{1.842}} \quad (4.5.3.38)$$

$$B = \frac{6.639}{(1.0 + h_2 / h)^{1.861}} \quad (4.5.3.39)$$

$$C = \frac{-2.291}{(1.0 + h_2 / h)^{1.90}} \quad (4.5.3.40)$$

$$g_e(s / h, h_2 / h) = 270.0 \left[1.0 - \tanh \left(D + E \sqrt{1.0 + h_2 / h} - \frac{F}{1.0 + h_2 / h} \right) \right] \quad (4.5.3.41)$$

$$D = 0.747 \csc \left(\frac{\pi x}{2.0} \right) \quad (4.5.3.42)$$

$$E = 0.725 \csc \left(\frac{\pi y}{2.0} \right) \quad (4.5.3.43)$$

$$F = 10^{(0.11 - 0.0947 \ s / h)} \quad (4.5.3.44)$$

$$x = 10^{(0.103 \ s / h - 0.159)} \quad (4.5.3.45)$$

$$y = 10^{(0.0492 \ s / h - 0.073)} \quad (4.5.3.46)$$

and odd mode corrections for the cover are

$$\Delta Z_{0,o}(\epsilon_r = 1.0) = f_o(w / h, h_2 / h) g_o(s / h, h_2 / h) \quad (4.5.3.47)$$

$$f_o(w / h, h_2 / h) = \left(\frac{w}{h} \right)^J \quad (4.5.3.48)$$

$$J = \tanh \left[\frac{(1.0 + h_2 / h)^{1.585}}{6.0} \right] \quad (4.5.3.49)$$

$$g_o(s/h, h_2/h) = 270.0 \left[1.0 - \tanh \left(G + K \sqrt{1.0 + \frac{h_2}{h}} - \frac{L}{1.0 + h_2/h} \right) \right] \quad (4.5.3.50)$$

$$G = 2.7178 - 0.796 \frac{s}{h} \quad (4.5.3.51)$$

for $s/h > 0.858$

$$K = \log \left[20.492 \left(\frac{s}{h} \right)^{0.174} \right] \quad (4.5.3.52)$$

for $s/h \leq 0.858$

$$K = 1.30 \quad (4.5.3.53)$$

for $s/h > 0.873$

$$L = 2.51 \left(\frac{s}{h} \right)^{-0.462} \quad (4.5.3.54)$$

and for $s/h \leq 0.873$

$$L = 2.674 \quad (4.5.3.55)$$

Dispersion can now be accounted for by the remaining relations. First we replace ϵ_{eff} (u, g, ϵ_r) and $\epsilon_{eff,o}$ (u, g, ϵ_r), which are zero-frequency relations [we will denote them $\epsilon_{eff,o}(0)$ and $\epsilon_{eff,e}(0)$] with frequency-dependent relations

$$\epsilon_{eff,o}(f) = \epsilon_r - \frac{\epsilon_r - \epsilon_{eff,o}(0)}{1.0 - G_o(f/f_{p,o})^2} \quad (4.5.3.56)$$

$$\epsilon_{eff,e}(f) = \epsilon_r - \frac{\epsilon_r - \epsilon_{eff,e}(0)}{1.0 - G_e(f/f_{p,e})^2} \quad (4.5.3.57)$$

where

$$G_e = \frac{\pi^2 (\epsilon_r - 1.0)}{12.0 \epsilon_{eff,e}(0)} \sqrt{\frac{Z_{0,e}}{120.0}} \quad (4.5.3.58)$$

$$G_o = \frac{\pi^2 (\epsilon_r - 1.0)}{12.0 \epsilon_{eff,o}(0)} \sqrt{\frac{Z_{0,o}}{120.0}} \quad (4.5.3.59)$$

$$f_{p,e} = \frac{Z_{0,e}}{4.0 \mu_0 H} \quad (4.5.3.60)$$

$$f_{p,o} = \frac{Z_{0,o}}{\mu_0 H} \quad (4.5.3.61)$$

The even and odd mode impedances are similarly corrected. The zero-frequency relations $Z_{0,e}(u, g)$ and $Z_{0,o}(u, g)$ are denoted $Z_{0,e}(0)$ and $Z_{0,o}(0)$. At high frequencies the fields become completely contained in the dielectric (ϵ_{eff} approaches ϵ_r) and the impedance of the line can be related to that of a stripline. Calculate the impedances of edge-coupled striplines (see later in this chapter) with the microstrip line dimensions as follows

$$Z_{0,e,sl}(w, b = 2.0 h, t, s) \quad (4.5.3.62)$$

$$Z_{0,o,sl}(w, b = 2.0 h, t, s) \quad (4.5.3.63)$$

We can now calculate the frequency-corrected impedances:

$$Z_{0,e}(f) = Z_{0,e,sl} - \frac{2.0 Z_{0,e,sl} - Z_{0,e}(0)}{1.0 + G_e(f/f_{p,e})^{1.6}} \quad (4.5.3.64)$$

$$Z_{0,o}(f) = Z_{0,o,sl} - \frac{2.0 Z_{0,o,sl} - Z_{0,o}(0)}{1.0 + G_o(f/f_{p,o})^{1.6}} \quad (4.5.3.65)$$

The accuracy of these corrections is better than $\pm 0.5\%$ to 18 GHz.

Although these equations are quite involved, we are rewarded with a general-purpose set of relations that describe a frequently encountered structure quite accurately. In fact since most microstrip lines are packaged one might use only this relation.

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4.5.4 Broadside Coupled Microstrip Lines

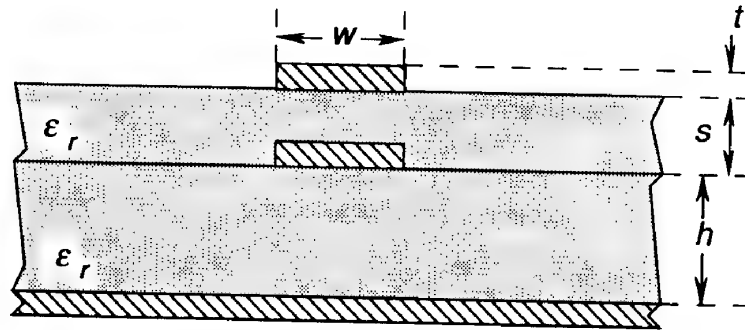


Figure 4.5.4.1: Broadside Coupled Microstrip Lines

The broadside coupled microstrip line structure is useful for tight coupling directional couplers and for differential transmission of signals.

Malone [1] reports good results with design of this structure by equating the capacitance of an equivalent broadside coupled stripline structure. If the desired coupled line parameters are calculated for broadside coupled striplines (Figure 4.6.2.1), the dimensions of the equivalent microstrip line structure are then calculated:

$$h = b_{sl} \left(\frac{1.0 - s_{sl}}{5.0} \right) \quad (4.5.4.1)$$

$$s = s_{sl} \quad (4.5.4.2)$$

$$w = w_{sl} \quad (4.5.4.3)$$

where the dimensions of the stripline are as shown in Figure 4.6.2.1 and are indicated above with the subscript *sl*.

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4.5.5 Three Coupled Microstrip Lines

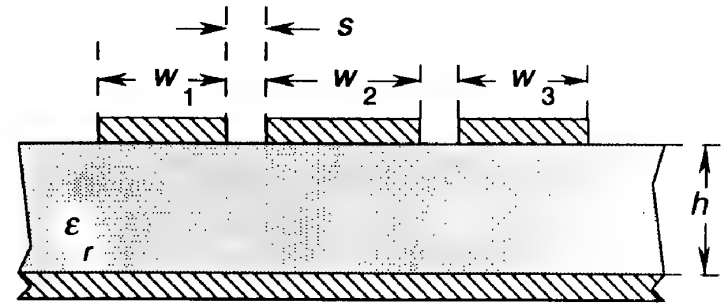


Figure 4.5.5.1: Three Coupled Microstrip Lines

Three coupled lines have the potential to reduce space over two-line couplers where the signal is to be split into three paths. This structure has five mode impedances.

Abdallah and El-Deeb [1] present the relation of the widths for equal mode impedances in terms of a polynomial fit to numerical data. If $w_1 = w_3$, then

$$w_2 / w_1 = \sum_{n=0}^3 B_n \left(\frac{s}{h} \right)^n \quad (4.5.5.1)$$

where

$$B_0 = 2.40599 - 3.06631 (w_1 / h) + 3.46389 (w_1 / h)^2 - 1.66833 (w_1 / h)^3 \quad (4.5.5.2)$$

$$B_1 = -1.41242 + 0.75437 (w_1 / h) + 1.98412 (w_1 / h)^2 - 1.49878 (w_1 / h)^3 \quad (4.5.5.3)$$

$$B_2 = 0.61033 - 0.38901 (w_1 / h) - 0.17853 (w_1 / h)^2 - 0.42474 (w_1 / h)^3 \quad (4.5.5.4)$$

$$B_3 = -0.20601 + 1.04401 (w_1 / h) - 2.23510 (w_1 / h)^2 + 1.74802 (w_1 / h)^3 \quad (4.5.5.5)$$

which is valid for $s/h \leq 0.8$ and $w_1/h \leq 0.8$ and accurate to better than 1%. For lines with $w_1/h > 0.8$, equal width lines ($w_1 = w_2$) create equal mode impedances.

Perlow and Presser [3] form a two-port from this structure by tying the outer lines' ends together. The resulting coupler can achieve tighter coupling than two-lines. Equations are derived by equating the capacitances of this structure to those of the two line coupler.

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4.5.6 Microstrip Line-to-Stripline Aperture Coupler

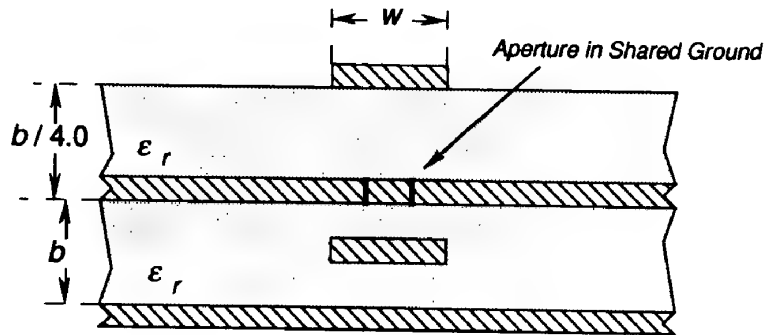


Figure 4.5.6.1: Microstripline-to-Stripline Aperture Coupler

The aperture coupler allows us to couple between layers of a multilayer structure. This is useful when we need a small amount of signal coupled to another layer. The equations may also be useful when evaluating coupling through areas of discontinuous ground plane.

Even and odd mode impedances for this structure are not found in the literature. We do find relations for the coupling, k . For a circular-shaped aperture in the shared ground plane,

$$k_{dB} = 20.0 \log_{10} \left[\frac{5.0 d_0^3 \left(\frac{-\epsilon_r \epsilon_{eff}}{\epsilon_r + \epsilon_{eff}} + 2.0 \sqrt{\epsilon_r \epsilon_{eff}} \right)}{9.0 \lambda b^2 \sqrt{F(m)}} \right] \quad (4.5.6.1)$$

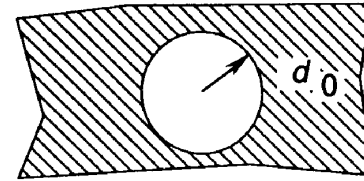


Figure 4.5.6.2: Circular Aperture Dimensions

For a slot aperture,

$$k_{dB} = 20.0 \log_{10} \left[\frac{20.0 \left(\frac{\alpha_c \epsilon_r \epsilon_{eff}}{\epsilon_r + \epsilon_{eff}} + \alpha_m \sqrt{\epsilon_r \epsilon_{eff}} \right)}{3.0 \lambda b^2 \sqrt{F(m)}} \right] \quad (4.5.6.2)$$

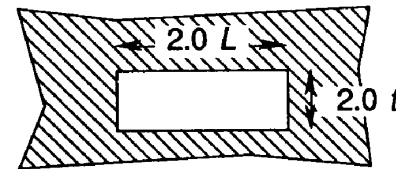


Figure 4.5.6.3: Rectangular Aperture Dimensions

where

$$\alpha_m = \frac{1.054 \beta^3}{\ln \left(1.0 + \frac{0.66}{\alpha} \right)} \quad (4.5.6.3)$$

$$\alpha_c = \frac{\pi \alpha^2 \beta^3 (1.0 - 0.5663 \alpha + 0.1398 \alpha^2)}{2.0} \quad (4.5.6.4)$$

$$\alpha = \frac{t}{l}$$

and for both of the equations:

$$F(m) = \frac{4.0}{b^2} \int_{-b/2}^{b/2} \int_{-\infty}^{\infty} \frac{dx dy}{\left| 1.0 - m^2 \cosh^2 \left[\frac{\pi (x + jy)}{b} \right] \right|} \quad (4.5.6.5, 4.5.6.6)$$

and

$$m = \operatorname{sech} \left(\frac{\pi w}{2.0 b} \right) \quad (4.5.6.7)$$

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4.6 COUPLED STRIPLINES

4.6.1 Nonhomogeneous Dielectric Broadside Coupled Striplines

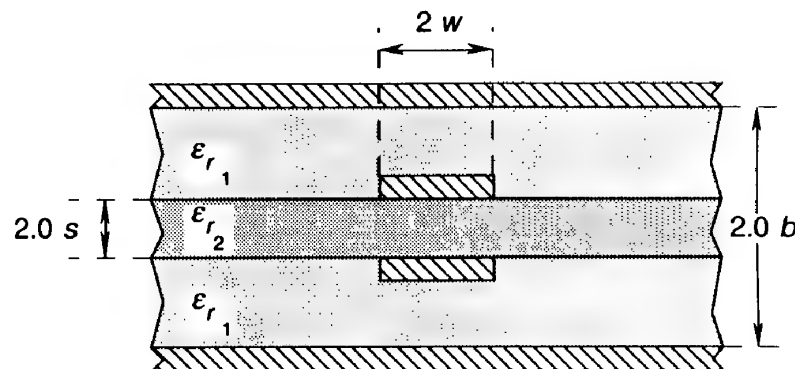


Figure 4.6.1.1: Nonhomogeneous Dielectric Broadside Coupled Striplines

This configuration is described in [2] and [5]. Broadside coupling allows tight coupling to be achieved with practical PCB dimensions. Bahl and Bhartia [2] treat this topology as two separate cases. The equations are in terms of the structure's capacitance:

$$Z_{0,e} = \frac{1.0}{c_0 \sqrt{C_e C_e^a}} \quad (4.6.1.1)$$

$$Z_{0,o} = \frac{1.0}{c_0 \sqrt{C_o C_o^a}} \quad (4.6.1.2)$$

$$\epsilon_{eff,e} = \frac{C_e}{C_e^a} \quad (4.6.1.3)$$

$$\epsilon_{eff,o} = \frac{C_o}{C_o^a} \quad (4.6.1.4)$$

where, for $\epsilon_{r,2} = 1.0$

$$\frac{1.0}{C_e} = \frac{1.0}{\pi \epsilon_0 Q^2} \int_0^\infty \frac{h_e \tilde{f}^2(\beta) d(\beta b)}{\left\{ \tanh(\beta s) + \epsilon_r \coth[\beta(b-s)] \right\} (\beta b)} \quad (4.6.1.5)$$

$$\frac{1.0}{C_o} = \frac{1.0}{\pi \epsilon_0 Q^2} \int_0^\infty \frac{h_o \tilde{f}^2(\beta) d(\beta b)}{\left\{ [\tanh(\beta s)]^{-1} + \epsilon_r \coth[\beta(b-s)] \right\} (\beta b)} \quad (4.6.1.6)$$

where

$$h_e = \frac{1.0}{2.0} \left\{ 1.0 + \frac{\cosh[\beta(s-t)]}{\cosh(\beta s)} \right\} \quad (4.6.1.7)$$

$$h_o = \frac{1.0}{2.0} \left\{ 1.0 + \frac{\sinh[\beta(s-t)]}{\sinh(\beta s)} \right\} \quad (4.6.1.8)$$

For $\epsilon_{r,1} = 1.0$,

$$\frac{1.0}{C_e} = \frac{1.0}{\pi \epsilon_0 Q^2} \int_0^\infty \frac{g_e \tilde{f}^2(\beta) d(\beta b)}{\left\{ \epsilon_r [\tanh(\beta s)] + \coth[\beta(b-s)] \right\} (\beta b)} \quad (4.6.1.9)$$

$$\frac{1.0}{C_o} = \frac{1.0}{\pi \epsilon_0 Q^2} \int_0^\infty \frac{g_o \tilde{f}^2(\beta) d(\beta b)}{\left\{ \epsilon_r [\tanh(\beta s)]^{-1} + \coth[\beta(b-s)] \right\} (\beta b)} \quad (4.6.1.10)$$

where

$$\begin{aligned} \frac{\tilde{f}(\beta)}{Q} &= 1.6 \left(\frac{\sin \beta w}{\beta w} \right) \\ &+ \frac{2.4}{(\beta w)^2} \left[\cos \beta w - \frac{2.0 \sin \beta w}{\beta w} + \frac{\sin^2(\beta w / 2.0)}{(\beta w / 2.0)^2} \right] \end{aligned} \quad (4.6.1.11)$$

$$g_e = \frac{1.0}{2.0} \left\{ 1.0 + \frac{\cosh[\beta(b-s-t)]}{\cosh[\beta(b-s)]} \right\} \quad (4.6.1.12)$$

$$g_o = \frac{1.0}{2.0} \left\{ 1.0 + \frac{\sinh[\beta(b-s-t)]}{\sinh[\beta(b-s)]} \right\} \quad (4.6.1.13)$$

and C_o^a and C_e^a are the odd and even mode capacitances with air replacing the dielectric in the equations. Integrals are evaluated numerically.

Conductor losses are calculated with the incremental inductance rule:

$$\begin{aligned} \alpha_{c,e} &= \frac{0.0116 R_s \sqrt{\epsilon_{eff,e}}}{Z_{0,e}^a} \left[\frac{\partial Z_{0,e}^a}{\partial(s/b)} (1.0 - s/b) \right. \\ &\quad \left. - \frac{\partial Z_{0,e}^a}{\partial(w/b)} (1.0 + w/b) - \frac{\partial Z_{0,e}^a}{\partial(t/b)} (0.5 + t/b) \right] \quad (\text{dB/length}) \end{aligned} \quad (4.6.1.14)$$

$$\begin{aligned} \alpha_{c,o} &= \frac{0.0116 R_s \sqrt{\epsilon_{eff,o}}}{Z_{0,o}^a} \left[\frac{\partial Z_{0,o}^a}{\partial(s/b)} (1.0 - s/b) \right. \\ &\quad \left. - \frac{\partial Z_{0,o}^a}{\partial(w/b)} (1.0 + w/b) - \frac{\partial Z_{0,o}^a}{\partial(t/b)} (0.5 + t/b) \right] \quad (\text{dB/length}) \end{aligned} \quad (4.6.1.15)$$

Dielectric losses are calculated with

$$\alpha_{d,e} = 27.3 \frac{\epsilon_r}{\sqrt{\epsilon_{eff,e}}} \frac{\epsilon_{eff,e} - 1.0 \tan \delta}{\epsilon_r - 1.0} \frac{\delta}{\lambda_0} \quad (4.6.1.16)$$

$$\alpha_{d,o} = 27.3 \frac{\epsilon_r}{\sqrt{\epsilon_{eff,o}}} \frac{\epsilon_{eff,o} - 1.0 \tan \delta}{\epsilon_r - 1.0} \frac{\delta}{\lambda_0} \quad (4.6.1.17)$$

The effect of an open end as a discontinuity in this structure is reported in [7]. See Chapter 5 for more information.

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4.6.2 Broadside Coupled Striplines

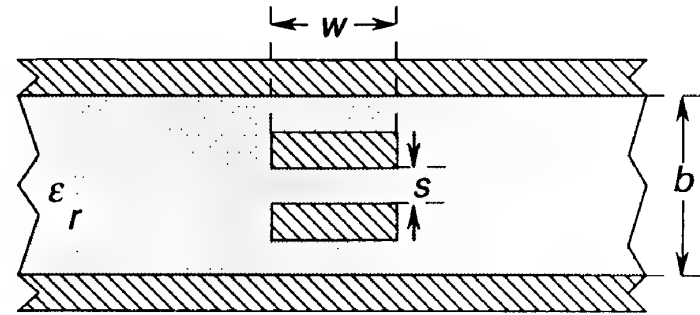


Figure 4.6.2.1: Broadside Coupled Striplines

This structure is described in [7] and [5]; however, broader range equations are available in Bahl and Bhartia [1]:

$$Z_{0,e} = \frac{\eta_0}{\sqrt{\epsilon_r}} \frac{K'(k)}{K(k)} \quad (\Omega) \quad (4.6.2.1)$$

$$Z_{0,o} = \frac{296.1}{\sqrt{\epsilon_r} \frac{b}{s} \tanh^{-1}(k)} \quad (\Omega) \quad (4.6.2.2)$$

where k is the solution of

$$\frac{w}{b} = \frac{1.0}{\pi} \left[\ln \left(\frac{1.0 + R}{1.0 - R} \right) - \frac{s}{b} \ln \left(\frac{1.0 + R/k}{1.0 - R/k} \right) \right] \quad (4.6.2.3)$$

and

$$R = \sqrt{\frac{\left(\frac{k}{s} \frac{b}{s} - 1.0 \right)}{\left(\frac{b}{k} \frac{s}{s} - 1.0 \right)}} \quad (4.6.2.4)$$

The accuracy is "virtually exact" for $w/s \geq 0.35$. These equations are more exact for high dielectric constant materials. Bahl and Bhartia [1] also do a sensitivity analysis of this structure and show that the sensitivity to dimensional changes increases with coupling.

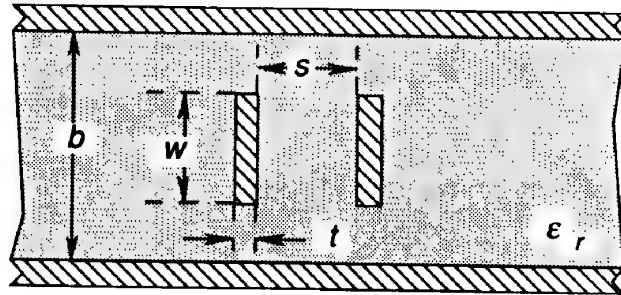


Figure 4.6.2.2: Vertical Plate Broadside Coupled Striplines

Another broadside coupled stripline configuration is described in Cohn, [2] and shown in the figure above. This configuration is less often encountered due to the difficulty of fabricating it with standard photolithographic techniques. The relevant design equations are

$$Z_{0,e} = \frac{\eta_0}{2.0 \sqrt{\epsilon_r}} \frac{K(k)}{K(k')} \quad (4.6.2.5)$$

$$Z_{0,o} = \frac{295.9}{\sqrt{\epsilon_r}} \frac{1.0}{\frac{b}{s} \cos^{-1} k + \ln \frac{1.0}{k}} \quad (4.6.2.6)$$

where

$$\frac{w}{b} = \frac{2.0}{\pi} \left[\frac{k'}{k} \tan^{-1} \left(\sqrt{\frac{1.0 - \frac{k}{k'} \frac{s}{b}}{1.0 + \frac{k}{k'} \frac{s}{b}}} \right) - \frac{s}{b} \tanh^{-1} \left(\sqrt{\frac{1.0 - \frac{k}{k'} \frac{s}{b}}{1.0 + \frac{k}{k'} \frac{s}{b}}} \right) \right] \quad (4.6.2.7)$$

which is accurate for

$$w/s > 1.0$$

and when $s/b \ll 1.0$ they are valid for smaller w/s . For thick strips the effect of t is compensated with equations of [3] by increasing the zero-thickness capacitance:

$$C'_{0,o} \left(\frac{w}{b}, \frac{s}{b}, \frac{t}{s} \right) = C'_{0,o} \left(\frac{w}{b}, \frac{s}{b}, 0.0 \right) + 4.0 \Delta C'_2 \left(\frac{t}{s} \right) \quad (4.6.2.8)$$

for $t \ll s$

$$C'_{0,e} \left(\frac{w}{b}, \frac{s}{b}, \frac{t}{s} \right) = C'_{0,e} \left(\frac{w}{b}, \frac{s}{b}, 0.0 \right) \quad (4.6.2.9)$$

and for $t \gg s$

$$C'_{0,e} \left(\frac{w}{b}, \frac{s}{b}, \frac{t}{s} \right) = C_0 \left(\frac{w'}{b}, \frac{t'}{w'} \right) \quad (4.6.2.10)$$

where

$$s' = s + 2.0 t \quad (4.6.2.11)$$

$$w' = s + 2.0 t \quad (4.6.2.12)$$

$$t' = w \quad (4.6.2.13)$$

$$\Delta C'_2(t/s) = \frac{\epsilon}{2.0 \pi} \left[\left(1.0 + \frac{t}{s} \right) \ln \left(1.0 + \frac{t}{s} \right) - \frac{t}{s} \ln \frac{t}{s} \right] \quad (4.6.2.14)$$

$$C_0 \left(\frac{w'}{b}, \frac{t'}{w'} \right) = \text{capacitance of a single strip of width } w', \text{ thickness } t', \text{ and having ground plane spacing } b. \quad (4.6.2.15)$$

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4.6.3 Offset Coupled Striplines

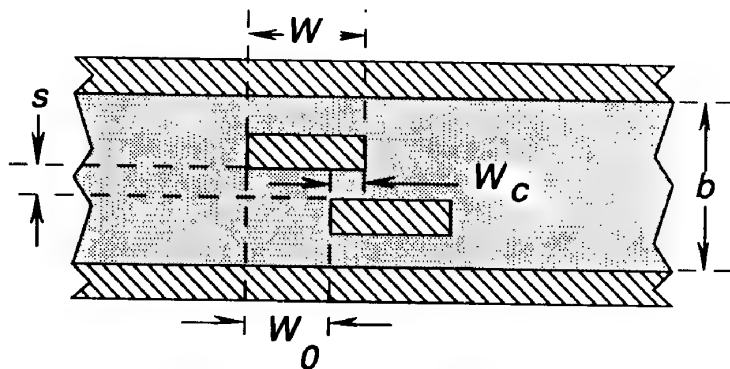


Figure 4.6.3.1: Offset Coupled Striplines

This configuration allows couplings less than approximately 6 dB to be achieved with practical PCB dimensions. For $w_c = 0$ use the equations of the last section. Shelton's [5] equations are normalized to $b = 1.0$.

$$Z_{0,e} = \frac{120.0 \pi}{\sqrt{\epsilon_r}} \frac{1.0}{C_e} \quad (\Omega) \quad (4.6.3.1)$$

$$Z_{0,o} = \frac{120.0 \pi}{\sqrt{\epsilon_r}} \frac{1.0}{C_o} \quad (4.6.3.2)$$

where for tight coupling $w_c \geq 0$ (overlapping traces as shown in Figure 4.6.3.1), $w / (1.0 - s) \geq 0.35$, and $w_c \geq 0.7$:

$$C_o = 2.0 w \left(\frac{1.0}{1.0 - s} + \frac{1.0}{s} \right) + C_{fo} = \frac{2.0 w}{s(1.0 - s)} + C_{fo} \quad (4.6.3.3)$$

$$C_e = \frac{2.0 w}{1.0 - s} + C_{fe} \quad (4.6.3.4)$$

$$C_{fo} = \frac{1.0}{\pi} \left\{ -\frac{2.0}{1.0 - s} \log s + \frac{1.0}{s} \log \left[\frac{p r}{(p + s)(1.0 + p)(r - s)(1.0 - r)} \right] \right\} \quad (4.6.3.5)$$

$$C_{fe} = \frac{1.0}{\pi} \left[\frac{-2.0}{1.0 - s} \log s - 2.0 \log s + 4.0 \log(s + p r) - \log p r (p + s)(p + 1.0)(r - s)(1.0 - r) \right] \quad (4.6.3.6)$$

$$r = \frac{p + \frac{1.0 + s}{2.0}}{1.0 + p \left(\frac{1.0 + s}{2.0 s} \right)} \quad (0 < p < \infty) \quad (4.6.3.7)$$

$$p = \frac{(B - 1.0) \left(\frac{1.0 + s}{2.0} \right) + \sqrt{\left(\frac{1.0 + s}{2.0} \right)^2 (B - 1.0)^2 + 4.0 s B}}{2.0} \quad (4.6.3.8)$$

$$B = \frac{p r}{s} = \frac{A - 2.0 + \sqrt{A^2 - 4.0 A}}{2.0} \quad (4.6.3.9)$$

$$A = \exp \left[\frac{60 \pi^2}{\sqrt{\epsilon_r} Z_0} \left(\frac{1.0 - p s}{\sqrt{\rho}} \right) \right] \quad (4.6.3.10)$$

$$\rho = \frac{Z_{0,e}}{Z_{0,o}} \quad (4.6.3.11)$$

and for loose coupling $w_c \leq 0$ (traces not overlapping), $w / (1.0 - s) \geq 0.35$, and $2.0 w_o / (1.0 + s) \geq 0.85$:

$$C_o = 2.0 w \left(\frac{1.0}{1.0 - s} + \frac{1.0}{1.0 + s} \right) + C_{fo} - C_f(a = \infty) \quad (4.6.3.12)$$

$$= \frac{4.0 w}{1.0 - s^2} + C_{fo} + C_f(a = \infty)$$

$$C_e = \frac{4.0 w}{1.0 - s^2} + C_{fe} + C_f(a = \infty) \quad (4.6.3.13)$$

$$C_{fo} = \frac{2.0}{\pi} \left[\frac{1.0}{1.0 + s} \log \frac{1.0 + a}{a(1.0 - q)} - \frac{1.0}{1.0 - s} \log q \right] \quad (4.6.3.14)$$

$$C_{fe} = \frac{2.0}{\pi(1.0 + s)} \log \left(\frac{1.0 + a}{a - a q} \right) - \frac{2.0}{1.0 - s} \log q - 2.0 \log \left(\frac{1.0 + a q}{a q} \right) \quad (4.6.3.15)$$

$$q = \left(\frac{s + 1.0}{2.0} \right) \left(\frac{a + \frac{2.0 s}{s + 1.0}}{a + \frac{s + 1.0}{2.0}} \right) \quad (4.6.3.16)$$

$$a = \sqrt{\left(\frac{s - k}{s + 1.0} \right)^2 + k} - \frac{s - k}{s + 1.0} \quad (4.6.3.17)$$

$$k = \frac{1.0}{\exp \left(\frac{\pi \Delta C}{2.0} \right) - 1.0} \quad (4.6.3.18)$$

$$\Delta C = \frac{120 \pi (\rho - 1.0)}{\sqrt{\epsilon_r} Z_0 \sqrt{\rho}} \quad (4.6.3.19)$$

$$\rho = \frac{Z_{0,e}}{Z_{0,o}} \quad (4.6.3.20)$$

$$C_f(a = \infty) = -\frac{2.0}{\pi} \left[\frac{1.0}{1.0 + s} \log \left(\frac{1.0 - s}{2.0} \right) + \frac{1.0}{1.0 - s} \log \left(\frac{1.0 + s}{2.0} \right) \right] \quad (4.6.3.21)$$

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4.6.4 Zero-Thickness Edge-Coupled Stripline

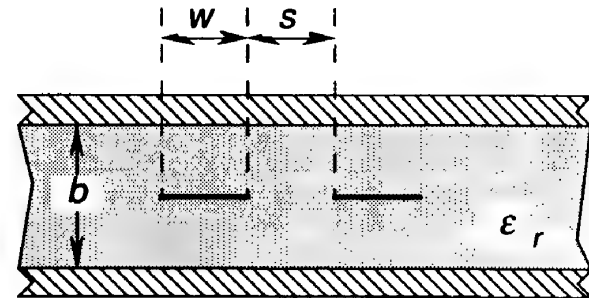


Figure 4.6.4.1: Edge-Coupled Stripline

This zero-thickness structure never occurs in real life; however it is often used to check the accuracy of finite-thickness equations.

$$Z_{0,e} = \frac{30 \pi K(k'_e)}{\sqrt{\epsilon_r} K(k_e)} (\Omega) \quad (4.6.4.1)$$

$$Z_{0,o} = \frac{30 \pi K(k_o)}{\sqrt{\epsilon_r} K(k_o)} (\Omega) \quad (4.6.4.2)$$

where

$$k_e = \tanh\left(\frac{\pi w}{2b}\right) \tanh\left[\frac{\pi(w+s)}{2b}\right] \quad (4.6.4.3)$$

$$k_o = \tanh\left(\frac{\pi w}{2b}\right) \coth\left[\frac{\pi(w+s)}{2b}\right] \quad (4.6.4.4)$$

$$k'_e = \sqrt{1 - k_e^2} \quad (4.6.4.5)$$

$$k'_o = \sqrt{1 - k_o^2} \quad (4.6.4.6)$$

Equations are exact.

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4.6.5 Finite-Thickness Edge-Coupled Striplines

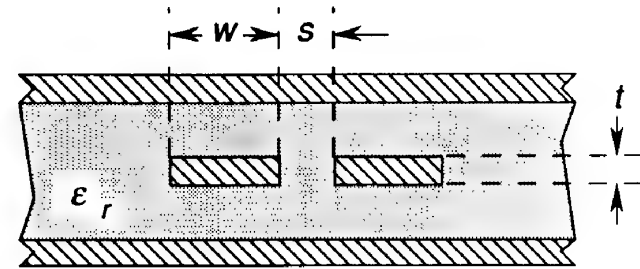


Figure 4.6.5.1: Edge-Coupled Stripline

The equations of Cohn [3] are

$$Z_{0,e}\left(\frac{w}{b}, \frac{t}{b}, \frac{s}{b}\right) = \left\{ \frac{1.0}{Z_0\left(\frac{w}{b}, \frac{t}{b}\right)} - \frac{C_f'(t/b)}{C_f'(0)} \left[\frac{1.0}{Z_0\left(\frac{w}{b}, 0\right)} - \frac{1.0}{Z_{0,e}\left(\frac{w}{b}, 0, \frac{s}{b}\right)} \right] \right\}^{-1} \quad (4.6.5.1)$$

for $s/t \geq 5.0$

$$Z_{0,o}\left(\frac{w}{b}, \frac{t}{b}, \frac{s}{b}\right) = \left\{ \frac{1.0}{Z_0\left(\frac{w}{b}, \frac{t}{b}\right)} - \frac{C_f'(t/b)}{C_f'(0)} \left[\frac{1.0}{Z_{0,o}\left(\frac{w}{b}, 0, \frac{s}{b}\right)} - \frac{1.0}{Z_0\left(\frac{w}{b}, 0\right)} \right] \right\}^{-1} \quad (4.6.5.2)$$

and for $s/t \leq 5.0$

$$Z_{0,o} \left(\frac{w}{b}, \frac{t}{b}, \frac{s}{b} \right) = \left\{ \frac{1.0}{Z_{0,o} \left(\frac{w}{b}, 0, \frac{s}{b} \right)} + \left[\frac{1.0}{Z_{0,o} \left(\frac{w}{b}, 0, \frac{s}{b} \right)} - \frac{1.0}{Z_0 \left(\frac{w}{b}, 0 \right)} \right] - \frac{2.0}{377} \left(\frac{C_f' (t/b)}{\epsilon} - \frac{C_f' (0)}{\epsilon} \right) + \frac{2.0}{377 s} \right\}^{-1} \quad (\Omega) \quad (4.6.5.3)$$

where

$$Z_0 \left(\frac{w}{b}, \frac{t}{b} \right) = \text{finite-thickness impedance of an isolated stripline (see Section 3.6.2)}$$

$$Z_0 \left(\frac{w}{b}, 0 \right) = \text{zero-thickness impedance of an isolated stripline (see Section 3.6.1)}$$

$$Z_{0,e} \left(\frac{w}{b}, 0, \frac{s}{b} \right) = \text{zero-thickness coupled stripline even mode impedance (see Section 4.6.4)}$$

$$Z_{0,o} \left(\frac{w}{b}, 0, \frac{s}{b} \right) = \text{zero-thickness coupled stripline odd mode impedance (see Section 4.6.4)}$$

$$C_f' (t/b) = \frac{0.0885 \epsilon_r}{\pi} \left\{ \frac{2.0}{1.0 - t/b} \ln \left(\frac{1.0}{1.0 - t/b} + 1.0 \right) - \left(\frac{1.0}{1.0 - t/b} - 1.0 \right) \ln \left[\frac{1.0}{(1.0 - t/b)^2} - 1.0 \right] \right\} \quad (\mu\text{F/cm})^\dagger \quad (4.6.5.4)$$

[†] [1]

$$C_f' (0) = \frac{0.0885 \epsilon_r}{\pi} 2.0 \ln 2.0 \quad (\mu\text{F/cm})^\ddagger \quad (4.6.5.5)$$

The equations [8] for even and odd mode conductor losses (incremental inductance rule) and the dielectric losses:

$$\alpha_{c,e} = \frac{2.0 R_s \epsilon_r Z_{0,e}}{\eta^2 b} \left[(s/b - 1.0) \frac{\partial C_{fe}' / \epsilon}{\partial s / b} + (1.0 + t/b) \left(\frac{\partial C_{fe}' / \epsilon}{\partial t / b} + \frac{\partial C_f' / \epsilon}{\partial t / b} \right) + 2.0 \frac{w/b + 1.0 - t/b}{(1.0 - t/b)^2} \right] \quad (4.6.5.6)$$

$$\alpha_{c,o} = \frac{2 R_s \epsilon_r Z_{0,o}}{\eta^2 b} \left[(s/b - 1.0) \frac{\partial C_{fo}' / \epsilon}{\partial s / b} + (1.0 + t/b) \left(\frac{\partial C_{fo}' / \epsilon}{\partial t / b} + \frac{\partial C_f' / \epsilon}{\partial t / b} \right) + 2.0 \frac{w/b + 1.0 - t/b}{(1 - t/b)^2} \right] \quad (4.6.5.7)$$

$$\alpha_{d,e} = \alpha_{d,o} = \frac{\pi \sqrt{\epsilon_r}}{\lambda} \tan \delta \quad (4.6.5.8)$$

where λ is the wavelength. The remaining capacitances are found with the following.

For small s

$$C_{fo}' \left(\frac{t}{b}, \frac{s}{b} \right) = C_f' \left(0, \frac{s}{b} \right) + \epsilon \frac{t}{s} \quad (4.6.5.9)$$

For $s/t \geq 10.0$

$$C_{fo}' \left(\frac{t}{b}, \frac{s}{b} \right) = C_f' \left(\frac{t}{b} \right) \frac{C_{fo}' \left(0, \frac{s}{b} \right)}{C_f' (0)} \quad (4.6.5.10)$$

For $s/t \leq 10.0$

[‡] [1], valid for $w/(b-t) \geq 0.35$ to 1.2% accuracy.

$$C_{f,o} = C_f' \left(0, \frac{s}{b} \right) + \frac{\epsilon Z_0 \left(\frac{s}{t}, 0 \right)}{377} \quad (4.6.5.11)$$

where

$Z_0 \left(\frac{s}{t}, 0 \right)$ = the impedance of a single zero-thickness strip having width s ,

ground plane spacing t , and $\epsilon_r = 1$

The mode velocities are slightly unequal and are calculated with the even and odd mode effective dielectric constants

$$\epsilon_{eff,e} = 1.0 \quad (4.6.5.12)$$

$$\epsilon_{eff,o} = \left[\frac{1.0 + 2.0 Z_{0,o} \left(\frac{w}{b}, 0, \frac{s}{b} \right) Z_0 \left(\frac{s}{b}, 0 \right) / \eta_0^2}{1.0 + 2.0 \epsilon_r Z_{0,o} \left(\frac{w}{b}, 0, \frac{s}{b} \right) Z_0 \left(\frac{s}{t}, 0 \right) / \eta_0^2} \right]^{-1} \quad (4.6.5.13)$$

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4.6.6 Shielded Edge-Coupled Striplines

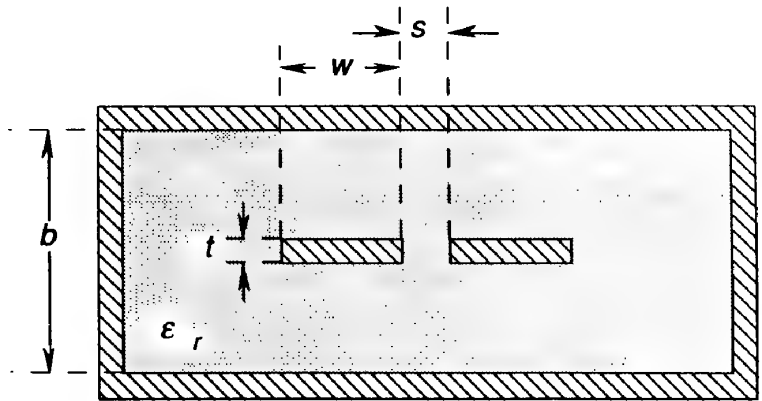


Figure 4.6.6.1: Shielded Edge-Coupled Striplines

Chao [1] describes a variable coupler using this configuration. By varying the side wall distance as the height, b , is varied, the coupling is controlled while keeping the characteristic impedance approximately constant. Chao's equations are

$$Z_{0,e} = \frac{\eta_0}{\sqrt{\epsilon_r}} \frac{\epsilon}{C_{0,e}} \quad (4.6.6.1)$$

$$Z_{0,o} = \frac{\eta_0}{\sqrt{\epsilon_r}} \frac{\epsilon}{C_{0,o}} \quad (4.6.6.2)$$

where

$$\frac{C_{0,e}}{\epsilon} = 2.0 \left(\frac{C_p}{\epsilon} + \frac{C_{fo}}{\epsilon} + \frac{C_{fom}}{\epsilon} \right) \quad (4.6.6.3)$$

$$\frac{C_{0\phi}}{\epsilon} = 2.0 \left(\frac{C_p}{\epsilon} + \frac{C_{fe'}}{\epsilon} + \frac{C_{fom'}}{\epsilon} \right) \quad (4.6.6.4)$$

$$\frac{C_p}{\epsilon} = \frac{2.0 w}{b - t} \quad (4.6.6.5)$$

$$\frac{C_{f\phi'}}{\epsilon} = \frac{2.0}{\pi} \ln \left[1.0 + \coth \frac{\pi s}{2.0 b} \right] + \frac{t}{s} \left(1.0 - \frac{s}{2.0 b} \right) \quad (4.6.6.6)$$

$$\frac{C_{fe'}}{\epsilon} = \frac{2.0}{\pi} \ln \left[1.0 + \tanh \frac{\pi s}{2.0 b} \right] \left(0.89 + \frac{2.14 t}{b} \right) \quad (4.6.6.7)$$

$$\frac{C_{fom'}}{\epsilon} = \frac{2.0}{\pi} \ln \left[1.0 + \coth \frac{\pi m}{2.0 b} + \frac{t}{m} \left(1.0 - \frac{m}{2.0 b} \right) \right] \quad (4.6.6.8)$$

The approximations for $C_{fe'}$ and $C_{fo'}$ are accurate to within a few percent for

$$0.1 \leq t/b \leq 0.5$$

$$s/b < 0.2$$

REFERENCES

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4.6.7 Asymmetric Edge-Coupled Striplines

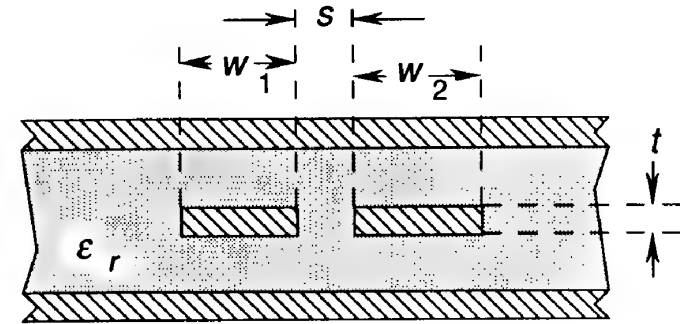


Figure 4.6.7.1: Asymmetric Edge-Coupled Stripline

The general equations for an asymmetric coupler are found in Cristal [3]. For this structure, we define four impedances:

$$Z_{0,e}^1 = \frac{1.0}{v_p C_1} \quad (4.6.7.1)$$

$$Z_{0,e}^2 = \frac{1.0}{v_p C_2} \quad (4.6.7.2)$$

$$Z_{0,o}^1 = \frac{1.0}{v_p (C_1 + 2.0 C_{12})} \quad (4.6.7.3)$$

$$Z_{0,o}^2 = \frac{1.0}{v_p (C_2 + 2.0 C_{12})} \quad (4.6.7.4)$$

where

$$\frac{C_1}{\epsilon} = \text{capacitance of line 1 to ground}$$

$$\frac{C_2}{\epsilon} = \text{capacitance of line 2 to ground}$$

$$\frac{C_{12}}{\epsilon} = \text{capacitance of line 1 to line 2}$$

These can be evaluated with the equations of Cohn [2] and in Section 4.6.5. The coupling and line impedances are defined as:

$$k = \frac{D}{\sqrt{AB}} \quad (4.6.7.5)$$

$$R_a = \sqrt{\frac{A}{B}} \sqrt{AB - D^2} \quad (4.6.7.6)$$

$$R_b = \sqrt{\frac{B}{A}} \sqrt{AB - D^2} \quad (4.6.7.7)$$

where A , B , and D are defined in terms of the admittances (reciprocal of impedances above):

$$A = \frac{Y_{0,e}^1 + Y_{0,e}^1}{2.0} \quad (4.6.7.8)$$

$$B = \frac{Y_{0,e}^2 + Y_{0,e}^2}{2.0} \quad (4.6.7.9)$$

$$D = \frac{Y_{0,e}^1 - Y_{0,e}^1}{2.0} \quad (4.6.7.10)$$

The coupling and the impedance transformation ratio have the restrictions that

$$\frac{1.0}{k^2} \geq \frac{R_2}{R_1} \quad (4.6.7.11)$$

$$\frac{R_2}{R_1} \geq 1.0 \quad (4.6.7.12)$$

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4.7 COUPLED SLABLINES

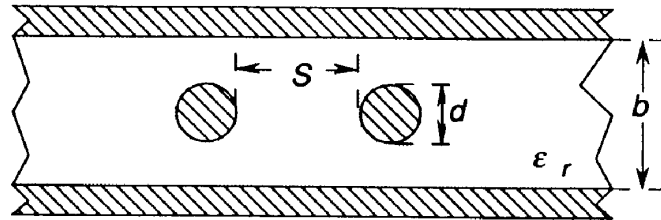


Figure 4.7.1: Coupled Slablins

This structure is used when analyzing coupled wires or coupled rods as used in interdigital filters.

$$Z_{0,e} = \frac{59.952}{\sqrt{\epsilon_r}} \ln \left[\frac{0.523962}{f_1(x) f_2(x, y) f_3(x, y)} \right] \quad (\Omega) \quad (4.7.1.1)$$

$$Z_{0,o} = \frac{59.952}{\sqrt{\epsilon_r}} \ln \left[\frac{0.523962 f_3(x, y)}{f_1(x) f_4(x, y)} \right] \quad (\Omega) \quad (4.7.1.2)$$

where

$$f_1(x) = \frac{x a(x)}{b(x)} \quad (4.7.1.3)$$

$$f_2(x) = \begin{cases} c(y) - x d(y) + e(x) g(y) & y < 0.9 \\ 1.0 + .004 e^{(0.9 - y)} & y \geq 0.9 \end{cases} \quad (4.7.1.4)$$

$$f_3(x) = \tanh \left[\frac{\pi(x + y)}{w} \right] \quad (4.7.1.5)$$

$$f_4(x) = \begin{cases} k(y) - x l(y) + m(x) n(y) & y < 0.9 \\ 1.0 & y \geq 0.9 \end{cases} \quad (4.7.1.6)$$

$$a(x) = 1.0 + e^{(16.0 x - 18.272)} \quad (4.7.1.7)$$

$$b(x) = \sqrt{5.905 - x^4} \quad (4.7.1.8)$$

$$c(y) = -0.8107 y^3 + 1.3401 y^2 - 0.6929 y + 1.0892 + \frac{0.014002}{y} - \frac{0.000636}{y^2} \quad (4.7.1.9)$$

$$d(y) = 0.11 - 0.83 y + 1.64 y^2 - y^3 \quad (4.7.1.10)$$

$$e(x) = -0.15 e^{(-13.0 x)} \quad (4.7.1.11)$$

$$g(y) = 2.23 e^{(-7.01 y + 10.24 y^2 - 27.58 y^3)} \quad (4.7.1.12)$$

$$k(y) = 1.0 + 0.01 \left(-0.0726 - \frac{0.2145}{y} + \frac{0.222573}{y^2} - \frac{0.012823}{y^3} \right) \quad (4.7.1.13)$$

$$l(y) = 0.01 \left(-0.26 + \frac{0.6866}{y} + \frac{0.0831}{y^2} - \frac{0.0076}{y} \right) \quad (4.7.1.14)$$

$$m(x) = -0.1098 + 1.2138 x - 2.2535 x^2 + 1.1313 x^3 \quad (4.7.1.15)$$

$$n(y) = -0.019 - \frac{0.016}{y} + \frac{0.0362}{y^2} - \frac{0.00243}{y^3} \quad (4.7.1.16)$$

where

$$x = d / b \quad (4.7.1.17)$$

$$y = s / b \quad (4.7.1.18)$$

The ground is required to be $\geq 5(2d + s)$ wide.

The equations for even and odd mode losses [6]:

$$\alpha_{c,e} = \frac{R_s \epsilon_r Z_{0,e}}{\eta^2 b} \left[(s/b - 1.0) \frac{\partial C_{oe}/\epsilon}{\partial s/b} + (1.0 + d/b) \left(\frac{\partial C_{oe}/\epsilon}{\partial d/b} \right) \right] \quad (4.7.1.19)$$

$$\alpha_{c,o} = \frac{R_s \epsilon_r Z_{0,o}}{\eta^2 b} \left[(s/b - 1.0) \frac{\partial C_{oo}/\epsilon}{\partial s/b} + (1.0 + d/b) \left(\frac{\partial C_{oe}/\epsilon}{\partial d/b} \right) \right] \quad (4.7.1.20)$$

$$\alpha_{d,e} = \alpha_{d,o} = \frac{\pi \sqrt{\epsilon_r}}{\lambda} \tan \delta \quad (4.7.1.21)$$

Rosloniec [5] suggests we set the initial guesses for synthesis to:

$$x^{(0)} = \frac{4.0}{\pi \exp \left(\frac{-Z_0}{59.952 \sqrt{0.987 - 0.171 k - 1.723 k^2}} \right)} \quad (4.7.1.22)$$

$$y^{(0)} = \frac{1.0}{\pi \ln \left(\frac{r+1.0}{r-1.0} \right)} - x^{(0)} \quad (4.7.1.23)$$

Newton's method is then used by the program given in [5] to solve for x and y based on the designer's required $Z_{0,o}$ and $Z_{0,e}$.

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4.8 COUPLED SUSPENDED SUBSTRATE OR COUPLED INVERTED MICROSTRIP LINES

The unshielded and offset versions of these structures are not available, however the following structures can be used when the line is centered in a shield.

4.8.1 Edge-Coupled Suspended Substrate Lines

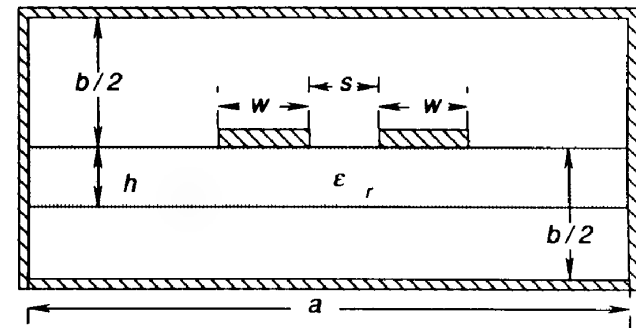


Figure 4.8.1.1: Edge-Coupled Suspended Substrate Lines

In Wang *et al.* [4] edge-coupled suspended substrate is analyzed. As for single lines, the suspended substrate configuration has lower loss and less sensitivity to physical dimensions than an equivalent microstrip or stripline.

The equations were derived by curve-fitting results of the method of lines. They are

$$Z_{0,e} = \frac{\eta_0 \pi}{2.0 \sqrt{\epsilon_{r,e}} \ln \left(2.0 \frac{1.0 + \sqrt{k_1'}}{1.0 - \sqrt{k_1'}} \right)} \frac{1.0}{1.0} \quad (4.8.1.1)$$

$$Z_{0,o} = \frac{\eta_0 \pi}{2.0 \sqrt{\epsilon_{r,o}} \ln \left(2.0 \frac{1.0 + \sqrt{k_2'}}{1.0 - \sqrt{k_2'}} \right)} \frac{1.0}{1.0} \quad (4.8.1.2)$$

and

$$\epsilon_{r,e} = \left\{ \left[1.0 - D \frac{s}{a} + H \ln \frac{w}{a} + R \right] (\ln \epsilon_r) \epsilon_r^{0.06} \left(\frac{\epsilon_r}{2.8} \right)^{(b/a - 0.4) / 3.0} \right\}^{-2} \quad (4.8.1.3)$$

$$\epsilon_{r\rho} = \left[1.0 - E \left(\ln \frac{s}{a} - F \right) \left(\ln \frac{w}{a} - G \right) (\ln \epsilon_r) \epsilon_r^{0.08} \right]^{-2} \quad (4.8.1.4)$$

where

$$k_1' = \sqrt{1.0 - k_1^2} \quad (4.8.1.5)$$

$$k_1 = \frac{\operatorname{tg} \vartheta_2'}{\operatorname{tg} \vartheta_1'} \quad (4.8.1.6)$$

$$k_2' = \sqrt{1.0 - k_2^2} \quad (4.8.1.7)$$

$$k_2 = \sqrt{\frac{1.0 + \operatorname{tg}^2 \vartheta_2'}{1.0 - \operatorname{tg}^2 \vartheta_1'}} \quad (4.8.1.8)$$

$$\vartheta_1' = Q_1 \operatorname{th} \left[\frac{Q_2 (2.0 w + s)}{b} \right] + Q_3 \quad (4.8.1.9)$$

$$\vartheta_2' = Q_1 \operatorname{th} \left(\frac{Q_2 s}{b} \right) + Q_3 \quad (4.8.1.10)$$

$$Q_1 = 0.0361168 e^{5.12 k} + 1.4404 \quad (4.8.1.11)$$

$$Q_2 = -0.0316177 e^{3.947 k} + 1.07319 \quad (4.8.1.12)$$

$$Q_3 = 0.15988 k^3 - 0.0895 k^2 + 0.02535 k \quad (4.8.1.13)$$

$$k = \sqrt{1.0 - \left(\frac{e^{\pi a/b} - 2.0}{e^{\pi a/b} + 2.0} \right)^4} \quad (4.8.1.14)$$

$$D = k_{11} \left(\frac{h}{a} \right)^3 + k_{12} \left(\frac{h}{a} \right)^2 + k_{13} \left(\frac{h}{a} \right) + k_{14} \quad (4.8.1.15)$$

$$E = 0.0001 \left(7.7291 \left\{ \frac{b}{a} + 6.9924 \left[9.5 \left(\frac{h}{a} \right)^2 - 1.0 \right] \right\} \right.$$

$$\times \left(\ln \frac{h}{a} + 2.1657 \right) + 561.62 \left(\frac{h}{a} \right)^2 - 6.7772 \left[\frac{h}{a} \right] \quad (4.8.1.16)$$

$$F = 60.4762 \left\{ \left(\frac{b}{a} - \frac{6.0 h}{a} \right)^2 \left(1.0 + \frac{2.0 h}{a} \right) - 1.6012 \left[1.0 - 8.0 \left(\frac{h}{a} \right)^2 \right] \right\}$$

$$\times (e^{h/a} - 1.04684) + 1721.61 \left(\frac{h}{a} \right)^2 + 19.027 \left(\frac{h}{a} \right) \quad (4.8.1.17)$$

$$G = 0.05154 \left\{ - \left(\frac{b}{a} - 0.8 \right)^2 + 0.3653 \left[1.0 - 8.0 \left(\frac{h}{a} \right)^2 \right] \right\}$$

$$\times (e^{h/a} + 54.554) + 6551.53 \left(\frac{h}{a} \right)^3 + 34.76 \left(\frac{h}{a} \right) \quad (4.8.1.18)$$

$$H = k_{21} \left(\frac{h}{a} \right)^3 + k_{22} \left(\frac{h}{a} \right)^2 + k_{23} \left(\frac{h}{a} \right) + k_{24} \quad (4.8.1.19)$$

$$R = k_{31} \left(\frac{h}{a} \right)^2 + k_{32} \left(\frac{h}{a} \right) + k_{33} \quad (4.8.1.20)$$

$$k_{11} = -28.794 \left(\frac{b}{a} \right)^2 + 29.958 \left(\frac{b}{a} \right) + 52.861 \quad (4.8.1.21)$$

$$k_{12} = 0.92658 \left(\frac{b}{a} \right)^2 + 3.2713 \left(\frac{b}{a} \right) - 20.653 \quad (4.8.1.22)$$

$$k_{13} = -1.5075 \left(\frac{b}{a} \right)^2 + 1.8914 \left(\frac{b}{a} \right) + 0.75778 \quad (4.8.1.23)$$

$$k_{14} = -0.25477 \left(\frac{b}{a} \right)^3 + 0.60480 \left(\frac{b}{a} \right)^2 - 0.48538 \left(\frac{b}{a} \right) + 0.16713 \quad (4.8.1.24)$$

$$k_{21} = -11.459 \left(\frac{b}{a} \right)^2 + 18.744 \left(\frac{b}{a} \right) - 17.109 \quad (4.8.1.25)$$

$$k_{22} = 7.4499 \left(\frac{b}{a} \right)^2 - 11.927 \left(\frac{b}{a} \right) + 7.4606 \quad (4.8.1.26)$$

$$\times \left(\ln \frac{1.0}{\sqrt{\epsilon_r}} + \frac{1.0}{\sqrt{\epsilon_r}} - 1.0 \right) \Big]^{-1} \quad (4.8.2.3)$$

$$\epsilon_{eff,e} = \frac{1.0}{\left[1.0 + (H - P \ln \frac{w}{b}) \ln \frac{1.0}{\sqrt{\epsilon_r}} \right]^2} \quad (4.8.2.4)$$

where

$$S_o = -0.1073 + 1.67080 \frac{h}{b} + 0.007484 \frac{a}{b} \quad (4.8.2.5)$$

$$T_o = 0.4768 + 2.1295 \frac{h}{b} - 0.01278 \frac{a}{b} \quad (4.8.2.6)$$

$$H_o = 0.7210 - 0.3568 \frac{h}{b} + 0.02132 \frac{a}{b} \quad (4.8.2.7)$$

$$P_o = -0.3035 + 0.3743 \frac{h}{b} + 0.07274 \frac{a}{b} \quad (4.8.2.8)$$

$$S_e = -2.6528 + 0.9452 \frac{h}{b} - 0.4531 \frac{a}{b} \quad (4.8.2.9)$$

$$P_e = 1.4793 - 1.1903 \frac{h}{b} - 0.04511 \frac{a}{b} \quad (4.8.2.10)$$

$$H_e = 0.2245 + 0.7192 \frac{h}{b} - 0.1022 \frac{a}{b} \quad (4.8.2.11)$$

$$P_e = 0.001356 + 0.06590 \frac{h}{b} + 0.01951 \frac{a}{b} \quad (4.8.2.12)$$

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- [1] Bhartia, Brakash, and Protap Pramanick, "Computer-Aided Design Models for Broadside-Coupled Striplines and Millimeter-Wave Suspended Substrate Microstrip Lines," *IEEE Transactions on Microwave Theory and Techniques*, Vol. MTT-36, No. 11, November 1988, pp. 1476-1481, and "Corrections to 'Computer-Aided Design Models for Broadside-Coupled Striplines and Millimeter-Wave Suspended Substrate Microstrip Lines,'" Vol. MTT-37, No. 10, October 1989, p. 1658.

- [2] Shu, Yong-hui, Xiao-xia Qi, and Yun-ji Wang, "Analysis Equations for Shielded Suspended Substrate Microstrip Line and Broadside-Coupled Stripline," 1980 *IEEE MTT-S International Microwave Symposium Digest*, pp. 693-696.

4.9 OTHER COUPLED LINES

4.9.1 Coupled Slotlines

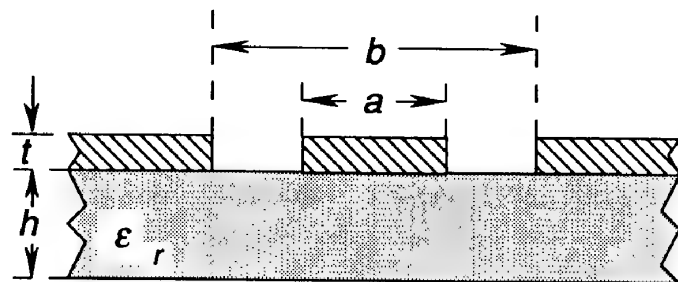


Figure 4.9.1.1: Coupled Slotlines

As can be seen from the figure, coupled slotlines are strongly related to the coplanar waveguide structure and the slotline structure. Knorr and Kuchler calculate the impedances and effective dielectric constant of this structure and present curves in [1]. Although closed-form equations are not yet available, some observations can be made about the structure at the extremes.

For a narrow center conductor or tight coupling ($a/h \ll 1.0$), the structure resembles a single slotline with slot width $b - a$ whose impedance is $Z_{0,sl}$ and

$$Z_{0,e} = \frac{Z_{0,sl}}{2.0} \quad (4.9.1.1)$$

For very wide center conductor ($a/h \gg 1.0$) or loose coupling between the slots, the structure becomes two uncoupled slotlines or CPW. In this situation, if we calculate the characteristic impedance of the coplanar waveguide structure, $Z_{0,cpw}$, the symmetric coupled slotline has

$$Z_{0,o} = 2.0 Z_{0,cpw} \quad (4.9.1.2)$$

REFERENCE

- [1] Knorr, Jeffrey B., and Klaus-Dieter Kuchler, "Analysis of Coupled Slots and Coplanar Strips on Dielectric Substrate," *IEEE Transactions on Microwave Theory and Techniques*, Vol. MTT-23, No. 7, July 1975, pp. 541–548.

Chapter 5

Transmission Line Components and Discontinuities

5.1 DISCONTINUITY MODELS

A discontinuity occurs whenever our assumption that the line can be modeled as a uniform line having length is violated. This can be a bend, a step in width, a T junction, a change in dielectric constant, a wall in close proximity, or a transition to another type of line. We need to calculate discontinuity effects to minimize them and in other cases to use them to create a desired circuit (e.g., a filter).

Often the discontinuity is small relative to the signal being propagated and can be accurately approximated by lumped-element equivalent circuits. Lumped-element models are best for understanding the effects of and remedies for discontinuities. For more accuracy over frequency the model may include transmission lines. As shown in Chapter 1, for small size relative to a wavelength, a transmission line can often be replaced by an equivalent lumped element. Several equivalent models are possible for each discontinuity.

In this chapter the equations for calculating the discontinuity models are given. These models are made up of the usual lumped elements and lengths of transmission lines. Reference planes (denoted T_n and a dotted line on each figure) indicate the physical location of the discontinuity model. The physical structure between the reference planes is replaced with the model.

5.2 COAXIAL DISCONTINUITIES

The original work in coaxial discontinuities was published in 1944 in two articles, by Whinnery *et al.*, [7] and [8]. Somlo followed these up with corrections and closed-form equations [5, 6] based on numerical calculations including higher order modes.

5.2.1 Coaxial Step in Inner Conductor

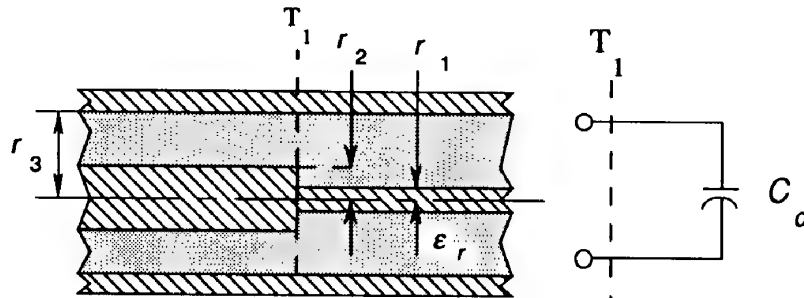


Figure 5.2.1.1: Coaxial Step in Inner Conductor

The discontinuity capacitance for a step on the inner conductor is [6, 8]

$$C_d = 2.0 \epsilon_r \pi r_2 C_{d1} \quad (5.2.1.1)$$

where

$$C_{d1} = \frac{\epsilon}{100.0 \pi} \left[\frac{\alpha^2 + 1.0}{\alpha} \ln \left(\frac{1.0 + \alpha}{1.0 - \alpha} \right) - 2.0 \ln \left(\frac{4.0 \alpha}{1.0 - \alpha^2} \right) \right] + 1.11 \times 10^{-15} (1.0 - \alpha) (\tau - 1.0) \quad (\text{F/cm}) \quad (5.2.1.2)$$

$$\alpha = \frac{r_3 - r_2}{r_3 - r_1} \quad (5.2.1.3)$$

$$\tau = \frac{r_3}{r_1} \quad (5.2.1.4)$$

which is accurate to ± 0.3 fF/cm for

$$0.01 \leq \alpha \leq 1.0$$

$$1.0 \leq \tau \leq 6.0$$

In Young [9], Green notes that the effect of the discontinuity on the fields is minimal at a distance greater than $2c$ away from the step for this structure. The effect of frequency can be approximately corrected [8] with the formula

$$\frac{C(f)}{C_d} = \frac{1.0}{\sqrt{1.0 - (2.0 r_2 / \lambda)^2}} \quad (5.2.1.5)$$

which is valid for

$$0.0 \leq r_2 \leq \lambda / 2.0$$

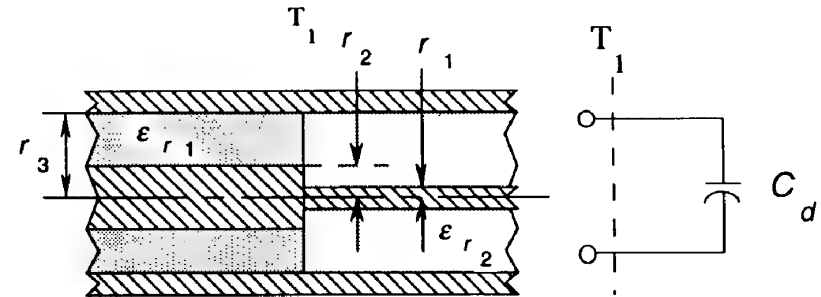


Figure 5.2.1.2: Coaxial Step on Inner Conductor with Dielectric Change

For a dielectric discontinuity in the same plane as the inner conductor step (Figure 5.2.1.2), the approximating equations are

$$C_d \cong 2.0 \pi r_3 \epsilon_r C_{d1} \quad (5.2.1.6)$$

where

$$C_{d1} = \frac{\epsilon}{100.0 \pi} \left[\frac{\alpha^2 + 1.0}{\alpha} \ln \left(\frac{1.0 + \alpha}{1.0 - \alpha} \right) - 2.0 \ln \left(\frac{4.0 \alpha}{1.0 - \alpha^2} \right) \right] + 1.11 \times 10^{-15} (1.0 - \alpha) (\tau - 1.0) \quad (\text{F/cm}) \quad (5.2.1.7)$$

$$\alpha = \frac{r_3 - r_2}{r_3 - r_1} \quad (5.2.1.8)$$

$$\tau = \frac{r_3}{r_1}$$

(5.2.1.9)

This structure is encountered in coaxial air lines that are supported by dielectric beads. The step on the inner conductor is introduced to compensate for the dielectric discontinuity.

In this situation and referring to the figure above, $\epsilon_{r1} = \epsilon_0$ and ϵ_{r2} is the bead dielectric. Two of these discontinuities occur as shown in the figure below.

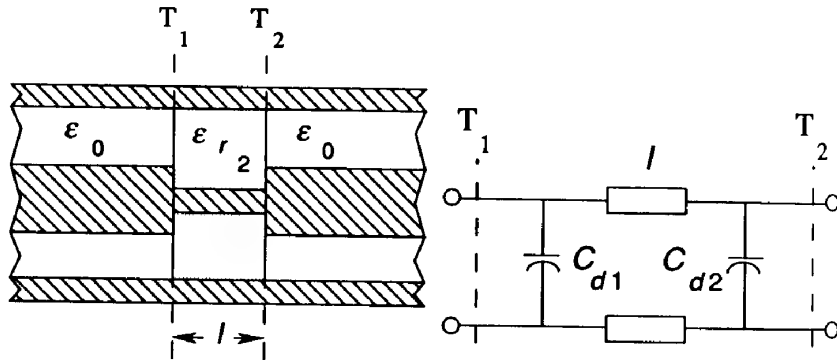


Figure 5.2.1.3: Compensated Coaxial Dielectric Bead Support

REFERENCES

- [1] Gwarek, Wojciech K., "Computer-Aided Analysis of Arbitrarily Shaped Coaxial Discontinuities," *IEEE Transactions on Microwave Theory and Techniques*, Vol. MTT-36, No. 2, February 1988, pp. 337-342.
- [2] Marcuvitz, N., *Waveguide Handbook*, McGraw-Hill, New York, 1955. Reprint with errata and preface by N. Marcuvitz, Peter Peregrinus Ltd., 1986.
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- [5] Somlo, P.I., "Calculating Coaxial Transmission-Line Step Capacitance," *IEEE Transactions on Microwave Theory and Techniques*, Vol. MTT-11, September 1963, p. 454. (Gives curves fit to charts of Whinnery *et al.*)

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- [7] Whinnery, J.R., *et al.*, "Coaxial-Line Discontinuities," *Proceedings of the IRE*, November 1944, pp. 695-709.
- [8] Whinnery, J.R., *et al.*, "Equivalent Circuits for Discontinuities in Transmission Lines," *Proceedings of the IRE*, February 1944, pp. 99-114.
- [9] Young, Leo, ed., *Advances in Microwaves*, Volume 2, "The Numerical Solution of Transmission Line Problems," by Harry E. Green, Academic Press, New York, 1967. (Improves on [7] and [8] for $b/a > 2$ and gives results in equation and table form rather than graphically.)

5.2.2 Coaxial Step in Outer Conductor

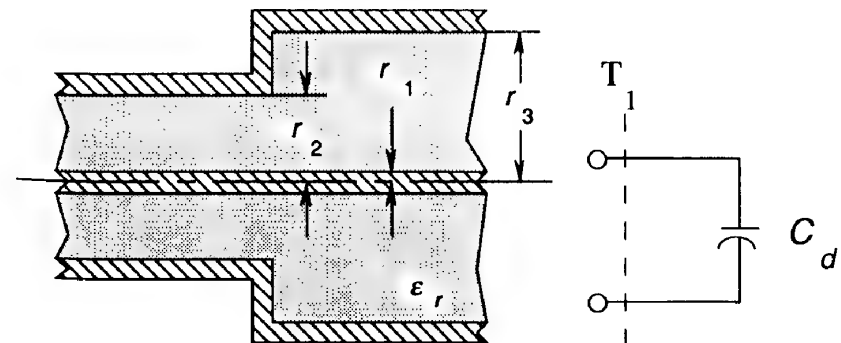


Figure 5.2.2.1: Coaxial Step in Outer Conductor

A step on the outer conductor is modeled by a capacitance

$$C_d = 2.0 \epsilon_r \pi r_1 C_{d2}' \quad (5.2.2.1)$$

where

$$C_{d2}' = \frac{\epsilon}{100.0 \pi} \left(\frac{\alpha^2 + 1.0}{\alpha} \ln \frac{1.0 + \alpha}{1.0 - \alpha} - 2.0 \ln \frac{4.0 \alpha}{1.0 - \alpha^2} \right) + 4.12 \times 10^{-15} (0.8 - \alpha) (\tau - 1.4) \quad (\text{F/cm}) \quad (5.2.2.2)$$

$$\alpha = \frac{r_2 - r_1}{r_3 - r_1} \quad (5.2.2.3)$$

$$\tau = \frac{r_3}{r_1} \quad (5.2.2.4)$$

The equation is accurate to ± 0.6 fF/cm for

$$0.01 \leq \alpha \leq 0.7$$

$$1.5 \leq \tau \leq 6.0$$

The effect of frequency can be approximately corrected [6] with the formula

$$\frac{C(f)}{C_d} = \frac{1.0}{\sqrt{1.0 - (2.0 r_3 / \lambda)^2}} \quad (5.2.2.5)$$

which is valid for

$$0.0 \leq r_3 \leq \lambda / 2.0$$

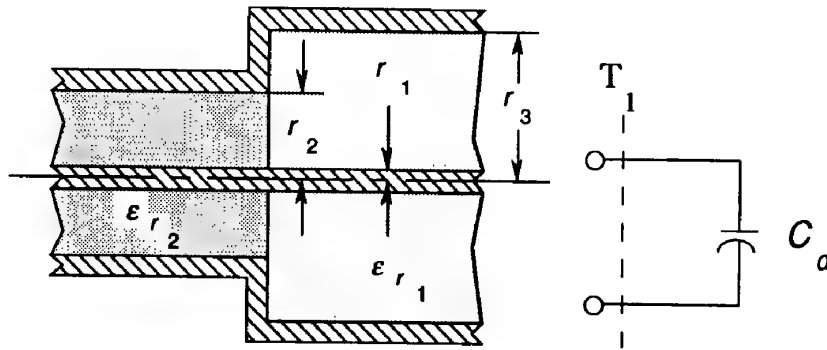


Figure 5.2.2.2: Coaxial Step in Outer Conductor with Dielectric Change

The capacitance of the step in the outer conductor with a corresponding dielectric step is approximated with

$$C_d = 2.0 \epsilon_{r2} \pi r_1 C_{d2} \quad (5.2.2.6)$$

where

$$C_{d2} = \frac{\epsilon}{100.0 \pi} \left(\frac{\alpha^2 + 1.0}{\alpha} \ln \frac{1.0 + \alpha}{1.0 - \alpha} - 2.0 \ln \frac{4.0 \alpha}{1.0 - \alpha^2} \right) + 4.12 \times 10^{-15} (0.8 - \alpha) (\tau - 1.4) \quad (\text{F/cm}) \quad (5.2.2.7)$$

$$\alpha = \frac{r_2 - r_1}{r_3 - r_1} \quad (5.2.2.8)$$

$$\tau = \frac{r_3}{r_1} \quad (5.2.2.9)$$

REFERENCES

- [1] Gwarek, Wojciech K., "Computer-Aided Analysis of Arbitrarily Shaped Coaxial Discontinuities," *IEEE Transactions on Microwave Theory and Techniques*, Vol. MTT-36, No. 2, February 1988, pp. 337-342.
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- [3] Saad, T., ed., *The Microwave Engineer's Handbook*, Vol. 1, Artech House, Norwood, MA, 1971.
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- [6] Whinnery, J.R., *et al.*, "Equivalent Circuits for Discontinuities in Transmission Lines," *Proceedings of the IRE*, February 1944, pp. 99-114.

5.2.3 Coaxial Simultaneous Step in Inner and Outer Conductors

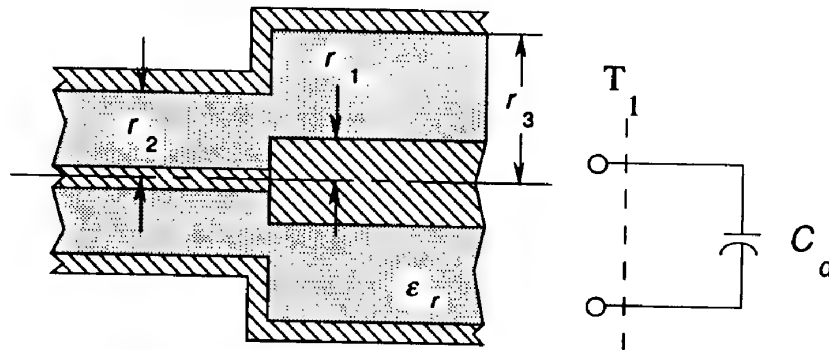


Figure 5.2.3.1: Coaxial Step in Inner and Outer Conductors

The discontinuity capacitance of this structure is modeled by the sum of two discontinuity capacitances:

$$C_d = 2.0 \epsilon_r \pi c \dot{C}_{d1} + 2.0 \epsilon_r \pi a \dot{C}_{d2} \quad (5.2.3.1)$$

where

$$\dot{C}_{d1} = \frac{\epsilon}{100.0 \pi} \left(\frac{\alpha^2 + 1.0}{\alpha} \ln \frac{1.0 + \alpha}{1.0 - \alpha} - 2.0 \ln \frac{4.0 \alpha}{1.0 - \alpha^2} \right) + 1.11 \times 10^{-15} (1.0 - \alpha) (\tau - 1.0) \quad (\text{F/cm}) \quad (5.2.3.2)$$

$$\dot{C}_{d2} = \frac{\epsilon}{100.0 \pi} \left(\frac{\alpha^2 + 1.0}{\alpha} \ln \frac{1.0 + \alpha}{1.0 - \alpha} - 2.0 \ln \frac{4.0 \alpha}{1.0 - \alpha^2} \right) + 4.12 \times 10^{-15} (0.8 - \alpha) (\tau - 1.4) \quad (\text{F/cm}) \quad (5.2.3.3)$$

The effect of frequency can be approximately corrected [6] with the formula

$$\frac{C(f)}{C_d} = \frac{1.0}{\sqrt{1.0 - (2.0 b / \lambda)^2}} \quad (5.2.3.4)$$

which is valid for

$$0.0 \leq b \leq \lambda / 2.0$$

Some simple equations that give good results for the discontinuity capacitance of this structure over frequency are given in [1]:

$$C(f) \equiv a + \frac{b}{\sqrt{1.0 - \left(\frac{f}{f_0}\right)^2}} \quad (5.2.3.5)$$

where

a, b are functions of the dimensions and dielectric constants ϵ_{r1} and ϵ_{r2} .

f_0 is the cutoff frequency of the first TM -mode.

The values of a, b , and f_0 from [1] are given in the table below.

ϵ_{r1}	ϵ_{r2}	f_0	a	b
1	1	49.42 GHz	2.125	12.88
1	2	34.94 GHz	4.339	25.69
1	10	15.63 GHz	20.93	128.9

Table 5.2.3.1: Step Capacitance Frequency Dependent Parameters

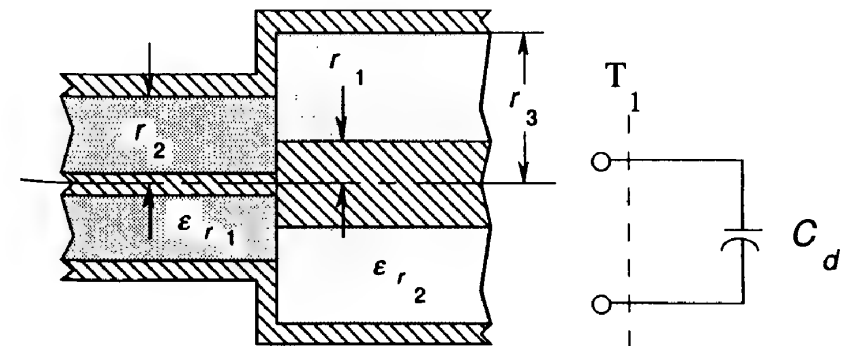


Figure 5.2.3.2: Coaxial Step in Inner and Outer Conductor with Dielectric Change

For lines with a discontinuity in the dielectric as well

$$C_d \equiv 2.0 \epsilon_{r1} \pi r_1 \dot{C}_{d2} + 2.0 \epsilon_{r2} \pi r_2 \dot{C}_{d1} \quad (5.2.3.6)$$

where

$$C'_{d1} = \frac{\epsilon}{100.0 \pi} \left[\frac{\alpha_1^2 + 1.0}{\alpha_1} \ln \left(\frac{1.0 + \alpha_1}{1.0 - \alpha_1} \right) - 2.0 \ln \left(\frac{4.0 \alpha_1}{1.0 - \alpha_1^2} \right) \right] + 1.11 \times 10^{-15} (1.0 - \alpha_1) (\tau_1 - 1.0) \quad (\text{F/cm}) \quad (5.2.3.7)$$

$$C'_{d2} = \frac{\epsilon}{100.0 \pi} \left[\frac{\alpha_2^2 + 1.0}{\alpha_2} \ln \frac{1.0 + \alpha_2}{1.0 - \alpha_2} - 2.0 \ln \frac{4.0 \alpha_2}{1.0 - \alpha_2^2} \right] + 4.12 \times 10^{-15} (0.8 - \alpha_2) (\tau_2 - 1.4) \quad (\text{F/cm}) \quad (5.2.3.8)$$

and

$$\alpha_1 = \frac{r_2 - r_1}{r_2 - r_0} \quad (5.2.3.9)$$

$$\alpha_2 = \frac{r_2 - r_1}{r_3 - r_1} \quad (5.2.3.10)$$

$$\tau_1 = \frac{r_2}{r_0} \quad (5.2.3.11)$$

$$\tau_2 = \frac{r_3}{r_1} \quad (5.2.3.12)$$

Another type of coaxial step (shown below) is modeled by Young [7].

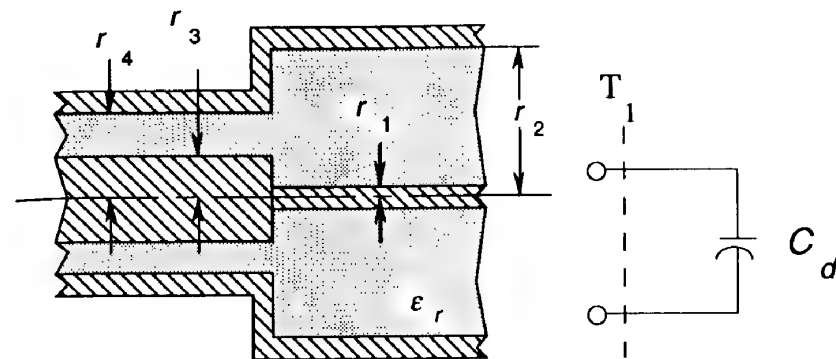


Figure 5.2.3.3: Another Coaxial Step in Inner and Outer Conductors

Young suggests modeling this with the series combination of two equivalent steps (similar to above). The equivalent steps are a step on the inner conductor and a step on the outer conductor.

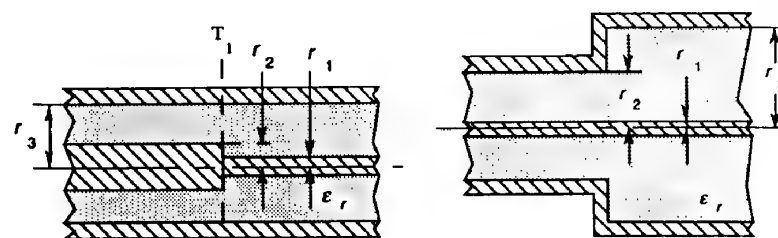


Figure 5.2.3.4: Double Step is Calculated from Equivalent Discontinuities

First we calculate from the original discontinuity's dimensions (subscripted o)

$$r = 0.5 \exp \left(\frac{\log r_1 \log \frac{r_4}{r_3} - \log r_3 \log \frac{r_2}{r_1}}{\log \frac{r_4}{r_3} - \log \frac{r_2}{r_1}} \right) \quad (5.2.3.13)$$

The step on the inner conductor portion of the equivalent model is now calculated with

$$r_1 = r_{1,o}$$

$$r_2 = r_{3,o}$$

$$r_3 = r$$

and the step on the outer conductor portion of the equivalent model is calculated with

$$r_1 = r$$

$$r_2 = r_{4,\rho}$$

$$r_3 = r_{2,\rho}$$

The accuracy of this approximation is within 10% of numeric calculations.

REFERENCES

- [1] Gogioso, L., *et al.*, "A Variational Approach to Compute the Equivalent Capacitance of Coaxial Line Discontinuities," *1979 IEEE MTT-S International Microwave Symposium Digest*, pp. 580–582.
- [2] Gwarek, Wojciech K., "Computer-Aided Analysis of Arbitrarily Shaped Coaxial Discontinuities," *IEEE Transactions on Microwave Theory and Techniques*, Vol. MTT-36, No. 2, February 1988, pp. 337–342.
- [3] Saad, T., ed., *The Microwave Engineer's Handbook*, Vol. 1, Artech House, Norwood, MA, 1971.
- [4] Silvester, P., and Ivan A. Cermak, "Analysis of Coaxial Line Discontinuities by Boundary Relaxation," *IEEE Transactions on Microwave Theory and Techniques*, Vol. MTT-17, No. 8, August 1969, pp. 489–495.
- [5] Somlo, P.I., "The Computation of Coaxial Line Step Capacitances," *IEEE Transactions on Microwave Theory and Techniques*, Vol. MTT-15, No. 1, January 1967, pp. 48–53. (Corrections to Whinnery *et al.* and Somlo 1963.)
- [6] Whinnery, J.R., *et al.*, "Equivalent Circuits for Discontinuities in Transmission Lines," *Proceedings of the IRE*, February 1944, pp. 99–114.
- [7] Young, Leo, ed., *Advances in Microwaves*, Volume 2, "The Numerical Solution of Transmission Line Problems," by Harry E. Green, Academic Press, New York, 1967. (Improves on [82] and [83] for $b/a > 2$ and gives results in equation form rather than graphically)

5.2.4 Gap in Coaxial Center Conductor

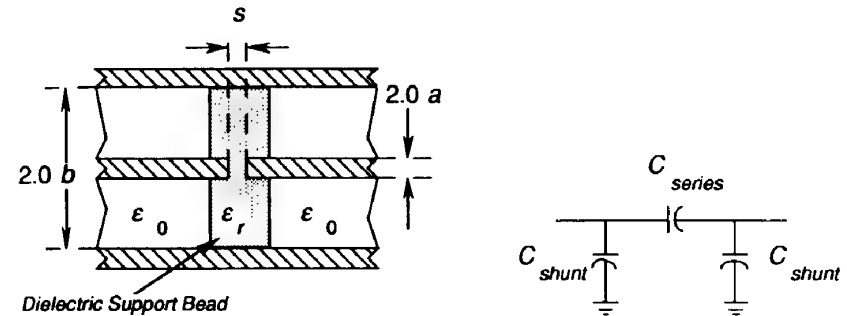


Figure 5.2.4.1: Gap in Coaxial Center Conductor

Dawirs [1] shows that the coaxial gap can be modeled as a π capacitor circuit. By comparison of various source's equations with measurement data it is determined that the best model is to use the C_{series} of Young [5], and the C_{shunt} of Whinnery *et al.* [4]. The value of C_{shunt} is most easily calculated with the closed-form equations of Somlo [3].

For the purposes of calculation the field is assumed to concentrate in the dielectric support (higher ϵ_r than the surrounding dielectric). The ϵ_r of the support bead is used for the discontinuity calculations.

$$C_{series} = \frac{0.7056 \epsilon_r a^2}{s} \text{ (pF)} \quad (5.2.4.1)$$

where a and s are in inches. Equation (5.2.4.1) is valid for

$$b - a < \lambda$$

$$s \ll \lambda$$

$$s \ll b - a$$

where λ is the free-space wavelength.

$$C_{shunt} = 2.0 \epsilon_r \pi b C_{d1}' \quad (5.2.4.2)$$

where

$$C'_{d1} = \frac{2.0 \epsilon}{100.0 \pi} (1.477 - \ln 4.0 \alpha) \quad (5.2.4.3)$$

$$\alpha = 1.0 - \frac{a}{b} \quad (5.2.4.4)$$

Also see the coaxial open end discontinuity, which is the above structure with a short inserted in the gap. The open end, therefore, has half the value of C_{series} and the same value for C_{shunt} (however, there is only one C_{shunt} in the model).

REFERENCES

- [1] Dawirs, Harvel N., "Equivalent Circuit of a Series Gap in the Center Conductor of a Coaxial Transmission Line," *IEEE Transactions on Microwave Theory and Techniques*, MTT-17, No. 2, February 1969, pp. 127-129.
- [2] Gwarek, Wojciech K., "Computer-Aided Analysis of Arbitrarily Shaped Coaxial Discontinuities," *IEEE Transactions on Microwave Theory and Techniques*, Vol. MTT-36, No. 2, February 1988, pp. 337-342.
- [3] Somlo, P.I., "The Computation of Coaxial Line Step Capacitances," *IEEE Transactions on Microwave Theory and Techniques*, Vol. MTT-15, No. 1, January 1967, pp. 48-53. (Corrections to Whinnery, *et al.* and Somlo 1963.)
- [4] Whinnery, J.R., *et al.*, "Coaxial-Line Discontinuities," *Proceedings of the IRE*, Vol. 32, November 1944, pp. 695-709.
- [5] Young, Leo, "The Practical Realization of Series-Capacitive Couplings for Microwave Filters," *Microwave Journal*, December 1962, pp. 79-81.

5.2.5 Coaxial Open Circuit

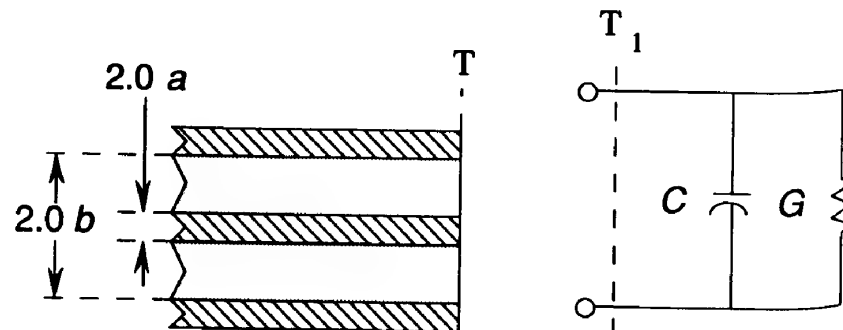


Figure 5.2.5.1: Coaxial Open Circuit

The coaxial open circuit can be part of a filter or other circuit. The fringing fields of an open circuit can also be used to probe a dielectric sample and measure its permittivity.

The equivalent circuit's conductance is calculated with:

$$\frac{G}{Y_0} = \frac{1.0}{\ln \frac{a}{b}} \int_0^{\pi/2.0} \frac{d\theta}{\sin \theta} [J_0(k a \sin \theta) - J_0(k b \sin \theta)]^2 \quad (5.2.5.1)$$

$$\approx \frac{2}{3 \ln \frac{a}{b}} \left[\frac{\pi^2 (b^2 - a^2)}{\lambda^2} \right]^2 \quad (5.2.5.2)$$

The susceptance is:

$$\begin{aligned} \frac{B}{Y_0} = \frac{1.0}{\pi \ln \frac{a}{b}} \int_0^{\pi} \left[2.0 \operatorname{Si} \left(k \sqrt{a^2 + b^2 - 2.0 a b \cos \phi} \right) \right. \\ \left. - \operatorname{Si} \left(2.0 k a \sin \frac{\phi}{2.0} \right) - \operatorname{Si} \left(2.0 k b \sin \frac{\phi}{2.0} \right) d\phi \right] \end{aligned} \quad (5.2.5.3)$$

$$\approx \frac{8.0 (a + b)}{\lambda_0 \ln \frac{a}{b}} \left[E \left(\frac{2.0 \sqrt{a b}}{a + b} \right) - 1.0 \right] \quad (5.2.5.4)$$

where

$E(x)$ = complete elliptic integral of the second kind

$\operatorname{Si}(x)$ = sine integral function

$J_0(x)$ = Bessel function of the first kind and order 0

λ_0 = free-space wavelength

$$k = \frac{2.0 \pi}{\lambda_0}$$

The approximations are valid for

$$a \ll \lambda_0$$

$$b \ll \lambda$$

Stated accuracy of the exact equations is better than 10%. The approximations agree to better than 15% for

$$\frac{a-b}{\lambda} < 0.10$$

$$\frac{a}{b} \geq 2.0$$

with increasing accuracy for increasing a/b (higher impedances).¹

REFERENCES

- [1] Gwarek, Wojciech K., "Computer-Aided Analysis of Arbitrarily Shaped Coaxial Discontinuities," *IEEE Transactions on Microwave Theory and Techniques*, Vol. MTT-36, No. 2, February 1988, pp. 337-342.
- [2] Marcuvitz, N., *Waveguide Handbook*, McGraw-Hill, New York, 1955. Reprint with errata and preface by N. Marcuvitz, Peter Peregrinus Ltd., 1986.

5.2.6 Enclosed Coaxial Open Circuit

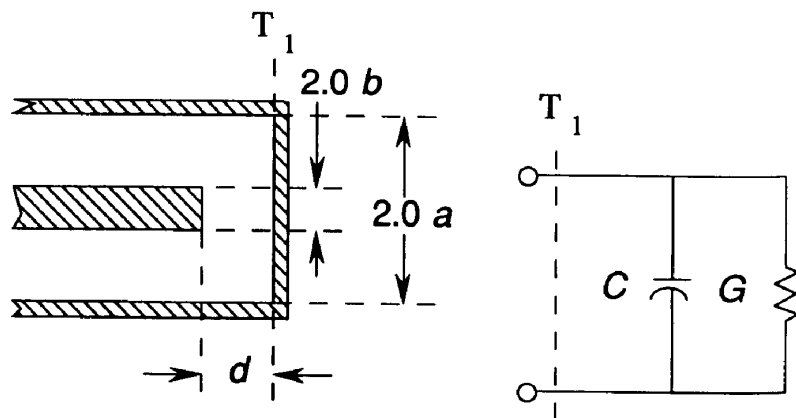


Figure 5.2.6.1: Enclosed Coaxial Open Circuit

¹A 50 Ω line has $a/b = 2.3$ for an air dielectric; a/b increases for other dielectrics.

The equivalent circuit of the coaxial open into its shield is a capacitance from the center conductor to the shield at the reference plane T_1 . The normalized susceptance is [2]

$$\frac{B}{Y_0} = \left(\frac{4.0 b}{\lambda_0} \ln \frac{a}{b} \right) \left[\frac{\pi b}{4.0 d} + \ln \left(\frac{a-b}{d} \right) \right] \quad (5.2.6.1)$$

where

λ_0 = free-space wavelength

This model is valid for:

$$2 \pi d \ll \lambda_0$$

$$d \ll a - b$$

$$\lambda_0 > 2(a - b)$$

REFERENCES

- [1] Gwarek, Wojciech K., "Computer-Aided Analysis of Arbitrarily Shaped Coaxial Discontinuities," *IEEE Transactions on Microwave Theory and Techniques*, Vol. MTT-36, No. 2, February 1988, pp. 337-342.
- [2] Marcuvitz, N., *Waveguide Handbook*, McGraw-Hill, New York, 1955. Reprint with errata and preface by N. Marcuvitz, Peter Peregrinus Ltd., 1986.
- [3] Young, Leo, ed., *Advances in Microwaves*, Volume 2, Academic Press, New York, 1967.

5.3 PAIRED-LINE DISCONTINUITIES

Analysis of paired-line transmission line discontinuities are not generally included in available literature. They are generally not used in filters and circuits. The designer of these structures may be able to make use of the equations of a similar structure for approximate design. The design's success would then rely on further experimentation.

5.3.1 Open End

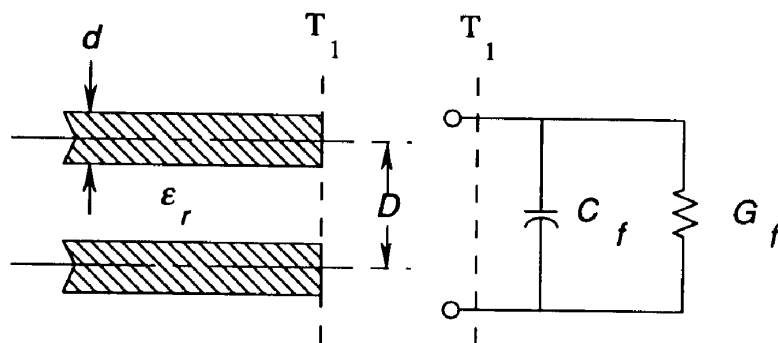


Figure 5.3.1.1: Paired-Line Open End

In [1] the equations for the excess capacitance of the open end of a two-wire transmission line were experimentally determined due to the intractability of an exact solution to be

$$\frac{C_f}{C_0} = \frac{D}{-3.954 + \sqrt{\left[2.564 \cosh^{-1} \left(\frac{D}{d}\right)\right]^2 + (3.954)^2}} \quad (\text{unitless}) \quad (5.3.1.1)$$

where

C_0 is the capacitance per unit length of the transmission line

d and D are defined in Section 3.3.1 and have the same dimensions

which is valid for

$$\frac{C_f}{C_0} \ll \lambda$$

and results from curve-fitting data taken over the range of $1.67 \leq \cosh^{-1} \left(\frac{D}{d}\right) \leq 5.0$. A parallel conductance, G_f , is also present and can be calculated with

$$G_f = \frac{\eta (\beta D)^4 \left[\frac{(C_f / C_0)}{D}\right]^2}{12.0 \pi} \quad (\text{Siemens}) \quad (5.3.1.2)$$

where η is defined in Section 1.1 and the other terms have already been defined.

Green [1] notes that the fringing field lines from the open extend approximately $D/4$ back from the open end and to infinity ($\gg D$) on the open end.

REFERENCES

- [1] Green, H.E., and J.D. Cashman, "End Effect in Open-Circuited Two-Wire Transmission Lines," *IEEE Transactions on Microwave Theory and Techniques*, Vol. MTT-34, No. 1, January 1986, pp. 180–182.

5.4 COPLANAR WAVEGUIDE DISCONTINUITIES

The coplanar waveguide structure is growing in popularity because of its single-sided nature; however closed-form design equations for discontinuities are not yet readily available. Design curves are available in some references. Houdart [1] catalogs many of the structures shown here and others.

REFERENCES

- [1] Houdart, M., "Coplanar Lines: Application to Broadband Microwave Integrated Circuits," *6th European Microwave Conference Proceedings*, 1976, pp. 49–53.

5.4.1 Coplanar Waveguide Shunt Capacitance

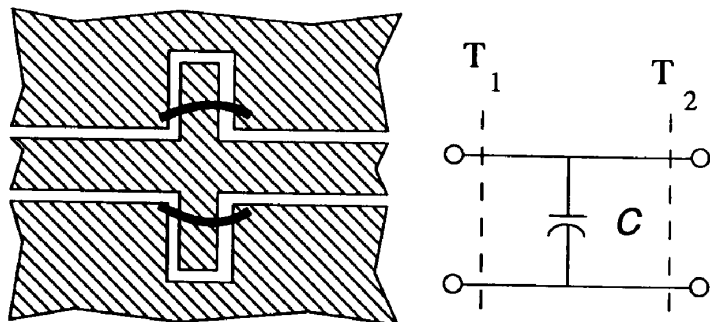


Figure 5.4.1.1: Shunt Capacitance in CPW

This line can be modeled as ideal transmission lines in the absence of a model for the cross junction. The open end effects can be modeled with the aid of curves in [2] and [3] as described later. Multiple lines can be repeated to build up higher capacitances. Bond wires (shown as dark lines in Figure 5.4.1.1) between the grounds are used to suppress undesired modes.

REFERENCES

- [1] Houdart, M., "Coplanar Lines: Application to Broadband Microwave Integrated Circuits," *6th European Microwave Conference Proceedings*, 1976, pp. 49–53.
- [2] Naghed, Mohsen, and Ingo Wolff, "A Three-Dimensional Finite-Difference Calculation of Equivalent Capacitances of Coplanar Waveguide Discontinuities," *1980 IEEE MTT-S International Microwave Symposium Digest*, pp. 1143–1146.

- [3] Simons, Rainee N., and George E. Ponchak, "Modeling of Some Coplanar Waveguide Discontinuities," *IEEE Transactions on Microwave Theory and Techniques*, Vol. MTT-36, No. 12, December 1988, pp. 1796–1803.

5.4.2 Coplanar Waveguide Series Capacitance

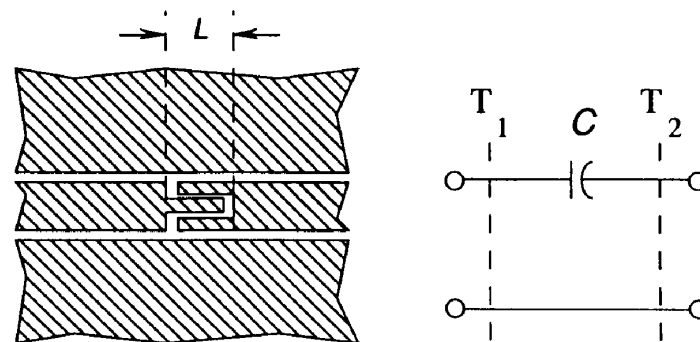


Figure 5.4.2.1: Series Capacitance in CPW

Again, no general closed-form equations are available. Multiple fingers may be interleaved to increase the capacitance. Williams and Schwarz [5] use the above configuration in a bandpass filter design and note that the radiation losses of the above structure are lower than the simpler end-coupled CPW series capacitance. Dib *et al.* [2] give a graph for one set of dimensions. They found that the structure resonated when L was slightly less than $\lambda/4$. In Dib *et al.* [3] closed-form discontinuity model equations are given for a single set of physical dimensions as a function of the finger length L .

In lieu of design equations a starting point may be the relations for coupled CPW as described in Chapter 4 together with a technique after Bates [1].

REFERENCES

- [1] Bates, R.N., "Design of Microstrip Spur-Line Band-Stop Filters," *Microwaves, Optics and Acoustics*, Vol. 1, No. 6, November 1977, pp. 209–214.
- [2] Dib, N.I., *et al.*, "Coplanar Waveguide Discontinuities for P-I-N Diode Switches and Filter Applications," *1990 IEEE MTT-S International Microwave Symposium Digest*, pp. 399–402.
- [3] Dib, N.I., *et al.*, "Theoretical and Experimental Characterization of Coplanar Waveguide Discontinuities for Filter Applications," *IEEE Transactions on Microwave Theory and Techniques*, Vol. MTT-39, No. 5, May 1991, pp. 873–882.

- [4] Houdart, M., "Coplanar Lines: Application to Broadband Microwave Integrated Circuits," *6th European Microwave Conference Proceedings*, 1976, pp. 49–53.
- [5] Williams, Dylan F., and S.E. Schwarz, "Design and Performance of Coplanar Waveguide Bandpass Filters," *IEEE Transactions on Microwave Theory and Techniques*, Vol. MTT-31, No. 7, July 1983, pp. 558–566.

5.4.3 Coplanar Waveguide Shunt Inductance

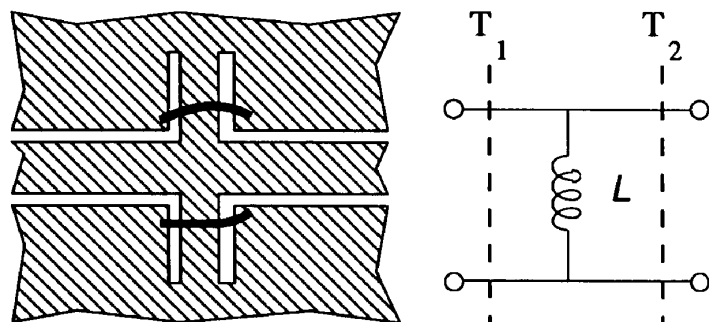


Figure 5.4.3.1: Shunt Inductance in CPW

This structure is cataloged in Houdart [1]. Kibuuka *et al.* [2] use this in a snaked or bent form as part of a band-pass filter. No cross model is available, so ideal equations for the short-circuited lines must be used. The bond wires (shown as dark lines in Figure 5.4.3.1) are used to suppress higher order modes.

REFERENCES

- [1] Houdart, M., "Coplanar Lines: Application to Broadband Microwave Integrated Circuits," *6th European Microwave Conference Proceedings*, 1976, Rome, Italy, pp. 49–53.
- [2] Kibuuka, Godfrey, *et al.*, "Coplanar Lumped Elements and their Application in Filters on Ceramic and Gallium Arsenide Substrates," *19th European Microwave Conference Proceedings*, September 4–7, 1989, London, England, pp. 656–661.

5.4.4 Coplanar Waveguide Series Inductance

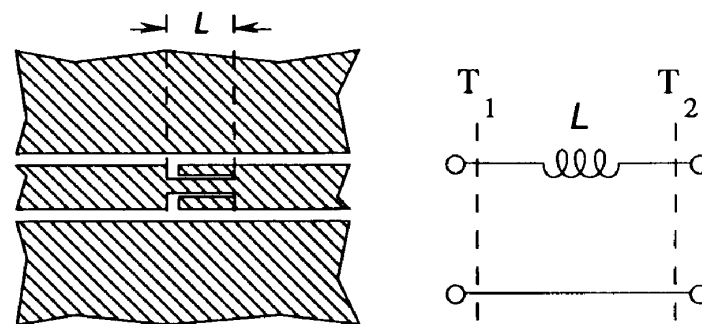


Figure 5.4.4.1: Series Inductance in CPW

In Dib *et al.* [2], this structure is shown to resonate somewhat below the frequency where $L = \lambda / 4.0$. In Dib *et al.* [3] closed-form discontinuity model equations are given for a single set of physical dimensions as a function of the finger length L . In lieu of a full set of design equations a starting point may be the relations for coupled CPW as described in Chapter 4 together with a technique after Bates [1].

REFERENCES

- [1] Bates, R.N., "Design of Microstrip Spur-Line Band-Stop Filters," *Microwaves, Optics and Acoustics*, Vol. 1, No. 6, November 1977, pp. 209–214.
- [2] Dib, N.I., *et al.*, "Coplanar Waveguide Discontinuities for P-I-N Diode Switches and Filter Applications," *1990 IEEE MTT-S International Microwave Symposium Digest*, pp. 399–402.
- [3] Dib, N.I., *et al.*, "Theoretical and Experimental Characterization of Coplanar Waveguide Discontinuities for Filter Applications," *IEEE Transactions on Microwave Theory and Techniques*, Vol. MTT-39, No. 5, May 1991, pp. 873–882.
- [4] Houdart, M., "Coplanar Lines: Application to Broadband Microwave Integrated Circuits," *6th European Microwave Conference Proceedings*, 1976, Rome, Italy, pp. 49–53.
- [5] Houdart, M., *et al.*, "Coplanar Lines: Application to Lumped and Semi-Lumped Microwave Integrated Circuits," *7th European Microwave Conference Proceedings*, 1977, Copenhagen, Denmark, pp. 450–454.

- [6] Simons, Rainee N., and George E. Ponchak, "Modeling of Some Coplanar Waveguide Discontinuities," *IEEE Transactions on Microwave Theory and Techniques*, Vol. MTT-36, No. 12, December 1988, pp. 1796–1803.

5.4.5 Coplanar Waveguide Open End with Connected Grounds

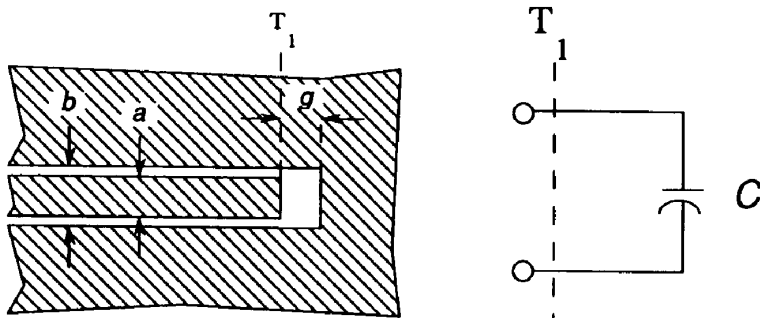


Figure 5.4.5.1: CPW Open End with Connected Grounds

The coplanar waveguide open end occurs in filters and matching circuits. Curves of the capacitance are found in Naghed and Wolff [1] and Simons and Ponchak [3] for a limited number of cases. For $g/h \geq 0.3$ the capacitance is approximately independent of g . Curve-fitting the results of Naghed and Wolff [2] gives the following equations for design. For $a/b = 0.75$:

$$C = 13.49807 + \frac{0.4487826}{g/h} - \frac{0.0047486308}{(g/h)^2} \quad (\text{fF}) \quad (5.4.5.1)$$

For $a/b = 0.50$:

$$C = 9.7179749 + \frac{0.34173829}{g/h} - \frac{0.0042329111}{(g/h)^2} \quad (\text{fF}) \quad (5.4.5.2)$$

For $a/b = 0.33$:

$$C = 7.3706512 + \frac{0.20880575}{g/h} - \frac{0.002343378}{(g/h)^2} \quad (\text{fF}) \quad (5.4.5.3)$$

For $a/b = 0.25$:

$$C = 5.9622775 + \frac{0.16905027}{g/h} - \frac{0.0019673178}{(g/h)^2} \quad (\text{fF}) \quad (5.4.5.4)$$

Equations are valid to approximately $\pm 2.5\%$ over the range of the graphs

$$0.03 \leq g/h \leq 1.0$$

Interpolation between values of a/b can be used elsewhere.

REFERENCES

- [1] Naghed, Mohsen, and Ingo Wolff, "A Three-Dimensional Finite-Difference Calculation of Equivalent Capacitances of Coplanar Waveguide Discontinuities," *1980 IEEE MTT-S International Microwave Symposium Digest*, pp. 1143–1146.
- [2] Naghed, Mohsen, and Ingo Wolff, "Equivalent Capacitances of Coplanar Waveguide Discontinuities and Interdigitated Capacitors Using a Three-Dimensional Finite Difference Method," *IEEE Transactions on Microwave Theory and Techniques*, Vol. MTT-38, No. 12, December 1990, pp. 1808–1815.
- [3] Simons, Rainee N., and George E. Ponchak, "Modeling of Some Coplanar Waveguide Discontinuities," *IEEE Transactions on Microwave Theory and Techniques*, December 1988, pp. 1796–1803.

5.4.6 CPW Open End

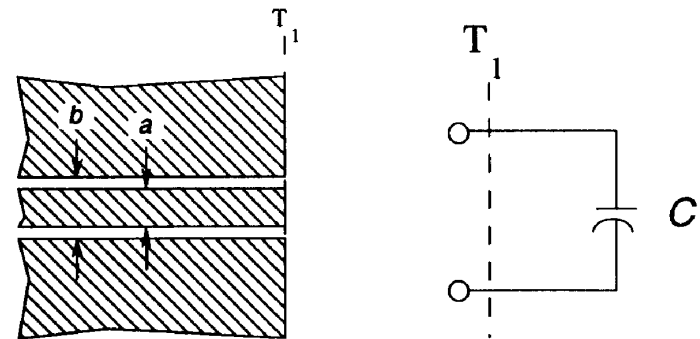


Figure 5.4.6.1: Coplanar Waveguide Open End

Analysis techniques are described in [1]. Design curves are available in Naghed and Wolff [2]. For the case of $\epsilon_r = 9.8$ and $h = 0.635$ mm, curve fits to their curves are given below. These curves are valid to within approximately $\pm 5\%$ over the range

$$0.1 \leq w/d \leq 0.9$$

For $h/b = 1.0$:

$$C = -5.3001726 + 33.062752 \sqrt{a/b} \quad (\text{fF}) \quad (5.4.6.1)$$

For $h/b = 0.2$:

$$C = -1.4051152 + 22.07531 \sqrt{a/b} \quad (\text{fF}) \quad (5.4.6.2)$$

For $h/b = 0.1$:

$$C = -0.99986182 + 17.285616 \sqrt{a/b} \quad (\text{fF}) \quad (5.4.6.3)$$

REFERENCES

- [1] Jansen, R.H., "Hybrid Mode Analysis of End Effects of Planar Microwave and Millimetrewave Transmission Lines," *IEE Proceedings*, Vol. 128, Pt. H, No. 2, April 1981, pp. 77–86.
- [2] Naghed, Mohsen, and Ingo Wolff, "A Three-Dimensional Finite-Difference Calculation of Equivalent Capacitances of Coplanar Waveguide Discontinuities," *1980 IEEE MTT-S International Microwave Symposium Digest*, pp. 1143–1146.
- [3] Naghed, Mohsen, and Ingo Wolff, "Equivalent Capacitances of Coplanar Waveguide Discontinuities and Interdigitated Capacitors Using a Three-Dimensional Finite Difference Method," *IEEE Transactions on Microwave Theory and Techniques*, Vol. MTT-38, No. 12, December 1990, pp. 1808–1815.

5.4.7 Coplanar Waveguide Series Gap

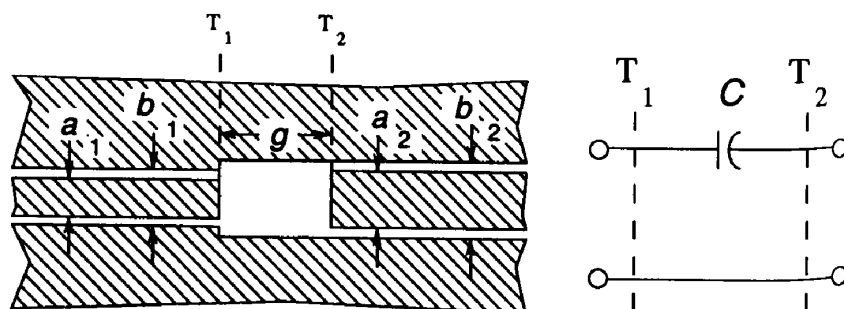


Figure 5.4.7.1: CPW Series Gap

Williams and Schwarz [4] compare this structure to simple interdigital CPW capacitors (as in Figure 5.4.2.1). They also experimented with overlaying this structure with dielectric material and dielectric material with a conductive mode suppressing ribbon.

Some design curves are also given for the equivalent circuit. Simons and Ponchak [3] include shunt capacitances in their model and also present curves.

Naghd [1] gives curves for the case of $\epsilon_r = 9.8$, $h = 0.635$, $a_1/h = 0.2$, and $a_1/b_1 = a_2/b_2 = 0.56$. Curve fits are given below.

For $a_2/a_1 = 1.0$:

$$C_{p1} = C_{p2} = \frac{-0.47822096 + 57.893523 \frac{g}{h}}{1.0 + 5.8451737 \frac{g}{h}} \quad (\text{fF}) \quad (5.4.7.1)$$

$$C_g = \frac{14.099684 - 11.843073 \frac{g}{h} + 1.0433114 \left(\frac{g}{h}\right)^2}{1.0 + 10.502696 \frac{g}{h} + 18.503307 \left(\frac{g}{h}\right)^2} \quad (\text{fF}) \quad (5.4.7.2)$$

For $a_2/b_1 = 2.0$:

$$C_{p1} = \frac{-4.0235493 + 35.576495 \frac{g}{h}}{1.0 + 2.7541323 \frac{g}{h}} \quad (\text{fF}) \quad (5.4.7.3)$$

$$C_{p2} = \frac{1.0 - 7.9340513 \frac{g}{h} + 37.923732 \left(\frac{g}{h}\right)^2}{0.087137596 - 0.59125367 \frac{g}{h} + 2.941263 \left(\frac{g}{h}\right)^2} \quad (\text{fF}) \quad (5.4.7.4)$$

$$C_g = \frac{18.99756 - 1.8997741 \frac{g}{h}}{1.0 + 9.9022677 \frac{g}{h} + 6.8034362 \left(\frac{g}{h}\right)^2} \quad (\text{fF}) \quad (5.4.7.5)$$

For $a_2/b_1 = 3.0$:

$$C_{p1} = \frac{-6.2682507 + 32.193198 \frac{g}{h}}{1.0 + 3.6295239 \frac{g}{h}} \quad (\text{fF}) \quad (5.4.7.6)$$

$$C_{p2} = \sqrt{315.94893 + 9.6539534/x} \quad (5.4.7.7)$$

$$C_g = \frac{20.102184 - 2.00947 \frac{g}{h}}{1.0 + 7.7736146 \frac{g}{h}} \quad (\text{fF}) \quad (5.4.7.8)$$

It was attempted to make these equations correct asymptotically; however, the accuracy and range of the data made this impossible to do in every case. These CAE equations are valid in general to within < 5% over the range

$$0.02 \leq g/h \leq 0.3$$

REFERENCES

- [1] Drissi, M'hamed, *et al.*, "Analysis of Coplanar Waveguide Radiating End Effects Using the Integral Equation Technique," *IEEE Transactions on Microwave Theory and Techniques*, Vol. MTT-39, No. 1, January 1991.
- [2] Naghed, Mohsen, and Ingo Wolff, "A Three-Dimensional Finite-Difference Calculation of Equivalent Capacitances of Coplanar Waveguide Discontinuities," *1980 IEEE MTT-S International Microwave Symposium Digest*, pp. 1143–1146.
- [3] Naghed, Mohsen, and Ingo Wolff, "Equivalent Capacitances of Coplanar Waveguide Discontinuities and Interdigitated Capacitors Using a Three-Dimensional Finite Difference Method," *IEEE Transactions on Microwave Theory and Techniques*, Vol. MTT-38, No. 12, December 1990, pp. 1808–1815.
- [4] Simons, Rainee N., and George E. Ponchak, "Modeling of Some Coplanar Waveguide Discontinuities," *IEEE Transactions on Microwave Theory and Techniques*, Vol. MTT-36, No. 12, December 1988, pp. 1796–1803.
- [5] Williams, Dylan F., and S.E. Schwarz, "Design and Performance of Coplanar Waveguide Bandpass Filters," *IEEE Transactions on Microwave Theory and Techniques*, Vol. MTT-31, No. 7, July 1983, pp. 558–566.

5.4.8 Coplanar Waveguide Abrupt Width Change

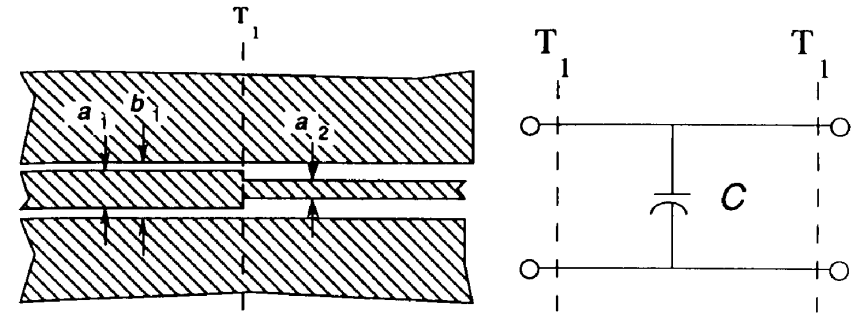


Figure 5.4.8.1: CPW Abrupt Width Change

A limited number of curves for this discontinuity are given in Simons and Ponchak [1]. Curve fits to this data are given here.

For $\epsilon_r = 2.1$, $a_1 = 10.5$ mil, $a_1/b_1 = 0.86$, $h = 62$ mil, and $1.1 < S_1/S_2 \leq 2.0$:

$$C_s = 0.01335 + 0.00772 \left(\frac{S_1}{S_2} \right) \quad (5.4.8.1)$$

For $\epsilon_r = 2.2$, $a_1 = 11$ mil, $a_1/b_1 = 0.9$, $h = 125$ mil, and $1.1 < S_1/S_2 \leq 3.0$:

$$C_s = \sqrt{0.0062268592 - \frac{0.0048471339}{(S_1/S_2)^2}} \quad (5.4.8.2)$$

For $\epsilon_r = 2.2$, $a_1 = 9$ mil, $a_1/b_1 = 0.9$, $h = 125$ mil, and $1.1 < S_1/S_2 \leq 3.0$:

$$C_s = \sqrt{0.0024722655 - \frac{0.0023832659}{(S_1/S_2)^2}} \quad (5.4.8.3)$$

These equations fit the curves of [1] over the specified range to better than $\pm 6\%$.

REFERENCES

- [1] Simons, Rainee N., and George E. Ponchak, "Modeling of Some Coplanar Waveguide Discontinuities," *IEEE Transactions on Microwave Theory and Techniques*, Vol. MTT-36, No. 12, December 1988, pp. 1796–1803.

5.4.9 CPW Abrupt Ground Plane Width Change

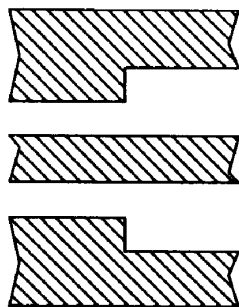


Figure 5.4.9.1: Coplanar Waveguide Abrupt Ground Plane Step

There is very little information on this structure in the literature. Kuo and Itoh [1] give plots of s -parameters for a range of step sizes.

REFERENCE

- [1] Kuo, Chih-Wen, and Tatsuo Itoh, "Characterization of the Coplanar Waveguide Step Discontinuity Using the Transverse Resonance Method," *19th European Microwave Conference Proceedings*, September 4–7, 1989, London, England, pp. 662–665.

5.4.10 Abrupt Coplanar Waveguide Bend

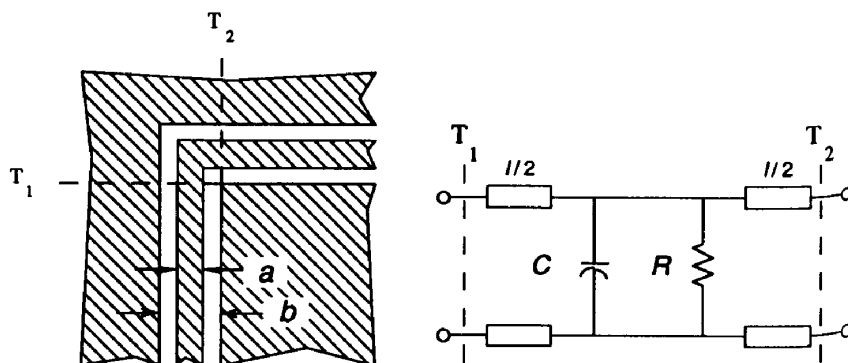


Figure 5.4.10.1: CPW Abrupt Bend

The capacitive discontinuity at a bend in CPW is caused by the difference in path length between the inner and outer metallization slots. Simons and Ponchak [1] report very good results compensating this difference with a small piece of dielectric on the inner slot of the bend to slow down this part of the wave. The model is described and was measured for one set of parameters in [2]. The resistive component is due to radiation losses at the corner.

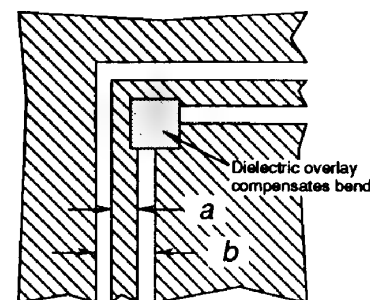


Figure 5.4.10.2: Compensated Abrupt Bend in CPW

REFERENCES

- [1] Simons, Rainee N., and George E. Ponchak, "Modeling of Some Coplanar Waveguide Discontinuities," *IEEE Transactions on Microwave Theory and Techniques*, Vol. MTT-36, No. 12, December 1988, pp. 1796–1803.
- [2] Simons, Rainee N., *et al.*, "Channelized Coplanar Waveguide: Discontinuities, Junctions, and Propagation Characteristics," *1989 IEEE MTT-S Symposium Digest*, pp. 915–918.

5.4.11 CPW Short Circuit

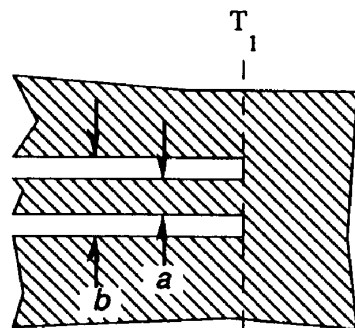


Figure 5.4.11.1: CPW Shorted End

The CPW short is not well covered in the literature. It can be modeled as an ideal line without discontinuities. A similar structure, coupled slots, is analyzed in [3].

REFERENCES

- [1] Bartolucci, G., and J. Piotrowski, "Full-Wave Analysis of Shielded Coplanar Waveguide Short-End," *Electronics Letters*, Vol. 26, No. 19, September 13, 1990, pp. 1615–1616.
- [2] Drissi, M'hamed, *et al.*, "Analysis of Coplanar Waveguide Radiating End Effects Using the Integral Equation Technique," *IEEE Transactions on Microwave Theory and Techniques*, Vol. MTT-39, No. 1, January 1991.
- [3] Jansen, R.H., "Hybrid Mode Analysis of End Effects of Planar Microwave and Millimetrewave Transmission Lines," *IEE Proceedings*, Vol. 128, Pt. H, No. 2, April 1981, pp. 77–86.

5.4.12 CPW Air Bridge

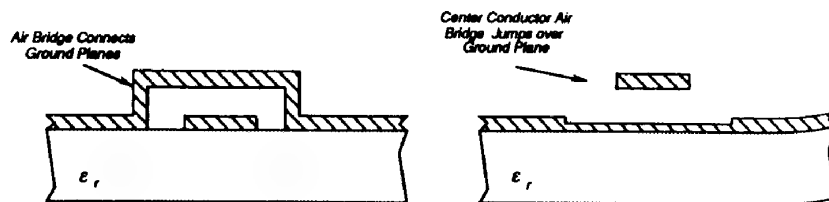


Figure 5.4.12.1: CPW Air Bridges

The air bridge is used in MMIC CPW to bridge signal conductors and also to break up modes by shorting the ground planes together. Because of the small dimensions involved in MMICs and multiple crossings, the capacitance between the lines becomes significant and must be corrected. Koster [1] measured the two types of bridges and various compensation schemes to reduce the capacitance's effects.

The second configuration of Figure 5.4.12.1 was shown to have significantly lower losses. Koster also recommends a similar MMIC coplanar waveguide T-junction as having fewer problems with higher order modes.

REFERENCE

- [1] Koster, N.H.L., *et al.*, *19th European Microwave Conference Proceedings*, September 4–7, 1989, London, England, pp. 666–671.

5.4.13 CPW Interdigital Capacitor

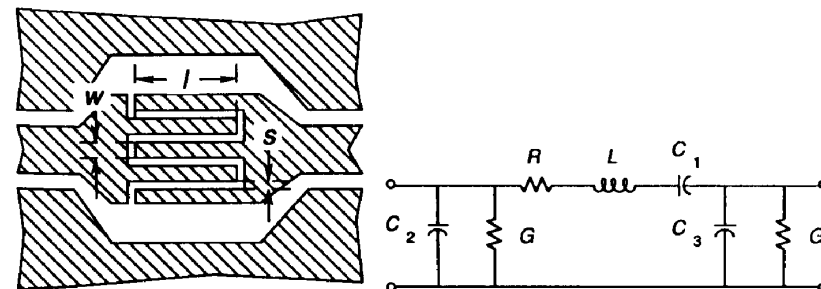


Figure 5.4.13.1: CPW Interdigital Capacitor

Design equations are not available for this configuration. Kibuuka *et al.* [1] found that the ϵ_{eff} of the modes of this structure are nearly equal and gives the model shown in Figure 5.4.13.1 for one set of dimensions. Naghed and Wolff [2] calculated this structure using the finite difference technique. The inductance element of the model, L , becomes significant for long fingers.

The interdigital capacitor can be analyzed by coupled lines with the correct conditions set at the ends. As an alternative, the interdigital capacitor of Chapter 7, could be used.

REFERENCES

- [1] Kibuuka, Godfrey, *et al.*, "Coplanar Lumped Elements and their Application in Filters on Ceramic and Gallium Arsenide Substrates," *19th European Microwave Conference Proceedings*, September 4–7, 1989, London, England, pp. 656–661.

- [2] Naghed, Mohsen, and Ingo Wolff, "A Three-Dimensional Finite-Difference Calculation of Equivalent Capacitances of Coplanar Waveguide Discontinuities," *1980 IEEE MTT-S International Microwave Symposium Digest*, pp. 1143–1146.
- [3] Naghed, Mohsen, and Ingo Wolff, "Equivalent Capacitances of Coplanar Waveguide Discontinuities and Interdigitated Capacitors Using a Three-Dimensional Finite Difference Method," *IEEE Transactions on Microwave Theory and Techniques*, Vol. MTT-38, No. 12, December 1990, pp. 1808–1815.

5.5 MICROSTRIP LINE DISCONTINUITIES

5.5.1 Microstrip Line Bend

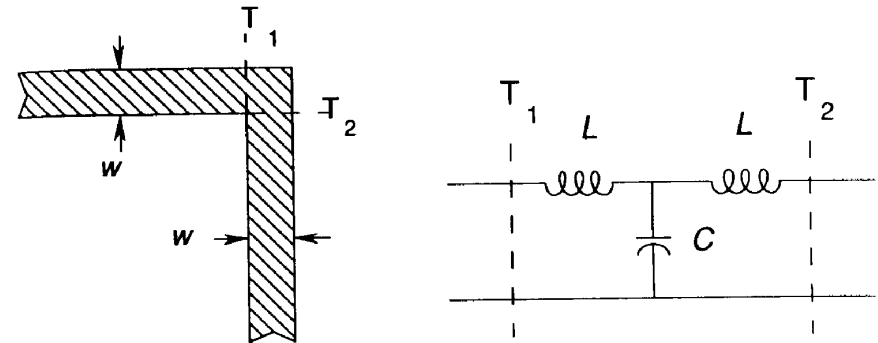


Figure 5.5.1.1: Microstrip Line Abrupt Bend

The discontinuity created by a bend in a microstrip line is given by the model above [2] where:

$$L/h = 100.0 (4.0 \sqrt{w/h} - 4.21) \text{ (nH/m)} \quad (5.5.1.1)$$

and for $w/h < 1$:

$$C/w = \frac{(14.0 \epsilon_r + 12.5)(w/h) - (1.83 \epsilon_r - 2.25)}{\sqrt{w/h}} + \frac{0.02 \epsilon_r}{w/h} \text{ (pF/m)} \quad (5.5.1.2)$$

or for $w/h \geq 1$:

$$C/w = (9.5 \epsilon_r + 1.25)(w/h) + 5.2 \epsilon_r + 7.0 \quad (5.5.1.3)$$

These equations are accurate to within 5% for

$$2.5 \leq \epsilon_r \leq 15.0$$

$$0.1 \leq w/h \leq 5.0$$

Kirschning [7] gives equations for C and L as

$$C = 0.001 h \left[(10.35 \epsilon_r + 2.5) (w/h)^2 + (2.6 \epsilon_r + 5.44) (w/h) \right] \text{ (pF)} \quad (5.5.1.4)$$

$$L = 0.22 \left[1.0 - 1.35 e^{-0.18(w/h)^{1.39}} \right] \text{ (nH)} \quad (5.5.1.5)$$

where w and h are in mm. Equations are valid for

$$2.0 \leq \epsilon_r \leq 13.0$$

$$0.2 \leq w/h \leq 6.0$$

Agreement of calculated and measured resonant frequencies using these equations was to 0.3%.

REFERENCES

- [1] Douville, Rene J.P., and David S. James, "Experimental Study of Microstrip Bends and Their Compensation," *IEEE Transactions on Microwave Theory and Techniques*, Vol. MTT-26, No. 3, March 1978, pp. 175-181.
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- [3] Hammerstad, E., "Computer-Aided Design of Microstrip Couplers with Accurate Discontinuity Models," *1981 IEEE MTT-S International Microwave Symposium Digest*, pp. 54-56.
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- [5] Hill, A., and V.K. Tripathi, "Analysis and Modeling of Coupled Right Angle Microstrip Bend Discontinuities," *1989 IEEE MTT-S Symposium Digest*, pp. 1143-1146.
- [6] Hill, Achim, "An Efficient Algorithm of the Analysis of Passive Microstrip Discontinuities for Microwave and Millimeter Wave Integrated Circuits in a Shielding Box," *Compact Software*, 1990, pp. 9-10.
- [7] Kirschning, M., et al., "Measurement and Computer-Aided Modeling of Microstrip Discontinuities by an Improved Resonator Method," *IEEE Microwave Theory and Techniques Symposium Digest 1983*, pp. 495-497.

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- [10] Menzel, Wolfgang and Ingo Wolff, "A Method for Calculating the Frequency-Dependent Properties of Microstrip Discontinuities," *IEEE Transactions on Microwave Theory and Techniques*, Vol. MTT-25, No. 2, February 1977, pp. 107-112.
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5.5.2 Optimal Right-Angle Mitered Bend

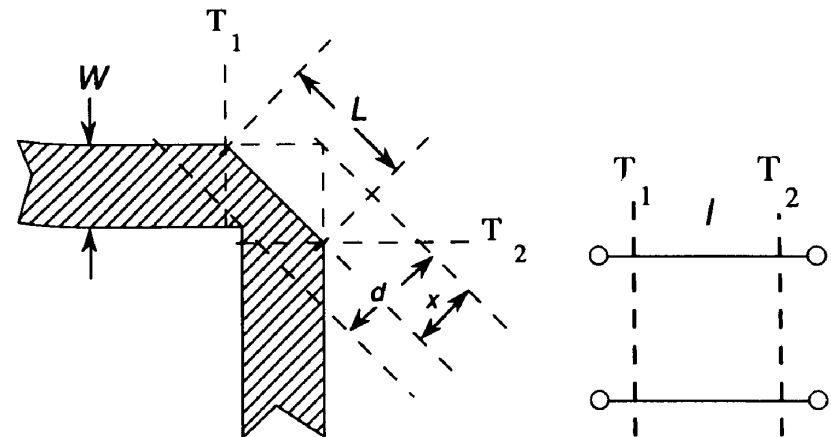


Figure 5.5.2.1: Optimal Microstrip Line Mitered Bend

The optimal mitre occurs for:

$$M = 52 + 65 e^{-1.35 w/h} = \frac{100x}{d} \quad (\%) \quad (5.5.2.1)$$

or, more conveniently,

$$\frac{L}{w} = \sqrt{2} [1.04 + 1.3 e^{-1.35 w/h}] \quad (5.5.2.2)$$

$$d - x = \sqrt{2} w - \frac{L}{2} \quad (5.5.2.3)$$

These equations are valid for $w/h \geq .25$, $\epsilon_r \leq 25$, ± 4 (sic, %?) accuracy. In Chadha [2], the optimal mitre is found to be $w/h = 2.895$. Chadha also recommends that for bends made of lines of unequal widths that more material be removed from the wider line. Figure 5.5.3.2 compares the references.

The equivalent electrical length of the mitre, l , is calculable from arbitrary bend equations of [1] by setting the angle α to 90° resulting in

$$l = \frac{L}{\sqrt{2}} \quad (5.5.2.4)$$

Other techniques of equalizing the bend are shown in Figure 5.5.2.2 below. Analytic equations for these structures are not available. Douville [3] shows that the 45° double bend is always better than a single right-angle bend.

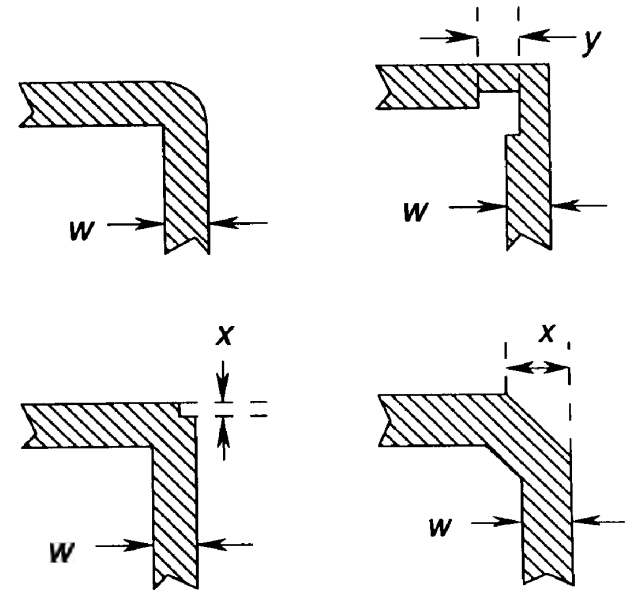


Figure 5.5.2.2: Various Techniques of Bend Compensation

REFERENCES

- [1] Anders, Peter, and Fritz Arndt, "Microstrip Discontinuity Capacitances and Inductances for Double Steps, Mitred Bends with Arbitrary Angle, and Asymmetric Right-Angle Bends," *IEEE Transactions on Microwave Theory and Techniques*, Vol. MTT-28, No. 11, November 1988, pp. 1213–1217.
- [2] Chadha, Rakesh, and K.C. Gupta, "Compensation of Discontinuities in Planar Transmission Lines," *IEEE Transactions on Microwave Theory and Techniques*, Vol. MTT-30, No. 12, December 1982, pp. 2151–2156.
- [3] Douville, Rene J.P., and David S. James, "Experimental Study of Microstrip Bends and Their Compensation," *IEEE Transactions on Microwave Theory and Techniques*, Vol. MTT-26, No. 3, March 1978, pp. 175–181.
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- [8] Zheng, J.X., and David C. Chang, "Numerical Modelling of Chamfered Bends and Other Microstrip Junctions of General Shape in MMICs," *1990 IEEE MTT-S Symposium Digest*, pp. 709–712.

5.5.3 Optimal Arbitrary Angle Mitered Bend

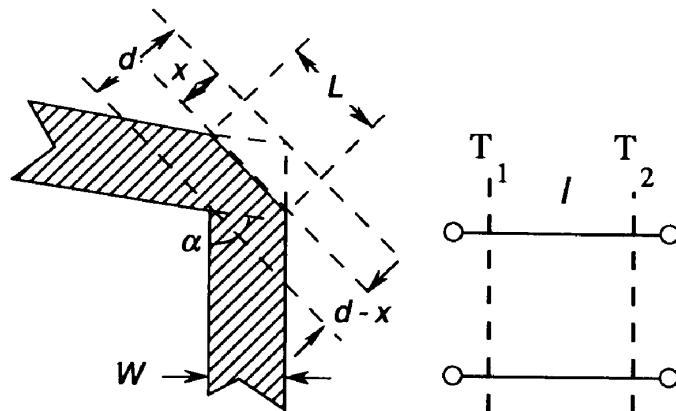


Figure 5.5.3.1: Arbitrary Mitered Bend

The optimal mitre dimensions are shown in Figure 5.5.3.2 for a range of angles and w/h ratios. This figure combines the results of a number of references for comparison and convenience.

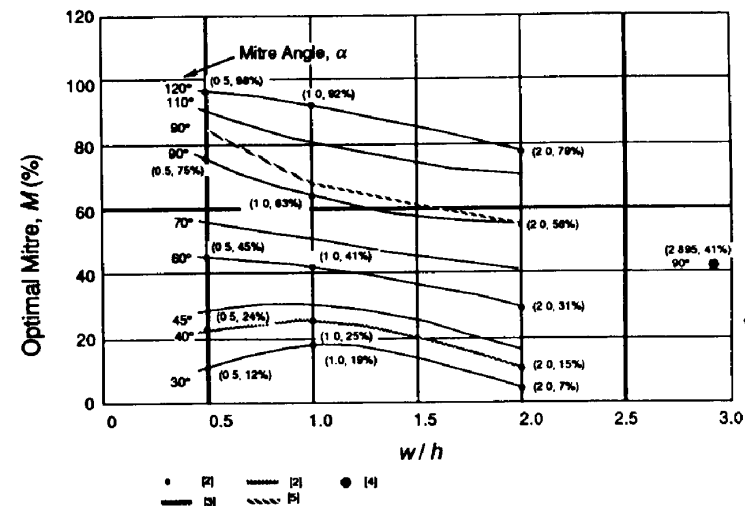


Figure 5.5.3.2: Optimum Mitres for Arbitrary Bend Angles

$$M = \frac{100x}{d} \quad (\%) \quad (5.5.3.1)$$

$$d - x = \sqrt{2} w \left(1.0 - \frac{M}{100} \right) \quad (5.5.3.2)$$

The arbitrary angle bend can be compensated by taking material off of the outside of the corner. This can be done by mitering. Figure 5.5.3.2 shows the optimal mitre percentage to achieve a matched impedance as a function of the line's w/h ratio. The accuracy is expected to be within a few percent [3]. For comparison the equation of [5] is plotted for the optimal 90° mitered bend. The other references are also shown.

The electrical length of the bend is calculated:

$$\frac{l_k}{w} = \frac{2M}{100 \sin \alpha} \quad (5.5.3.3)$$

REFERENCES

- [1] Anders, Peter, and Fritz Arndt, "Beliebig abgeknickte Mikrostrip-Leitungen mit bogenförmigen Übergang," *Archiv Für Elektronik und Übertragungstechnik*, Band 33, Heft 3, March 1979, pp. 93–99.

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- [3] Anders, Peter, and Fritz Arndt, "Microstrip Discontinuity Capacitances and Inductances for Double Steps, Mitred Bends with Arbitrary Angle, and Asymmetric Right-Angle Bends," *IEEE Transactions on Microwave Theory and Techniques*, Vol. MTT-28, No. 11, November 1980, pp. 1213–1217.
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5.5.4 Microstrip Line Rounded Bend

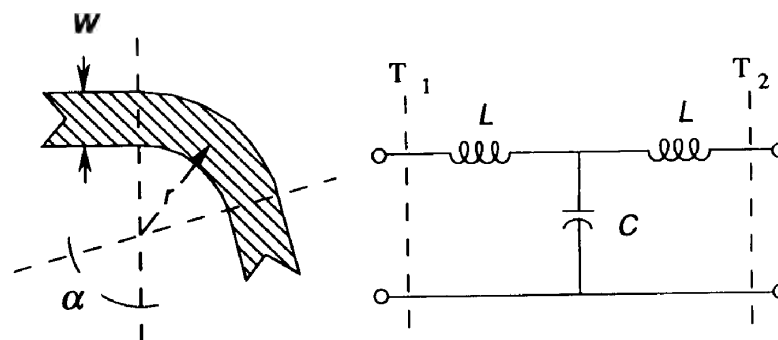


Figure 5.5.4.1: Microstrip Line Rounded Bend

Curves of model parameters are available in Weissnar *et al.* [2]; however, no closed-form equations have been derived to date. The dispersion relative to a straight line has been derived by curve-fitting experimental results in Roy *et al.* [1] and are:

$$\epsilon_{eff, bend}(f) = \epsilon_r - \frac{\epsilon_r - \epsilon_{eff}(0) \exp\left(\frac{4\sqrt{h/r}}{12.0}\right)}{1.0 + \left(\frac{f}{f_c}\right)^2} \quad (5.5.4.1)$$

where

$$f_c = \frac{c}{4.0 h \sqrt{\epsilon_r - 1.0}} \quad (5.5.4.2)$$

$\epsilon_{eff}(0)$ is defined in (3.5.1.3)

r is the mean radius of the curve

which is valid for

$$2.0 \leq \epsilon_r \leq 12.0$$

$$0.8 \leq w/h \leq 5.0$$

$$\frac{\lambda_g}{4.0 \pi h} < r/h \leq \infty$$

$$0.0 \leq f \leq 12 \text{ GHz}$$

This equation showed good agreement with dispersion equations of a straight line when r is infinity.

REFERENCES

- [1] Roy, Jibendu Sekhar, *et al.*, "Dispersion Characteristics of Curved Microstrip Transmission Lines," *IEEE Transactions on Microwave Theory and Techniques*, Vol. MTT-38, No. 8, August 1990, pp. 1366–1370.
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5.5.5 Optimal Microstrip Line Rounded Bend

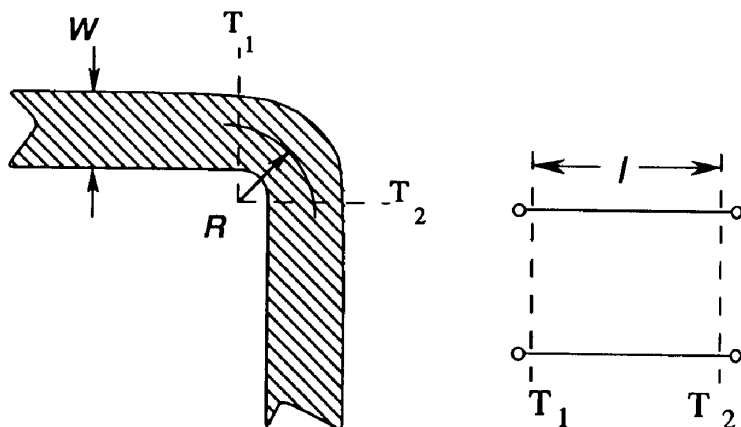


Figure 5.5.5.1: Optimal Microstrip Line Rounded Bend

As an alternative to the mitred corners of the previous sections, a microstrip line may be gradually bent to make a corner. In [2] it is stated that $R \geq 4.0 h$ will give $VSWR < 1.05$. Generally, this is not a space efficient technique to use. U-bends can be made in a similar fashion as shown in Figure 5.5.5.2 below.

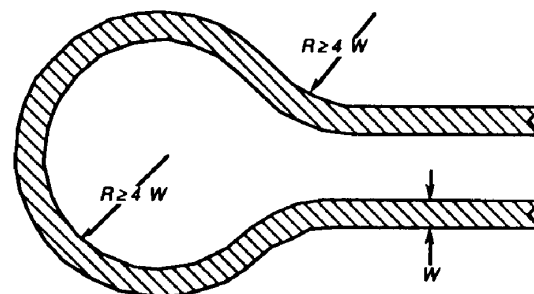


Figure 5.5.5.2: Microstrip Line U-bend

The line can be modeled as a section of line having length equal to the center line length of the curved section. This model does not include the changes in ϵ_{eff} and the effects of coupling between parts of the curve.

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5.5.6 Unequal Width Bend

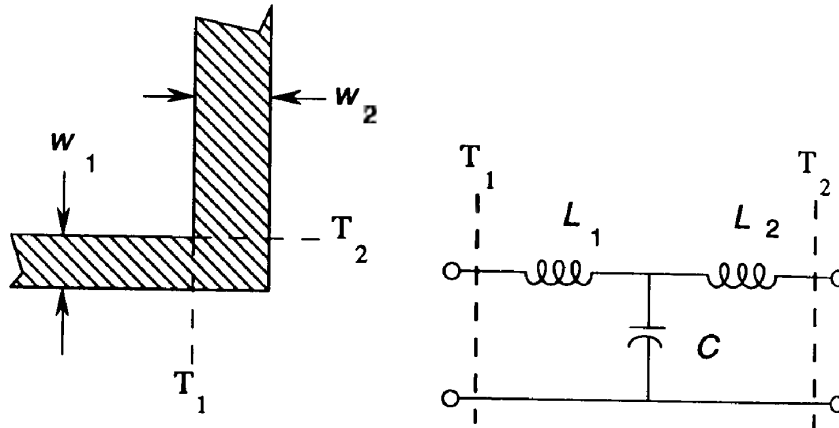


Figure 5.5.6.1: Unequal Width Bend

Graphs of s parameters of this junction are available in [1] and [2]. The curves of [2] have been fit with closed-form equations. These can be interpolated for other values of w_1/h . For this model (low frequency) $L_1 = L_2$.

For $w_1/h = 2.0$,

$$\frac{C}{C'h} = 0.23073798 + 1.5514797 \left(\frac{w_2}{w_1}\right) - 0.052050577 \left(\frac{w_2}{w_1}\right)^2 \quad (5.5.6.1)$$

$$\frac{2.0 L}{L'h} = \frac{0.84199442 - 1.4452286 \left(\frac{w_2}{w_1}\right) + 2.3264141 \left(\frac{w_2}{w_1}\right)^2}{1.0 - 0.79920542 \left(\frac{w_2}{w_1}\right) + 2.7511988 \left(\frac{w_2}{w_1}\right)^2} \quad (5.5.6.2)$$

For $w_1/h = 1.0$,

$$\frac{C}{C'h} = 0.23550754 + 0.41160143 \left(\frac{w_2}{w_1}\right) + 0.21758643 \left(\frac{w_2}{w_1}\right)^2 \quad (5.5.6.3)$$

$$\frac{2.0 L}{L'h} = \frac{0.27972652 - 1.0950536 \left(\frac{w_2}{w_1}\right) + 1.4761673 \left(\frac{w_2}{w_1}\right)^2}{1.0 + 1.0158358 \left(\frac{w_2}{w_1}\right) + 4.9673273 \left(\frac{w_2}{w_1}\right)^2} \quad (5.5.6.4)$$

For $w_1/h = 0.5$,

$$\frac{C}{C'h} = 0.15274236 - 0.033719111 \left(\frac{w_2}{w_1}\right) - 0.23887456 \left(\frac{w_2}{w_1}\right)^2 \quad (5.5.6.5)$$

$$\frac{2.0 L}{L'h} = \frac{0.4550564 - 3.1343404 \left(\frac{w_2}{w_1}\right) + 2.3297251 \left(\frac{w_2}{w_1}\right)^2}{1.0 + 8.2371314 \left(\frac{w_2}{w_1}\right) + 4.1304198 \left(\frac{w_2}{w_1}\right)^2} \quad (5.5.6.6)$$

The variables C' and L' are the incremental capacitance and inductance of the single strip. These equations fit the curves to better than approximately 2.5% for the valid region of

$$0.1 \leq w_2/w_1 \leq 0.9$$

The original curve accuracy is a few percent. Frequency dependency of the parameters is also available in Mehran [3].

REFERENCES

- [1] Anders, Peter, and Fritz Arndt, "Microstrip Discontinuity Capacitances and Inductances for Double Steps, Mitred Bends with Arbitrary Angle, and Asymmetric Right-Angle Bends," *IEEE Transactions on Microwave Theory and Techniques*, Vol. MTT-28, No. 11, November 1988, pp. 1213–1217.
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5.5.7 Microstrip Line Radial Stub

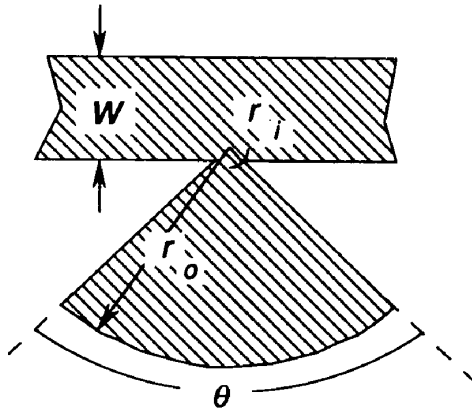


Figure 5.5.7.1: Radial Stub

This element presents a low impedance to ground at a precisely located point. It is useful because it is physically shorter than the equivalent transmission line. The two basic types are series and shunt as depicted below:

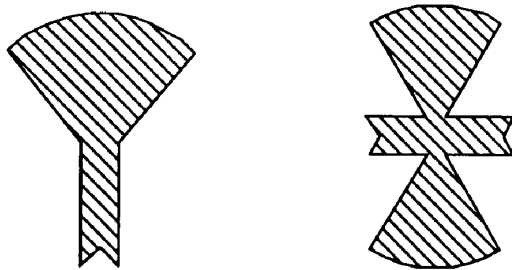


Figure 5.5.7.2: Radial Stubs: (a) Series (b) Shunt

Multiple stubs can be present on either the series or shunt configuration. In the shunt configuration, the added symmetry of multiple stubs may improve bandwidth.

The relations of March [8] include the effect of losses:

$$Z_{in} = -j \frac{120 \pi h}{r_i \theta \sqrt{\epsilon_{eff}}} \text{Ct}(\beta r_i, \beta r_o) \quad (\Omega) \quad (5.5.7.1)$$

or

$$C = \frac{1.0}{\omega \frac{120 \pi h}{r_i \theta \sqrt{\epsilon_{eff}}} \text{Ct}(\beta r_i, \beta r_o)} \quad (F) \quad (5.5.7.2)$$

where,

$$\epsilon_{eff} = \epsilon_{eff} \text{ of microstripline having width, } w_{eq,stub} = (r_i + r_o) \sin\left(\frac{\theta}{2}\right) \quad (5.5.7.3)$$

$$\text{Ct}(x, y) \equiv \frac{Y_0(x) J_1(y) - J_0(x) Y_1(y)}{J_1(x) Y_1(y) - Y_1(x) J_1(y)} = \text{large radial cotangent function [9]} \quad (5.5.7.4)$$

For small conductor losses we define

$$k = \beta - j \alpha \quad (5.5.7.5)$$

$$\beta = \frac{2 \pi \sqrt{\epsilon_{eff}}}{\lambda_0} \quad (5.5.7.6)$$

$$\alpha \equiv \frac{R(r)}{\lambda_0} \quad (5.5.7.7)$$

and the Bessel functions of complex argument can be written as

$$J_0(kr) = J_0(\beta r) + j \alpha r J_1(\beta r) \quad (5.5.7.8)$$

$$J_1(kr) = J_1(\beta r) - j \alpha r \left[J_0(\beta r) - \frac{J_1(\beta r)}{\beta} \right] \quad (5.5.7.9)$$

$$Y_0(kr) = Y_0(\beta r) + j \alpha r Y_1(\beta r) \quad (5.5.7.10)$$

$$Y_1(kr) = Y_1(\beta r) - j \alpha r \left[Y_0(\beta r) - \frac{Y_1(\beta r)}{\beta} \right] \quad (5.5.7.11)$$

$J_i(\beta r)$ = i th order Bessel function of the first kind having real argument

$Y_i(\beta r)$ = i th order Bessel function of the second kind having real argument

Atwater [3] reports good results modeling the radial stub as a series of short transmission lines similar to the delta stub model depicted in Figure 5.5.8.2. A complete model should include a microstrip line open end as described elsewhere in this chapter.

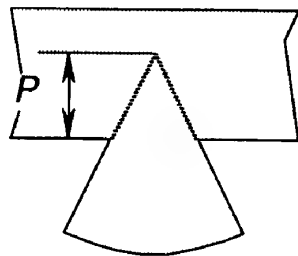


Figure 5.5.7.3: Definition of Radial Stub Insertion Depth, P

Giannini [5] notes that for shunt stubs, high frequencies, high dielectric constants, and r_i / w large the above relations may need to be corrected with effective dimensions. The effective radii r_i and r_o replace the actual radii in the calculations above.

$$r_{i,e} = r + \frac{2.0 \Delta}{\sin \frac{\theta}{2.0}} \quad (5.5.7.12)$$

$$r_{o,e} = r_o \sqrt{1.0 + \frac{2.0 h}{\pi r_o} \ln \left(\frac{\pi r_o}{2.0 h} \right) + 1.7726} \quad (5.5.7.13)$$

where

$$r = \frac{2.0 P + w_{eff} - w}{2.0 \cos \frac{\theta}{2.0}} \quad (5.5.7.14)$$

$$\Delta = \frac{w_{eff,2} - \alpha r}{2.0} \quad (5.5.7.15)$$

and

w_{eff} = the effective width of a microstrip line having width w

$w_{eff,2}$ = the effective width of a microstrip line having width αr

P = depth of insertion, defined in Figure 5.5.7.3 above

Angles greater than 180° have the disadvantage of requiring an additional section of line to access the radial stub, and the advantage of greater bandwidth. Atwater [2] has shown data that such stubs also have a wider bandwidth than the corresponding quarter-wave line. The loss can be calculated by analogy with the standard transmission line [8]:

$$\alpha_c = \frac{R_s \sqrt{\epsilon_{eff}}}{120 \pi h} \quad (5.5.7.16)$$

Giannini [4] reported measurement data comparing 180° stubs to three 60° stubs. The three 60° stub configuration yielded approximately 20% wider bandwidths.

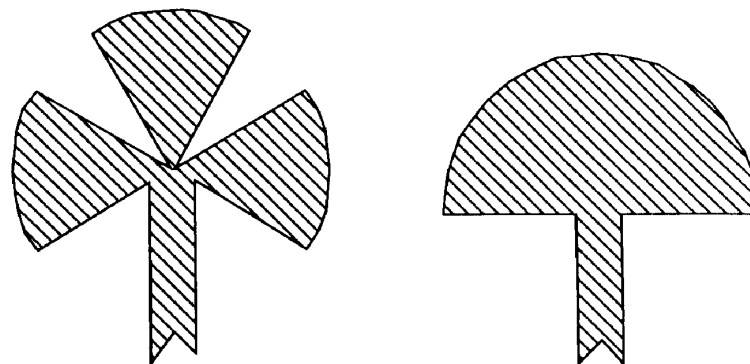


Figure 5.5.7.4: (a) Three 60° Radial Stubs (b) A Single 180° Radial Stub

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5.5.8 Delta Stub

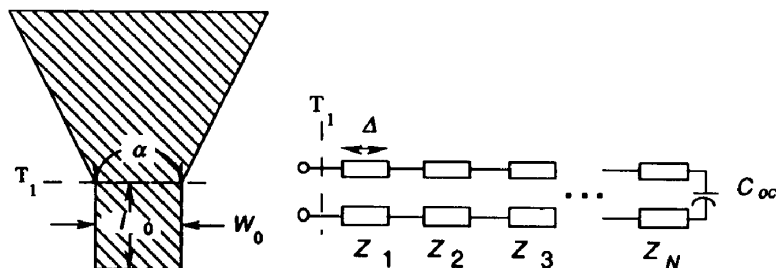


Figure 5.5.8.1: Delta Stub

The delta stub is used as a circuit element where a short is desired. Because of the straight sides it may be easier to lay out than the radial stub and it is smaller than an open-circuited transmission line.

The open end of the delta stub is transferred by the equivalent line length to a short circuit at the reference plane T. Coimbra [1, 2] recommends modeling the delta stub by a series of n transmission line segments of length Δ as shown below. The number, n , is chosen so each segment is $\ll \lambda_g$ in length. Each segment's characteristic impedance is

calculated with the standard microstrip line equations with the exception of the final section, which includes the open end discontinuity.

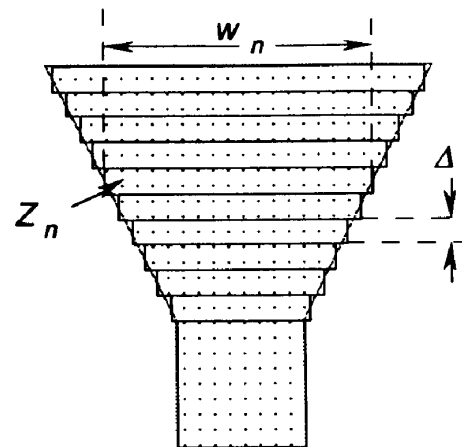


Figure 5.5.8.2: Modeling the Delta Stub

Coimbra used an iterative technique for synthesis. The impedance of the cascaded lines is calculated as each section is added until the desired open-circuit impedance is achieved.

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5.5.9 Microstrip Line Cross

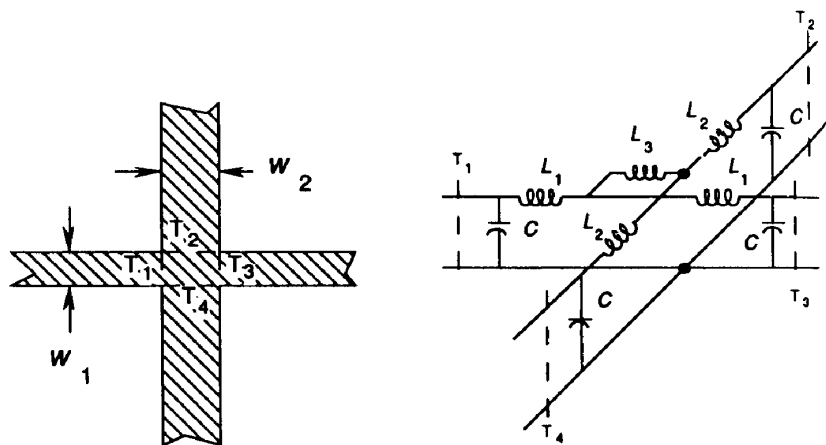


Figure 5.5.9.1: Microstrip Line Cross

The cross junction is used when we wish to parallel two stubs to get a lower impedance. The design equations [3] are:

$$\frac{C}{w_1} = \left\{ \left[\frac{86.6 w_2}{h} - 30.9 \left(\frac{w_2}{h} \right)^{1/2} + 367.0 \right] \log \left(\frac{w_1}{h} \right) + \left(\frac{w_2}{h} \right)^3 + \frac{74.0 w_2}{h} + 130.0 \right\} \left(\frac{w_1}{h} \right)^{-1/3} - 240.0 + \frac{2.0}{w_2/h} - \frac{1.5 w_1}{h} \left(1.0 - \frac{w_2}{h} \right) \quad (\text{pF/m}) \quad (5.5.9.1)$$

$$L_1/h = \left\{ \left[\frac{165.6 w_2}{h} + 31.2 \sqrt{\frac{w_2}{h}} - 11.8 \left(\frac{w_2}{h} \right)^2 \right] \frac{w_1}{h} - \frac{32.0 w_2}{h} + 3.0 \right\} \left(\frac{w_1}{h} \right)^{-3/2} \quad (\text{nH/m}) \quad (5.5.9.2)$$

$$L_2/h = \left\{ \left[\frac{165.6 w_1}{h} + 31.2 \sqrt{\frac{w_1}{h}} - 11.8 \left(\frac{w_1}{h} \right)^2 \right] \frac{w_2}{h} - \frac{32.0 w_1}{h} + 3.0 \right\} \left(\frac{w_2}{h} \right)^{-3/2} \quad (\text{nH/m}) \quad (5.5.9.3)$$

$$L_3/h = 337.5 + \left(1.0 + \frac{7}{w_1/h} \right) \frac{1.0}{w_2/h} - \frac{5.0 w_2}{h} \cos \left[\frac{\pi (1.5 h - w_1)}{2 h} \right] \quad (\text{nH/m}) \quad (5.5.9.4)$$

The equation for C is valid for

$$\epsilon_r = 9.9$$

$$0.3 \leq w_1/h \leq 3.0$$

$$0.1 \leq w_2/h \leq 3.0$$

Equations for L_1 , L_2 , and L_3 are valid for

$$0.5 \leq w_1/h \leq 2.0$$

$$0.5 \leq w_2/h \leq 2.0$$

The equations are accurate to within 5%.

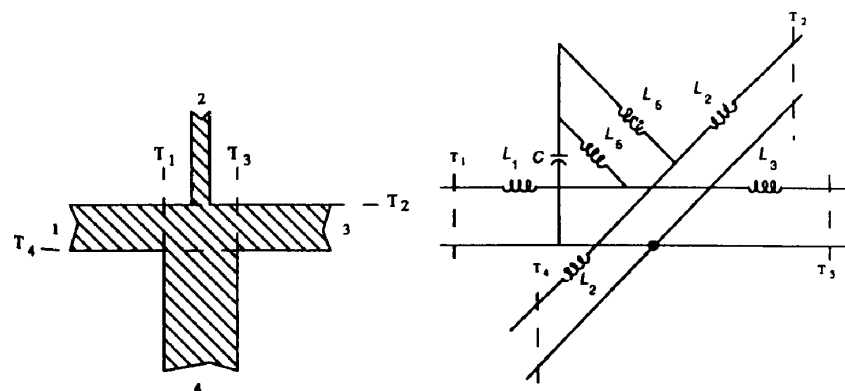


Figure 5.5.9.2: Asymmetric Cross Junction

Akello [1] reports on measurements of the asymmetric cross junction. For $Z_2 \gg Z_1$, Z_3 , Z_4 the affect of line 2 is shown to be small. This allows the discontinuity at the junction to be approximated with the equations for a microstrip line T.

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5.5.10 Microstrip Line T Junction

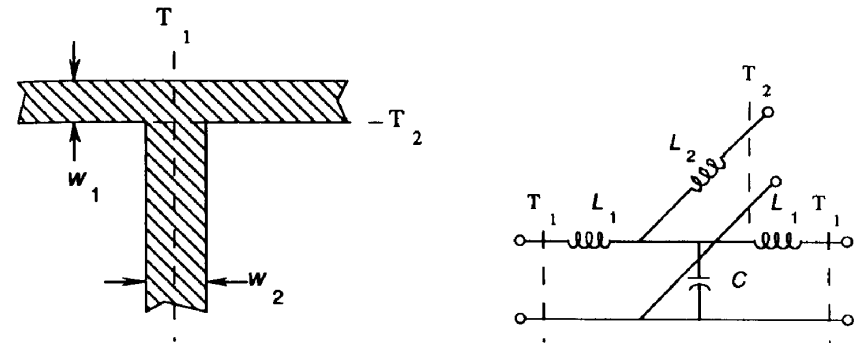


Figure 5.5.10.1: Microstrip Line T Junction

The T junction is used when we wish to split a signal into two paths. This is used in various combiner/dividers. A number of references give design curves for the discontinuity. Design equations are available in [3]

$$\frac{C_T}{w_1} = \frac{100.0}{\tanh(0.0072 Z_0)} + 0.64 Z_0 - 261.0 \quad (\text{pF/m})$$

For $0.5 \leq w_1/h \leq 2.0$ and $0.5 \leq w_2/h \leq 2.0$,

$$\frac{\Delta L_1}{h} = \frac{-w_2}{h} \left[\frac{w_2}{h} \left(-0.016 \frac{w_1}{h} + 0.064 \right) + \frac{0.016}{\frac{w_1}{h}} \right] L_{w,1}$$

For $1.0 \leq w_1/h \leq 2.0$ and $0.5 \leq w_2/h \leq 2.0$,

$$\frac{\Delta L_2}{h} = \left[\left(\frac{0.12 w_1}{h} - 0.47 \right) \frac{w_2}{h} + 0.195 \frac{w_1}{h} - 0.357 \right. \\ \left. + 0.0283 \sin \left(\frac{\pi w_1}{h} - 0.75 \pi \right) \right] L_{w,2}$$

For each line width we calculate

$$L_{w,n} = \frac{Z_0(w_n) \sqrt{\epsilon_{eff}}}{c}$$

The equations are valid to an accuracy of better than 5% for $\epsilon_r = 9.9$ and

$$25.0 \leq Z_0 \leq 100.0$$

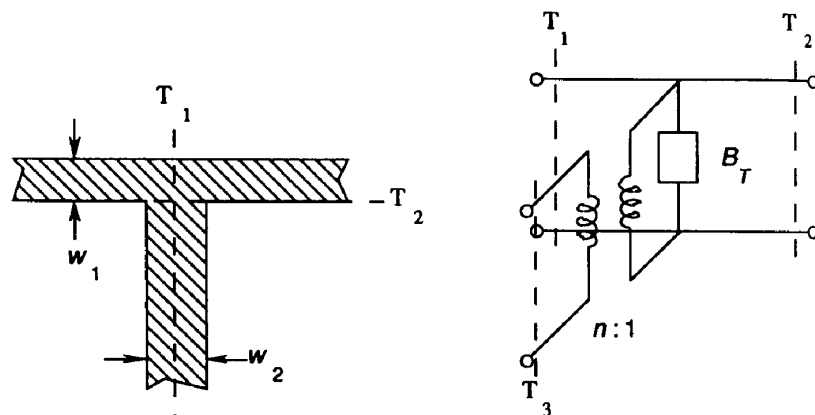


Figure 5.5.10.2: Microstrip Line T Junction

An alternative model, shown in Figure 5.5.10.2, is given in Hammerstad [4]. The design equations are

$$d_1 = \frac{0.05 D_2 Z_1 n^2}{Z_2}$$

$$d_2 = \frac{D_1}{2.0} - 0.16 \frac{D_1 Z_1}{Z_2} \left[1.0 + \left(\frac{2.0 D_1}{\lambda} \right)^2 - 2.0 \ln \frac{Z_1}{Z_2} \right]$$

$$n = \frac{\sin \left(\frac{2.0 \pi D_1 Z_1}{2.0 \lambda Z_2} \right)}{\frac{2.0 \pi D_1 Z_1}{2.0 \lambda Z_2}}$$

for $\frac{Z_1}{Z_2} \leq 0.5$

$$B_T = \frac{-D_1 (1.0 - 2.0 D_1 / \lambda)}{\lambda Z_2}$$

and for $\frac{Z_1}{Z_2} \geq 0.5$

$$B_T = \frac{-D_1 \left(\frac{3.0 Z_1}{Z_2} - 2.0 \right) (1.0 - 2.0 D_1 / \lambda)}{\lambda Z_2}$$

where

$$D_i = \frac{\eta_0 h}{\sqrt{\epsilon_{eff,i}} Z_i}$$

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5.5.11 Asymmetrical Microstrip Line T Junction

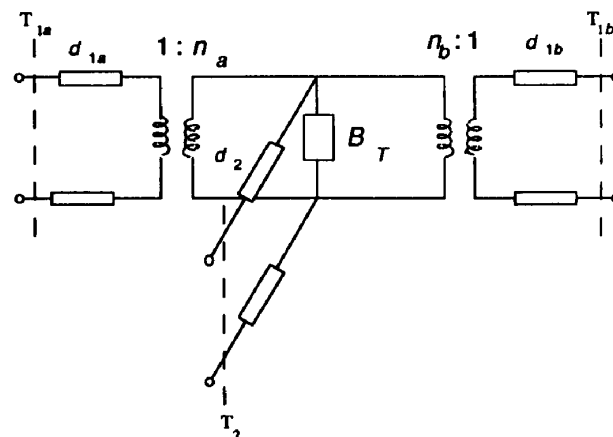
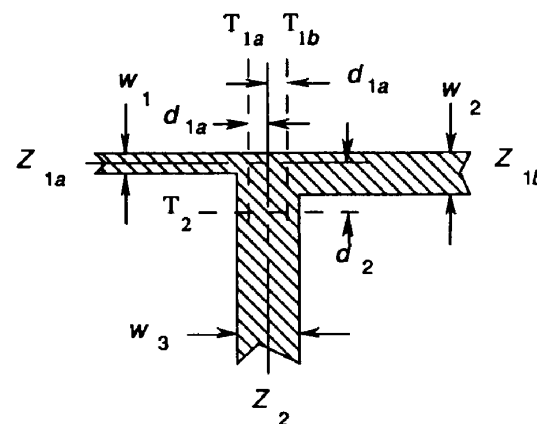


Figure 5.5.11.1: Asymmetric Microstrip Line T Junction and Model

Hammerstad [1] recommends that we modify the equations for a symmetric T junction to model the asymmetric junction. The equations are

$$\frac{d_{1a}}{D_2} = 0.055 \left[1.0 - 1.0 \frac{Z_{1a}}{Z_2} \left(\frac{f}{f_{p1}} \right)^2 \right] \frac{Z_{1a}}{Z_2} \quad (5.5.11.1)$$

$$\frac{d_{1b}}{D_2} = 0.055 \left[1.0 - 1.0 \frac{Z_{1b}}{Z_2} \left(\frac{f}{f_{p1}} \right)^2 \right] \frac{Z_{1b}}{Z_2} \quad (5.5.11.2)$$

$$\frac{d_2}{D_1} = 0.5 - \left[0.05 + 0.7 e^{-1.6 \sqrt{Z_{1a} Z_{1b}} / Z_2} + 0.25 \frac{\sqrt{Z_{1a} Z_{1b}}}{Z_2} \left(\frac{f}{f_{p1}} \right)^2 - 0.17 \ln \frac{\sqrt{Z_{1a} Z_{1b}}}{Z_2} \right] \frac{\sqrt{Z_{1a} Z_{1b}}}{Z_2} \quad (5.5.11.3)$$

$$n_a = \sqrt{1.0 - \pi \left(\frac{f}{f_{p1}} \right)^2 \left[\frac{1.0}{12.0} \frac{Z_{1a}^2}{Z_2^2} + \left(0.5 - \frac{d_2}{D_1} \right)^2 \right]} \quad (5.5.11.4)$$

$$n_b = \sqrt{1.0 - \pi \left(\frac{f}{f_{p1}} \right)^2 \left[\frac{1.0}{12.0} \frac{Z_{1b}^2}{Z_2^2} + \left(0.5 - \frac{d_2}{D_1} \right)^2 \right]} \quad (5.5.11.5)$$

$$\frac{B_T \lambda_1}{Y_2 D_1} = 5.5 \frac{\epsilon_r + 2.0}{\epsilon_r} \left[1.0 + 0.9 \ln \frac{\sqrt{Z_{1a} Z_{1b}}}{Z_2} + 4.5 \frac{\sqrt{Z_{1a} Z_{1b}}}{Z_2} \left(\frac{f}{f_{p1}} \right)^2 - 4.4 \exp \left(\frac{-1.3 \sqrt{Z_{1a} Z_{1b}}}{Z_2} \right) - 20.0 \left(\frac{Z_2}{\eta_0} \right)^2 \right] n^{-2.0} \frac{d_1}{D_2} \quad (5.5.11.6)$$

where

$$f_p = 0.4 Z / h$$

The variable f_p is in GHz, Z in Ω , and h in mm. The Z 's are calculated with the usual microstrip line equations.

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5.5.12 Optimal Microstrip Line T

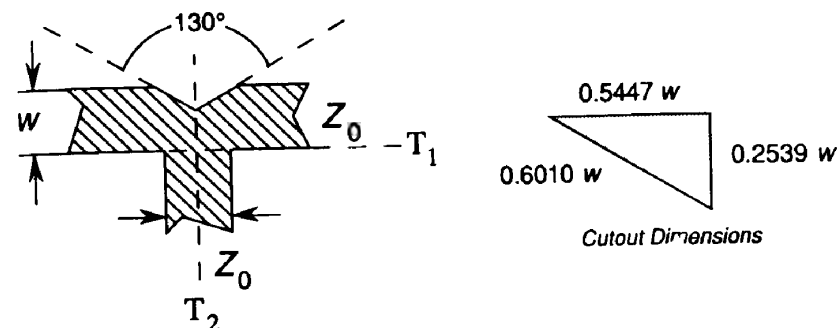


Figure 5.5.12.1: Optimal Microstrip Line T and Notch Dimensions

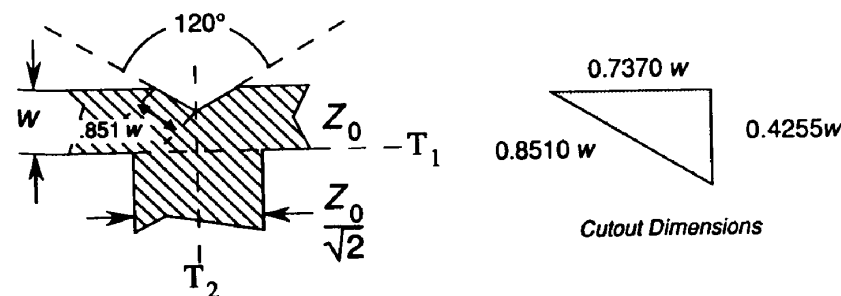


Figure 5.5.12.2: Optimal Microstrip Line T and Notch Dimensions

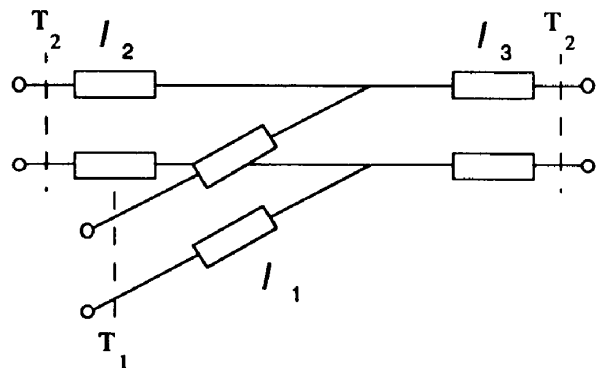


Figure 5.5.12.3: Compensated Microstrip Line T Model

The microstrip line T junction can be compensated by removing a triangular section of the line. For equal impedance lines the optimal compensating notch has an angle of $>120^\circ$ (approximately 130° from Figure 8 [1]) and the optimal length of a side, a , is $0.601w$ for 130° . For a main line impedance of $Z_0/\sqrt{2}$ the optimal angle is 120° and the optimal side is $0.851w$. The dimension a is independent of frequency, dielectric, and other parameters.

The compensated equal impedance T junction reference plane extensions are approximately independent of frequency [1]. For the case shown in Figure 5.5.12.1, the low frequency value is

$$\Delta l_1 / h \approx 0.45$$

$$\Delta l_2 / h = \Delta l_3 / h \approx 0.7$$

For the case shown in Figure 5.5.12.2 the low frequency value is

$$\Delta l_1 / h \approx 0.2$$

$$\Delta l_2 / h = \Delta l_3 / h \approx 0.2$$

Exact values and curves showing the frequency dependency are available in [1], Figures 7 and 9.

In Wu *et al.* [2], another technique of compensating a T is proposed and analyzed by the method of moments. Good results were achieved; however, no design equations are available.

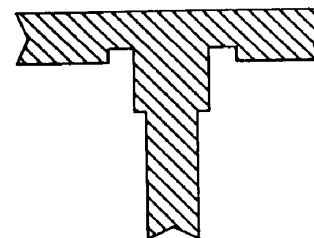


Figure 5.5.12.4: Shaped T junction

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5.5.13 Microstrip Line Step (Abrupt Change in Width)

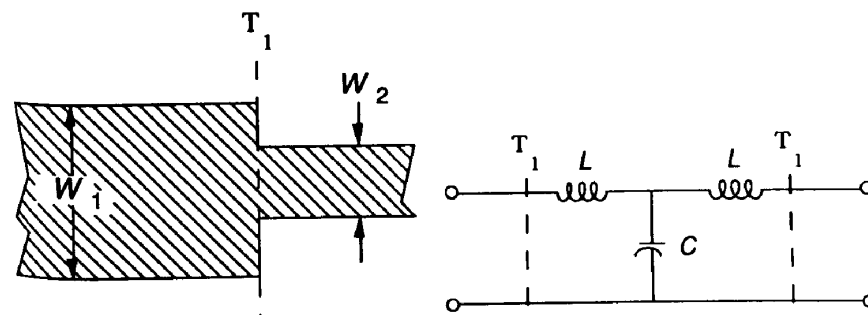


Figure 5.5.13.1: Microstrip Line Abrupt Width Change

The impedance step is modeled by the equivalent circuit of Figure 5.5.13.1 above. A number of references [1-5, 8, 13] give design curves. The design equations of [3] are given below.

$$\text{For } 1.5 \leq \frac{W_1}{W_2} \leq 3.5,$$

$$C_s = \sqrt{w_1 w_2} \left[(10.1 \log \epsilon_r + 2.33) \frac{w_1}{w_2} - 12.6 \log \epsilon_r - 3.17 \right] \quad (\text{pF/m}) \quad (5.5.13.1)$$

The stated accuracy is better than 10% for

$$w_1 > w_2$$

$$\epsilon_r \leq 10.0$$

For $3.5 \leq \frac{w_1}{w_2} \leq 10.0$ and $\epsilon_r = 9.6$,

$$C_s = \sqrt{w_1 w_2} \left[130.0 \log \left(\frac{w_1}{w_2} \right) - 44.0 \right] \quad (\text{pF/m}) \quad (5.5.13.2)$$

which is accurate to within 0.5%. The inductance is calculated with

$$L_s = h \left[40.5 \left(\frac{w_1}{w_2} - 1.0 \right) - 75.0 \log \left(\frac{w_1}{w_2} \right) + 0.2 \left(\frac{w_1}{w_2} - 1.0 \right) \right] \quad (\text{nH/m}) \quad (5.5.13.3)$$

This equation is accurate to within 5.0% for

$$\frac{w_1}{w_2} \leq 5.0$$

$$\frac{w_2}{h} = 1.0$$

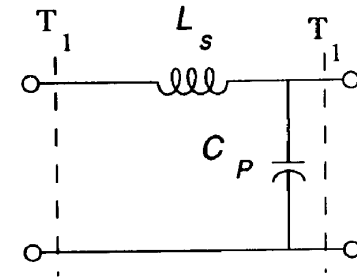
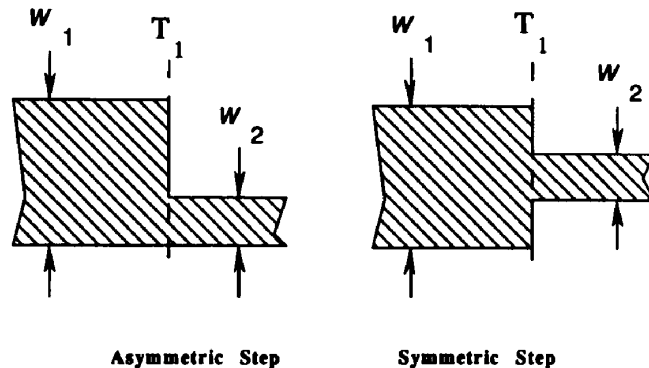


Figure 5.5.13.2: Microstrip Line Steps and Equivalent Circuit

A frequency-dependent model of the asymmetric step is presented in Hoffman [7]. The equivalent circuit is shown in Figure 5.5.13.2 and the model parameters are

$$L_s = \frac{Z_{L1}}{2.0 \pi f} \frac{4.0 w_{eff,1}}{\Delta \lambda_1} \left\{ \ln \left[\left(\frac{1.0 - a^2}{4.0 a} \right) \left(\frac{1.0 + a}{1.0 - a} \right)^{(a+1/a)/2.0} \right] + 2.0 \frac{A + B + 2.0 D}{A B - D^2} + \left(\frac{w_{eff,1}}{2.0 \Delta \lambda_1} \right)^2 \left(\frac{1.0 - a}{1.0 + a} \right)^{4.0 a} \right. \\ \left. \left(\frac{5.0 a^2 - 1.0}{1.0 - a^2} + \frac{4.0 a^2 D}{3.0 A} \right)^2 \right\} \quad (5.5.13.4)$$

$$C_p \equiv (0.012 + 0.0039 \epsilon_r) (w_1 - w_2) \quad (\text{pF}) \quad (5.5.13.5)$$

where

$$A = \left(\frac{1.0 + a}{1.0 - a} \right)^{2.0 a} \left[\frac{1.0 - \sqrt{1.0 - \left(\frac{2.0 w_{eff,1}}{\Delta \lambda_1} \right)^2}}{1.0 - \sqrt{1.0 - \left(\frac{2.0 w_{eff,1}}{\Delta \lambda_1} \right)^2}} \right] - \frac{1.0 + 3.0 a^2}{1.0 - a^2} \quad (5.5.13.6)$$

$$B = \left(\frac{1.0 + a}{1.0 - a} \right)^{a/2.0} \left[\frac{1.0 + \sqrt{1.0 - \left(\frac{2.0 w_{eff,2}}{\Delta \lambda_2} \right)^2}}{1.0 - \sqrt{1.0 - \left(\frac{2.0 w_{eff,2}}{\Delta \lambda_2} \right)^2}} \right] - \frac{3.0 + a^2}{1.0 - a^2}$$

$$D = \left(\frac{4.0 a}{1.0 - a^2} \right)^2$$

$$a = \frac{w_{eff,2}}{w_{eff,1}}$$

$$\lambda_n = \frac{c_0}{f \sqrt{\epsilon_{eff,n}}}$$

$$\Delta = \begin{cases} 1 & \text{for symmetrical steps} \\ 2 & \text{for asymmetrical steps} \end{cases}$$

and

$$w_{eff,n} = w_n + \frac{w_{eff,static} - w_n}{1.0 + f / f_{p5}}$$

$$w_{eff,static} = \frac{\eta_0 h}{Z_{L0}}$$

Z_{L0} = the impedance of the line calculated with ϵ_r set to 1.0

$$f_{p5} = \frac{c}{2.0 w \sqrt{\epsilon_r}}$$

Equations are accurate to within a few percent for

$$w_1 > w_2$$

$$0 < Z_{L1} < Z_{L2}$$

$$0 < f < f_{g,HE2}$$

where

$$f_{g,HE2} = \frac{c}{w_{eff,1} \sqrt{\epsilon_{eff}}}$$

(5.5.13.7)

(5.5.13.8)

(5.5.13.9)

(5.5.13.10)

(5.5.13.11)

(5.5.13.12)

(5.5.13.13)

(5.5.13.14)

(5.5.13.15)

(5.5.13.16)

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5.5.14 Compensated Microstrip Line Step

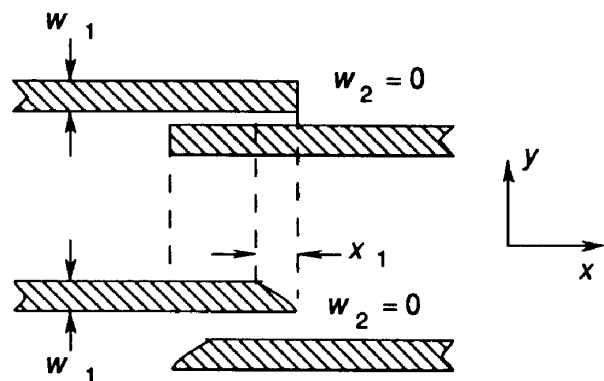


Figure 5.5.14.1: Compensated Open End Discontinuities

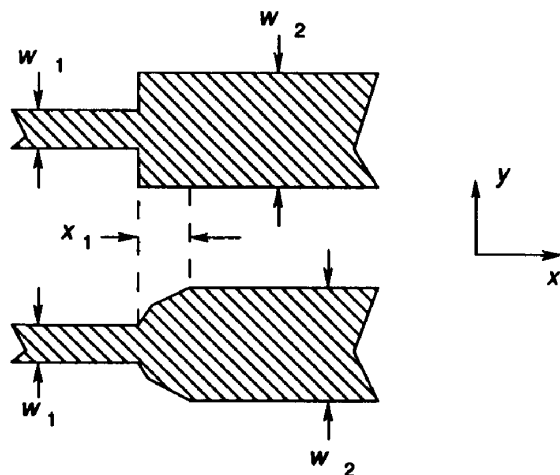


Figure 5.5.14.2: Compensated Step Discontinuities

Step discontinuities occur in impedance transformers, filters, and resonators, when interfacing to a device pin and in many other situations. Although the impedance step is desired, the discontinuity creates unwanted parasitics. The technique of this section removes this parasitic.

The presence of a step discontinuity in any TEM line can be compensated by the addition of a taper near the discontinuity as shown in the figures above (upper figures are uncompensated, lower are compensated). The cross-section is varied continuously to move the transmission line's inherent capacitance gradually into fringing capacitance, thus keeping the impedance constant until the step is complete. This does not remove the impedance change, just the parasitics associated with the step.

For microstrip line, this procedure of removing line capacitance and replacing it with fringing capacitance is used. Hoefler [1] gives

$$\frac{dy}{dx} = \pm \sqrt{\left[\frac{\frac{\eta_0 h}{2.0 y \epsilon_r} \frac{\epsilon_{eff}(w_2)}{Z_0(\epsilon_r = 1.0, w_2)} - 1.0}{\frac{\eta_0 h}{2.0 y \epsilon_r} \frac{\epsilon_{eff}(2.0 y)}{Z_0(\epsilon_r = 1.0, 2.0 y)} - 1.0} \right]^2 - 1.0} \quad (5.5.14.1)$$

where

$\epsilon_{eff}(w)$ = the effective relative dielectric constant of a microstrip line having width w (see Chapter 3)

$Z_0(\epsilon_r = 1.0, w)$ = the impedance of microstrip line having width w and the dielectric replaced with air (see Chapter 3)

y = distance from centerline of the microstrip line

To use this equation, Hoefler recommends starting with $y = w_1 / 2.0$ and using approximately 50 steps. The change in width dy / dx is calculated and added to the previous step width until the desired final width w_2 is reached.

Fewer steps can be used initially and then increased until the contour dimensions converge. This process and the equation above can be used for impedance steps and open end corrections. Asymmetric compensations are also described in Hoefler.

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5.5.15 Optimal Microstrip Line Linear Taper

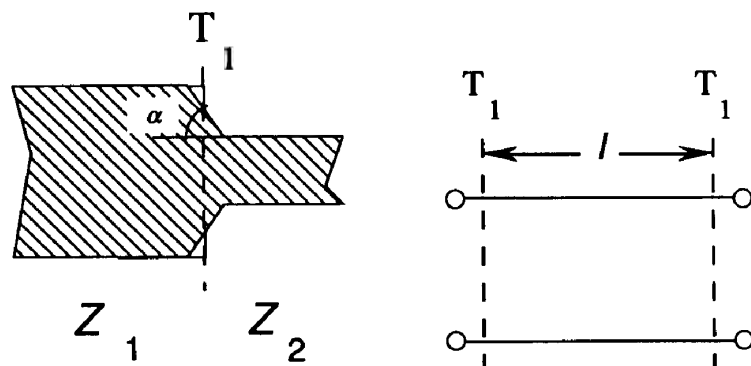


Figure 5.5.15.1: Optimal Linear Tapers

A simple approximate compensation for the strays of the impedance step is the linear taper. The optimal linear taper for 1:2 and $1:\sqrt{2}$ steps is found by Chadha and Gupta [1] to be 60° centered at the original location of the step. Chadha and Gupta also gives plots of reference plane shift amount (shifts towards low Z line).

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5.5.16 Microstrip Line Double Step

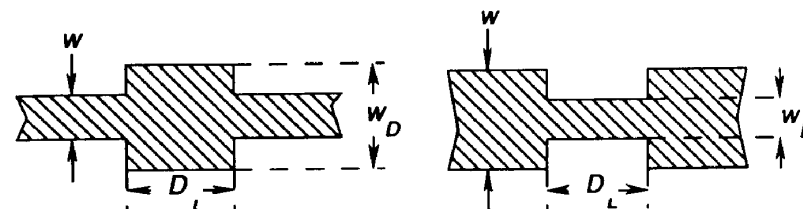


Figure 5.5.16.1: Microstrip Line Double Step

Impedance steps can occur repetitively in filters. Usually these are modeled as a series of transmission lines connected with step discontinuities. In cases where they are spaced closely, the steps will interact. In Hill [3], it is shown with a full wave analysis that two steps separated by $D_L = w_D$ interact significantly and cannot be modeled by cascading the lines and step discontinuities. Hill also found that for $D_L = 3.0 \times w_D$ and $D_L = 8.0 \times w_D$ the isolated models were valid. In Anders and Arndt [1], curves are shown for D_L / w_D less than 2.0. For the cases shown, uncoupled models appear to work well for $D_L > 2.0 \times w_D$.

Giannini *et al.* [2] describe an advanced model that can be used with existing CAE tools to model interacting steps.

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5.5.17 Microstrip Line Open

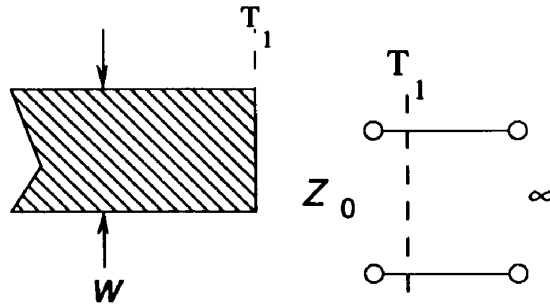


Figure 5.5.17.1: Microstrip Line Open

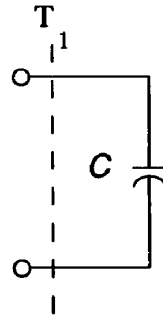


Figure 5.5.17.2: Alternative Microstrip Line Open Model

The open end adds a capacitance that effectively lengthens the transmission line. The additional length can be calculated (h is as shown in Figure 3.5.1.1) [10]:

$$\frac{\Delta l}{h} = \left(\frac{\xi_1 \xi_3 \xi_5}{\xi_4} \right) \quad (5.5.17.1)$$

$$\xi_1 = 0.434907 \left[\frac{\epsilon_{eff}^{0.81} - 0.26}{\epsilon_{eff}^{0.81} - 0.189} \right] \left[\frac{(w/h)^{0.8544} + 0.236}{(w/h)^{0.8544} + 0.87} \right] \quad (5.5.17.2)$$

$$\xi_2 = 1 + \frac{(w/h)^{0.371}}{2.358 \epsilon_r + 1} \quad (5.5.17.3)$$

$$\xi_3 = 1.0 + \frac{0.5247 \tan^{-1} [0.084 (w/h)^{1.9413} / \xi_2]}{\epsilon_{eff}^{0.9236}} \quad (5.5.17.4)$$

$$\xi_4 = 1.0 + 0.0377 \tan^{-1} [0.067 (w/h)^{1.456}] [6.0 - 5.0 e^{0.036(1.0 - \epsilon_r)}] \quad (5.5.17.5)$$

$$\xi_5 = 1.0 - 0.218 e^{-7.5 w/h} \quad (5.5.17.6)$$

When compared with numerical results at 1.0 GHz, accuracy is better than 2.5% for

$$0.01 \leq \frac{w}{h} \leq 100.0$$

$$\epsilon_r \leq 50.0$$

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5.5.18 Coupled Microstrip Line Open End

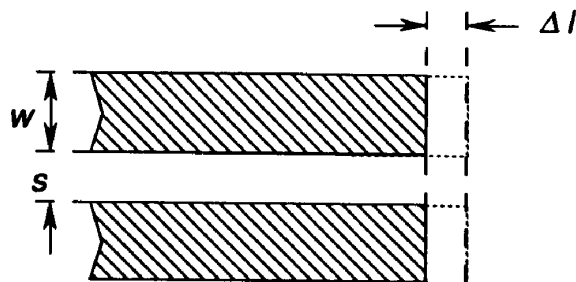


Figure 5.5.18.1: Coupled Microstrip Line Open End

Coupled microstrip open ends occur in couplers and coupled line filters. The effects of open ends in these circuits are modeled by an electrical lengthening of the physical strips after Kirschning and Jansen [3].

$$\Delta l_e = \left[\Delta l(2u, \epsilon_r) - \Delta l(u, \epsilon_r) + 0.0198 h g^{R_1} \right] e^{-0.328 g^{2.244}} + \Delta l(u, \epsilon_r)$$

$$\Delta l_0 = \left[\Delta l(u, \epsilon_r) - h R_3 \right] \left(1.0 - e^{-R_4} \right) + h R_3 \quad (5.5.18.1)$$

$$\Delta l_0 = \left[\Delta l(u, \epsilon_r) - h R_3 \right] \left(1.0 - e^{-R_4} \right) + h R_3 \quad (5.5.18.2)$$

where

$\Delta l(u, \epsilon_r)$ = length correction for a single line of width w

$\Delta l(2u, \epsilon_r)$ = length correction for a single line of width $2w$

$$R_1 = 1.187 \left(1.0 - e^{-0.069 u^{2.1}} \right) \quad (5.5.18.3)$$

$$R_2 = 0.343 u^{0.6187} + \frac{0.45 \epsilon_r}{1.0 + \epsilon_r} u^{1.357 + \frac{1.65}{1.0 + 0.7 \epsilon_r}} \quad (5.5.18.4)$$

$$R_3 = 0.2974 \left(1.0 - e^{-R_2} \right) \quad (5.5.18.5)$$

$$R_4 = (0.271 + 0.0281 \epsilon_r) g^{\left(\frac{1.167 \epsilon_r}{0.66 + \epsilon_r} \right)} + \left(\frac{1.025 \epsilon_r}{0.687 + \epsilon_r} \right) g^{\left(\frac{0.958 \epsilon_r}{0.706 + \epsilon_r} \right)} \quad (5.5.18.6)$$

The values of $\Delta l(u, \epsilon_r)$ and $\Delta l(2u, \epsilon_r)$ are calculated with the relationships in this book for the microstrip line open discontinuity. The stated accuracy to the numerical analysis for the length corrections is 5%.

Lau [4] compared simulations with network analysis data for coupled microstrip lines. The open end effect was found significant for frequencies over about 24 GHz. The effect of the open ends coupling to the opposite line was also modeled by the addition of 2 fF at each of the lines.

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5.5.19 Microstrip Line Gap

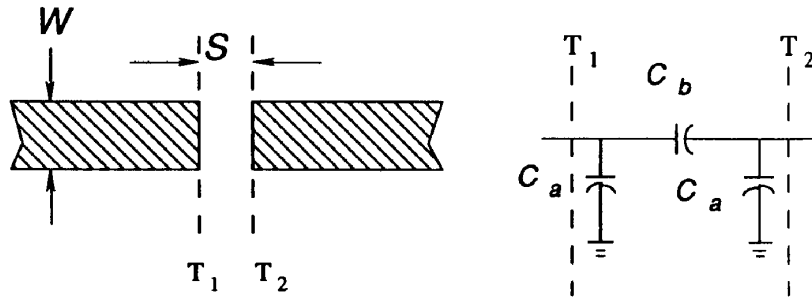


Figure 5.5.19.1: Microstrip Line Gap

The microstrip gap is used in filters and other circuits as a coupling element. It also occurs as a stray circuit element across surface mount components.

Hammerstad [5] fits the curves presented in [1] with the following simple relations:

$$B_g = 2.4 \frac{Y h}{\lambda} \frac{u + 0.1}{u + 1.0} \sqrt{\frac{\epsilon_r + 2.0}{\epsilon_r + 1.0}} \ln \left[\frac{\epsilon_r s}{h (\epsilon_r + 2.0)} \right] \quad (5.5.19.1)$$

The capacitance of C_a can be modeled by an equivalent length of transmission line, Δ_p :

$$\Delta_p = \Delta_e \tanh^2 \sqrt{0.5 s / \Delta_e} \quad (5.5.19.2)$$

$$\Delta_g = \Delta_p + \frac{B_g \lambda_g}{2.0 \pi Y} \quad (5.5.19.3)$$

where

$$Y = \text{line admittance (S)} \quad (5.5.19.4)$$

$$u = w / h \quad (5.5.19.5)$$

$$\Delta_e = 0.412 h \frac{\epsilon_{eff} + 0.3}{\epsilon_{eff} - 0.258} \frac{u + 0.262}{u + 0.813} = \text{length extension of an open end} \quad (5.5.19.6)$$

which are valid for

$$\epsilon_r = 9.8 \text{ and } 25.0 \leq Z_0 \leq 70.0$$

$$\epsilon_r = 2.2 \text{ and } 35.0 \leq Z_0 \leq 100.0$$

The accuracy is 0.05 for B_g / Y and 0.1 h for Δ_g .

For references such as [1] that give equations or charts of C_{even} and C_{odd} the relations to the model of Figure 5.5.14.1 are:

$$C_a = C_{even} / 2.0 \quad (5.5.19.7)$$

$$C_b = C_{odd} / 2.0 - C_{even} / 4.0 \quad (5.5.19.8)$$

The effect of C_a can also be modeled by replacing C_a with the equivalent line extension. The relation is

$$\Delta l = C_a Z_0 v_p \quad (5.5.19.9)$$

Özmehtmet [9] improves the above model by adding frequency dependence and the effects of radiation losses (a resistor R_r in parallel with C_b). The series capacitance C_b is split into three parts representing the parallel-plate capacitance of the facing trace ends and two fringing capacitances. The frequency-dependent ϵ_{eff} is used to calculate their values after [1]. The radiation resistance is given by

$$R_r = \frac{120 \pi^2}{K} \quad (5.5.19.10)$$

where

$$K = \frac{\sin 2a}{2.0a} + \cos 2.0a + 2.0a \text{ Si}(2.0a) - 1.0 \quad (5.5.19.11)$$

$$a = \pi W_e(f) / \lambda_0 \quad (5.5.19.12)$$

$W_e(f)$ = the frequency-dependent effective width. Calculate by equating the incremental capacitance of the frequency-dependent microstrip line to a parallel-plate capacitor having the same dielectric and plate spacing equal to the dielectric thickness.

(5.5.19.13)

$\text{Si}(x)$ = sine integral function

Garg and Bahl [4] recommend the following model

$$C_{\text{even}} / w = (\epsilon_r / 9.6)^{0.9} (s / w)^{m_e} e^{K_e} \quad (\text{pF} / \text{m}) \quad (5.5.19.14)$$

$$C_{\text{odd}} / w = (\epsilon_r / 9.6)^{0.8} (s / w)^{m_o} e^{K_o} \quad (\text{pF} / \text{m}) \quad (5.5.19.15)$$

where for $0.1 \leq s / w \leq 1.0$

$$m_o = \frac{w}{h} \left(0.619 \log \frac{w}{h} - 0.3853 \right) \quad (5.5.19.16)$$

$$K_o = 4.26 - 1.453 \log \frac{w}{h} \quad (5.5.19.17)$$

and for $0.1 \leq s / w \leq 0.3$

$$m_e = 0.8675 \quad (5.5.19.18)$$

$$K_e = 2.043 (w / h)^{0.12} \quad (5.5.19.19)$$

For $0.3 \leq s / w \leq 1.0$

$$m_e = \frac{1.565}{(w / h)^{0.16}} - 1.0 \quad (5.5.19.20)$$

$$K_e = 1.97 - \frac{0.03}{w / h} \quad (5.5.19.21)$$

which is valid for

$$2.5 \leq \epsilon_r \leq 15.0$$

The accuracy is within 7%.

REFERENCES

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- [8] Maeda, Minora, "An Analysis of Gap in Microstrip Transmission Lines," *IEEE Transactions on Microwave Theory and Techniques*, Vol. MTT-20, No. 6, June 1972, pp. 390-396.
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- [10] Yang, Hung-Yu, et al., "Microstrip Open-End and Gap Discontinuities in a Substrate-Superstrate Structure," *IEEE Transactions on Microwave Theory and Techniques*, Vol. MTT-37, No. 10, October 1989, pp. 1542-1546.

5.5.20 Asymmetric Microstrip Line Gap

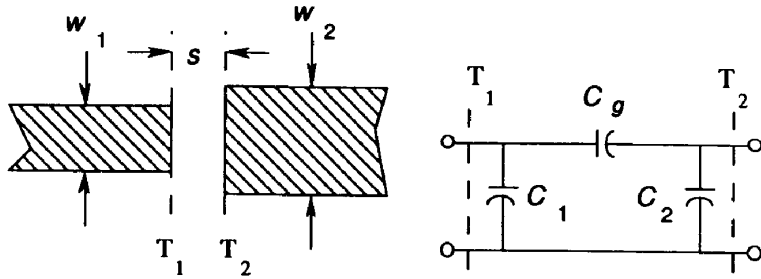


Figure 5.5.20.1: Asymmetric Microstrip Line Gap

Kirschning *et al.* [2] gives these equations for the asymmetric microstrip line gap:

$$B_g = Q_1 \pi f h e^{(-1.86 g/h)} \left\{ 1.0 + 4.19 \left[1.0 - e^{-0.785 (h/w_1)^{0.5} w_2/w_1} \right] \right\} \quad (\text{mS}) \quad (5.5.20.1)$$

$$B_1 = \frac{2.0 \pi f C_1 (Q_2 + Q_3)}{1.0 + Q_2} \quad (\text{mS}) \quad (5.5.20.2)$$

$$B_2 = \frac{2.0 \pi f C_2 (Q_2 + Q_4)}{1.0 + Q_2} \quad (\text{mS}) \quad (5.5.20.3)$$

where C_1 and C_2 are the open end capacitances of microstrip lines having widths w_1 and w_2 in pF. The frequency f is in GHz and h is in mm. The Q_n 's are given by

$$Q_1 = 0.04598 \left\{ 0.03 + (w_1/h)^{1.23} / \left[1.0 + 0.12 (w_2/w_1 - 1.0)^{0.9} \right] \right\} \times (0.272 + 0.07 \epsilon_r) \quad (5.5.20.4)$$

$$Q_2 = 0.10700 (w_1/h + 9.0) (g/h)^{3.23} + 2.09 (g/h)^{1.05} \left(\frac{1.5 + 0.3 w_1/h}{1.0 + 0.6 w_1/h} \right) \quad (5.5.20.5)$$

$$Q_3 = e^{[-0.5978 (w_2/w_1)^{1.35}] - 0.55} \quad (5.5.20.6)$$

$$Q_4 = e^{[-0.5978 (w_1/w_2)^{1.35}] - 0.55} \quad (5.5.20.7)$$

and the equations are valid for

$$0.1 \leq w_1/h \leq 3.0$$

$$0.1 \leq w_2/h \leq 3.0$$

$$1.0 \leq w_2/w_1 \leq 3.0$$

$$0.2 \leq g/h \leq \infty$$

$$6.0 \leq \epsilon_r \leq 13.0$$

Error is less than 0.1 mS for frequencies up to 18 GHz and $f \cdot h$ products up to 12 GHz · mm.

REFERENCES

- [1] Biswas, Animesh, and Vijai K. Tripathi, "Modeling of Asymmetric and Offset Gaps in Shielded Microstrips and Slotlines," *IEEE Transactions on Microwave Theory and Techniques*, Vol. MTT-38, No. 6, June 1990, pp. 818–822.
- [2] Kirschning, M., *et al.*, "Measurement and Computer-Aided Modeling of Microstrip Discontinuities by an Improved Resonator Method," *IEEE Microwave Theory and Techniques Symposium Digest 1983*, pp. 495–497.
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5.5.21 Microstrip Line Finger Break

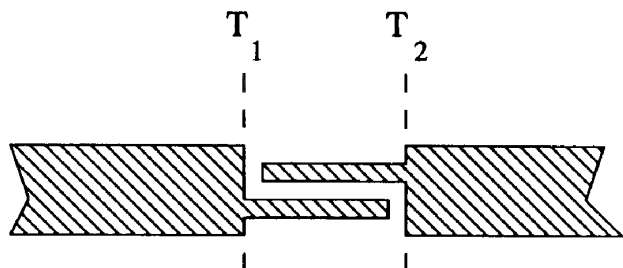


Figure 5.5.21.1: Microstrip Line Finger Break

5.5.22 Microstrip Line Zig-Zag Slit

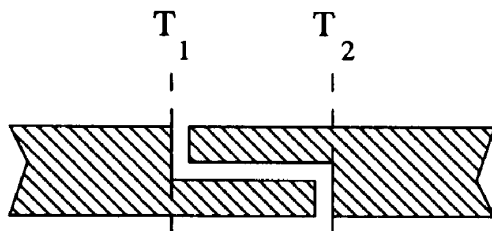


Figure 5.5.22.1: Microstrip Line Zig-Zag Slit

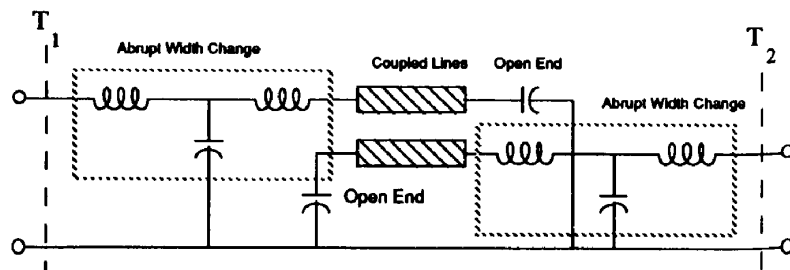


Figure 5.5.22.2: Finger Break and Zig-Zag Slit Model

In [1] data are presented showing that the above two structures can be modeled by the combination of 2 coupled lines, abrupt change in widths, and the end effect capacitances. Lancombe and Cohen [2] reports on experiments with the number of fingers in an interdigital capacitor. It was found that the two-finger configuration resulted in $>$ octave bandwidth (1.4:1 VSWR) whereas the multiple finger configurations did not (0.78 octave).

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5.5.23 Arbitrary Angle Microstrip Line Bend

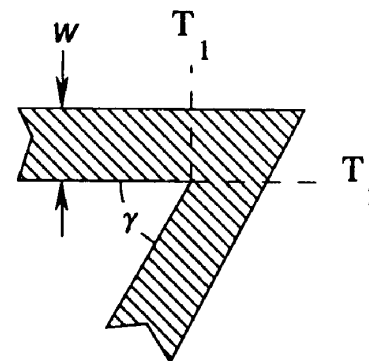


Figure 5.5.23.1: Arbitrary Angle Microstrip Line Bend

The bend in microstrip line is useful whenever we need to modify the signal path to accommodate a packaging or layout scheme. The compensated bend has been described earlier. Mehran [1] gives graphs of s -parameters for the uncompensated structure, but no closed form equations are given. For small angles he found that coupling can be significant. Menzel [2] gives similar curves for the particular case of the 120° bend.

REFERENCES

- [1] Mehran, Reza, "Calculations of Microstrip Bends and Y-Junctions with Arbitrary Angle," *IEEE Transactions on Microwave Theory and Techniques*, Vol. MTT-26, No. 6, June 1978, pp. 400-405.

- [2] Menzel, Wolfgang, "Frequency Dependent Transmission Properties of Microstrip Y-Junctions and 120° Bends," *Microwaves, Optics and Acoustics*, Vol. 2, No. 2, March 1978, pp. 55-59.

5.5.24 Microstrip Line Y Junction

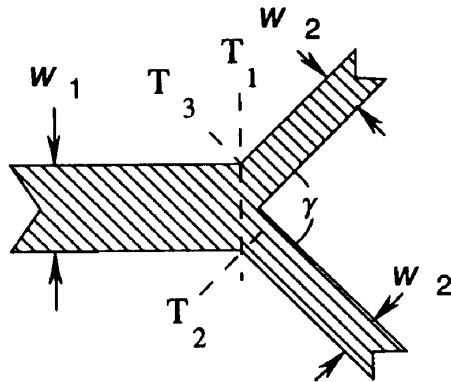


Figure 5.5.24.1: Microstrip Line Y Junction

This structure is useful in design of combiner/dividers. Mehran [2] and Menzel [3] give s -parameter curves for the discontinuity introduced.

Closed-form equations for the radiation of this structure when $w_1 = 2.0 w_2$ have been found by Lewin [1]. The radiated power is

$$P = 60.0 (k h)^2 F \quad (5.5.24.1)$$

where

$$k = \text{free-space wavenumber} \quad (5.5.24.2)$$

$$F = \frac{\epsilon_r + 1.0 - 2.0 c^2}{4.0 c^2} \frac{s}{\sqrt{\epsilon_r - c^2}} \ln \left(\frac{\sqrt{\epsilon_r - c^2} + s}{\sqrt{\epsilon_r - c^2} - s} \right) - \frac{1.0 - c}{1.0 + c} \left[\frac{(\epsilon_r + 1.0)(1.0 - c)(1.0 + 3.0 c)}{4.0 c^2} \right]$$

$$\frac{1.0}{\sqrt{\epsilon_r}} \ln \frac{\sqrt{\epsilon_r + 1.0}}{\sqrt{\epsilon_r - 1.0}} + \frac{(2.0 \epsilon_r - c)}{\Gamma} \ln \frac{\epsilon_r - c + \Gamma}{\epsilon_r - 1.0} \quad (5.5.24.3)$$

and

$$c = \cos \beta \quad (5.5.24.4)$$

$$s = \sin \beta \quad (5.5.24.5)$$

$$\beta = \frac{\gamma}{2.0} \quad (5.5.24.6)$$

$$\Gamma = \sqrt{2.0 \epsilon_r (1.0 - c) - s^2} \quad (5.5.24.7)$$

The equations were found by integration of the far fields of the structure. Coupling between the arms was neglected and can be expected to have an effect for small γ .

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- [3] Menzel, Wolfgang, "Frequency Dependent Transmission Properties of Microstrip Y-Junctions and 120° Bends," *Microwaves, Optics and Acoustics*, Vol. 2, No. 2, March 1978, pp. 55-59.
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5.5.25 Slit in Microstrip Line

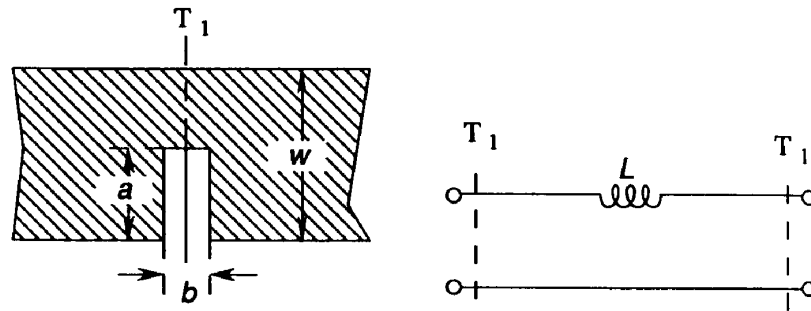


Figure 5.5.25.1: Microstrip Transverse Slit

The microstrip line slit is used for tuning, as a circuit element in filters, and in directional couplers for mode velocity equalization. A simple inductive model (Figure 5.5.25.1) is presented in Hofer [1]. His equations are

$$L = \frac{\mu_0 \pi}{2.0} \left(\frac{a'}{A} \right)^2 \quad (5.5.25.1)$$

where

$$\frac{a'}{A} = 1.0 - \frac{Z_{0,air}}{Z_{0,microstrip}} \quad (5.5.25.2)$$

and $Z_{0,air}$ is the impedance of a microstrip line of width w and dielectric replaced by air, and $Z_{0,microstrip}$ is the impedance of a microstrip line of width $w' = w - a$ and dielectric also replaced by air. The capacitance of the gap which is in parallel with the inductor may be calculated as half the capacitance of a microstrip line gap having dimensions $s = b$ and $w = 2.0 a$ (see Figure 5.5.19.1).

The above equations are valid for

$$b < h$$

$$b \ll \lambda_g$$

$$0.0 \leq a/w \leq 0.9$$

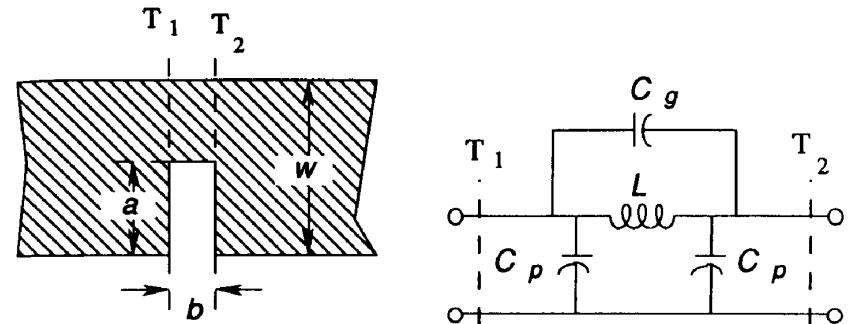


Figure 5.5.25.2: Microstrip Transverse Slit

Mosig [2] presents a more complete model (shown in Figure 5.5.25.2), which includes the capacitance across the slit and to ground. Design curves are given in this reference.

REFERENCES

- [1] Hofer, Wolfgang J., "Equivalent Series Inductivity of a Narrow Transverse Slit in Microstrip," *IEEE Transactions on Microwave Theory and Techniques*, Vol. MTT-25, No. 10, October 1977, pp. 822-824.
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5.5.26 Microstrip Line Spur Line

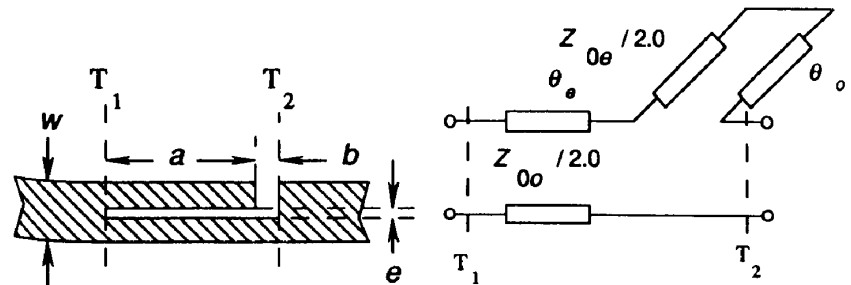


Figure 5.5.26.1: Spur Line

The spur line is a useful filter element first adapted to microstrip lines in Bates [1] for lines having equal mode velocities. This is rarely the case and the analysis of Nguyen [2]

extended this to lines of unequal mode velocities. Nguyen finds a model for the spur line in terms of equivalent transmission lines as shown above. The parameters $Z_{0,e}$, $Z_{0,\rho}$, θ_e , and θ_ρ are calculated with the equations for coupled microstrip lines.

REFERENCES

- [1] Bates, R.N., "Design of Microstrip Spur-Line Band-Stop Filters," *Microwaves, Optics and Acoustics*, Vol. 1, No. 6, November 1977, pp. 209–214.
- [2] Nguyen, C., et al., "Millimeter Wave Printed Circuit Spurline Filters," *1983 IEEE MTT-S Symposium Digest*, pp. 98–100.

5.5.27 Air Bridges

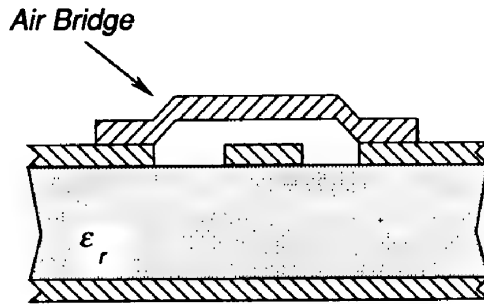


Figure 5.5.27.1: Air Bridge

The air bridge is used in microwave integrated circuits to bridge a signal over other signal conductors. This occurs commonly in spiral inductors. The resulting structure couples energy between the signal paths. Wiemer [2] mentions that he has had success modeling this with the parallel-plate capacitance multiplied by 1.55 to 1.7 to include the effect of fringing or stray capacitances. This is in agreement with the results of Jansen *et al.* [1] which suggest the factor 1.7.

REFERENCES

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5.5.28 Coupled Microstrip Line Right-Angle Bends

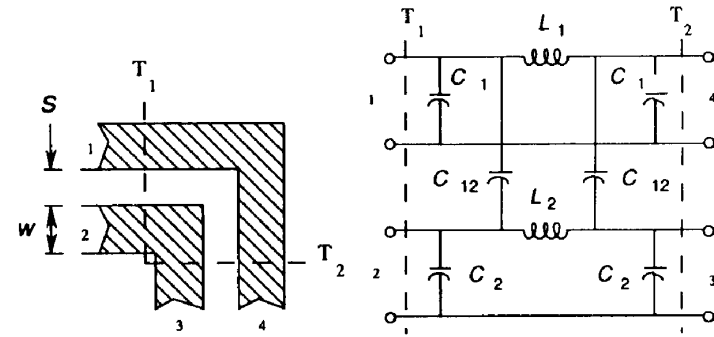


Figure 5.5.28.1: Coupled Microstrip Line Right-Angle Bends

This structure occurs in spiral inductors and is a significant error source in their simulation. Closed-form relations for CAE have been fit to the curves of Hill and Tripathi [2], which apply for $w = s$

$$\frac{C_{11,e}}{C_{11,\infty}h} = -0.18306061 + 5.0447137 \left(\frac{w}{h}\right) \quad (5.5.28.1)$$

$$\frac{C_{12,e}}{C_{12,\infty}h} = -0.039680602 + 2.4009002 \left(\frac{w}{h}\right) \quad (5.5.28.2)$$

$$\frac{C_{22,e}}{C_{22,\infty}h} = -0.12157902 + 1.0334401 \left(\frac{w}{h}\right) \quad (5.5.28.3)$$

$$\frac{L_{11,e}}{L_{11,\infty}h} = -0.50404018 + 4.4984334 \left(\frac{w}{h}\right) \quad (5.5.28.4)$$

$$\frac{L_{22,e}}{L_{22,\infty}h} = -0.53287219 + 0.48572951 \left(\frac{w}{h}\right) \quad (5.5.28.5)$$

$$\frac{L_{12,e}}{L_{12,\infty}h} = -0.74416701 + 0.44005033 \sqrt{\frac{w}{h}} \quad (5.5.28.6)$$

where

$$C_1 = \frac{C_{11,e} - C_{12,e}}{2.0} \quad (5.5.28.7)$$

$$C_2 = \frac{C_{22,e} - C_{12,e}}{2.0} \quad (5.5.28.8)$$

$$C_{12} = \frac{C_{12,e}}{2.0} \quad (5.5.28.9)$$

$$L_1 = L_{11,e} + L_{12,e} \quad (5.5.28.10)$$

$$L_2 = L_{22,e} + L_{12,e} \quad (5.5.28.11)$$

The variables subscripted with infinity (∞) are the parameter values corresponding to the isolated and coupled line without bend. These curve fits are valid to within a few percent for

$$0.2 \leq w/h \leq 3.0$$

Curved bend coupling is described in [1].

REFERENCES

- [1] Abouzahra, Mohamed, and Leonard Lewin, "Theory and Application of Coupling Between Curved Transmission Lines," *IEEE Transactions on Microwave Theory and Techniques*, Vol. MTT-30, No. 11, November 1982, pp. 1988–1995 and "Corrections" Vol. MTT-33, No. 1, January 1985, pp. 74–75.
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5.6 STRIPLINE DISCONTINUITIES

5.6.1 Abrupt Stripline 90° Bend

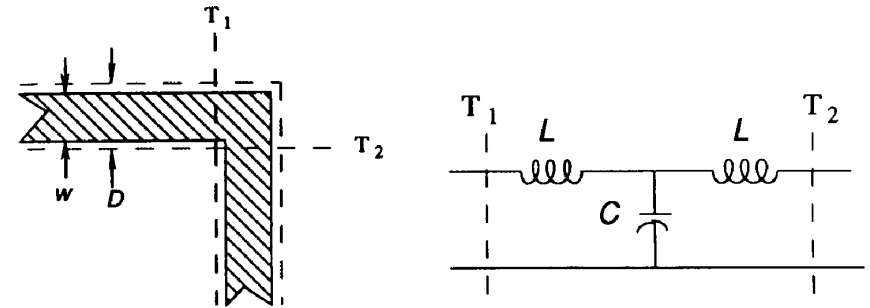


Figure 5.6.1.1: Abrupt Stripline Bend

According to Howe [4], an abrupt bend in stripline can be approximated with 0.2 pF to ground. In Oliner [5], a right-angle bend is modeled as the combination of series inductances and a shunt capacitance

$$L = \frac{2.0 D Z_0 \left[0.878 + 2.0 \left(\frac{D}{\lambda_g} \right)^2 \right]}{\lambda_g \omega} \quad (5.6.1.1)$$

$$C = \frac{2.0 \pi D}{\omega \lambda_g Z_0 \left[1.0 - 0.114 \left(\frac{2.0 D}{\lambda_g} \right)^2 \right]} \quad (5.6.1.2)$$

where

$$D = \frac{b K(k)}{K(k')} \equiv w + \frac{2.0 b}{\pi} \ln 2.0 \quad (5.6.1.3)$$

$$k = \tanh \left(\frac{\pi w}{2 b} \right) \quad (5.6.1.4)$$

and Z_0 is calculated with the stripline equations from Chapter 2.

REFERENCES

- [1] Altschuler, H.M., and A.A. Oliner, "Discontinuities in the Center Conductor of Symmetric Strip Transmission Line," *IRE Transactions on Microwave Theory and Techniques*, Vol. MTT-8, May 1960, pp. 328–339, and "Addendum to 'Discontinuities in the Center Conductor of Symmetric Strip Transmission Line,'" Vol. MTT-10, No. 2, March 1962, p. 143.
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- [4] Howe, Jr., Harlan, *Stripline Circuit Design*, Artech House, Norwood, MA, 1974.
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5.6.2 Arbitrary Angle Stripline Bend

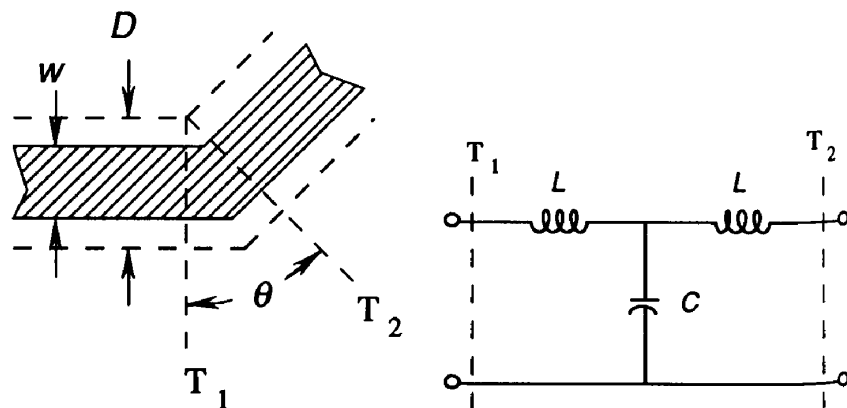


Figure 5.6.2.1: Arbitrary Angle Stripline Bend

Oliner [4] gives

$$L = \frac{2.0 D Z_0}{\omega \lambda_g} \left\{ \Psi \left[-0.5 \left(1.0 - \frac{\theta}{\pi} \right) \right] - \Psi(-0.5) \right\} \quad (5.6.2.1)$$

$$= \frac{2.0 D Z_0}{\omega \lambda_g} \left[\Psi \left(\frac{\theta}{2.0 \pi} - 0.5 \right) - \Psi(-0.5) \right] \quad (5.6.2.2)$$

$$C = \frac{2.0 \pi D}{\lambda Z_0 \omega \cot \left(\frac{\theta}{2.0} \right)} \quad (5.6.2.3)$$

where

$$D = \frac{b K(k)}{K(k')} \approx w + \frac{2.0 b}{\pi} \ln 2.0 \quad (5.6.2.4)$$

$$k = \tanh \left(\frac{\pi w}{2.0 b} \right) \quad (5.6.2.5)$$

θ = bend angle in radians

The psi or digamma function, $\Psi(x)$, is tabulated in [1, Table 6.1]

$\Psi(x)$ = psi or digamma function = logarithmic derivative of gamma function

$$\equiv \frac{\partial [\ln \Gamma(x)]}{\partial x} \quad (5.6.2.6)$$

A closed-form equation was found that is accurate to within 0.001% for $0.5 \leq x \leq 1.0$

$$\Psi(x) \approx 0.144202 + 0.023591 (x + 1.0) + 0.870648 \ln (x + 1.0) \\ - 0.744174 / (x + 1.0) - 1.0 / x \quad (5.6.2.7)$$

REFERENCES

- [1] Abramowitz, Milton, and Irene A. Stegun, ed., *Handbook of Mathematical Functions with Formulas, Graphs, and Mathematical Tables*, Dover Publications, 1972.
- [2] Altschuler, H.M., and A.A. Oliner, "Discontinuities in the Center Conductor of Symmetric Strip Transmission Line," *IRE Transactions on Microwave Theory and Techniques*, Vol. MTT-8, May 1960, pp. 328–339 and "Addendum to 'Discontinuities in the Center Conductor of Symmetric Strip Transmission Line,'" Vol. MTT-10, No. 2, March 1962, p. 143.

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5.6.3 Optimal Stripline Mitered 90° Bend

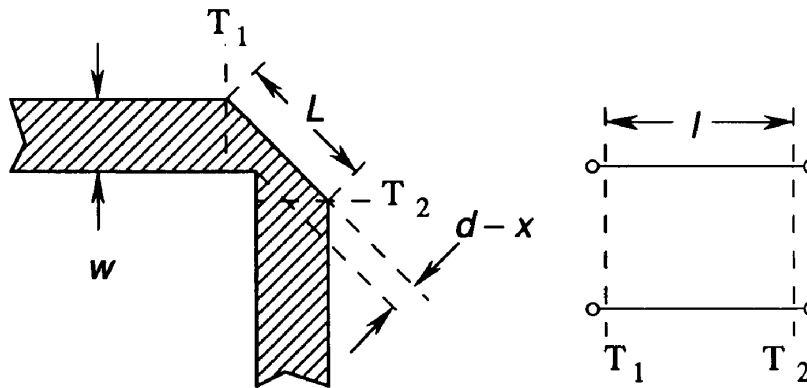


Figure 5.6.3.1: Optimal Stripline Mitered Bend

Mathaei, Young, and Jones [1] give the optimal mitre dimensions for $b/\lambda = 0.0847$ and negligible t . Howe [2] states that in his experience this curve applies over frequencies through C-band. An equation fitting this graph has been derived for CAE use

$$\frac{a}{w} = 1.33068 - 0.38174 \frac{w}{b} + 0.152215 \left(\frac{w}{b}\right)^2 - 0.024477 \left(\frac{w}{b}\right)^3 \quad (5.6.3.1)$$

which fits the curve to within 0.5% for $0.2 \leq w/b \leq 3.0$. The width of the narrow section of trace used for the bevel is

$$d - x = \frac{2.0 w - a}{\sqrt{2.0}} \quad (5.6.3.2)$$

where

$$a = \frac{L}{\sqrt{2.0}} \quad (5.6.3.3)$$

The effective length of the mitre is necessary for modeling and is useful when matching line electrical lengths. Some designers will use very large radius ($> 4 \rightarrow 5 w$) bends when matching electrical lengths rather than use mitered corners. A CAE equation has been derived for this curve also:

$$l/w = \left[1.7854346 - 0.015055579 \frac{w}{b} - 0.11334006 \left(\frac{w}{b}\right)^2 + 0.050432954 \left(\frac{w}{b}\right)^3 - 0.0097883487 \left(\frac{w}{b}\right)^4 \right]^{-1} \quad (5.6.3.4)$$

which fits the curve to within 0.3% for $0.2 \leq w/b \leq 3.0$.

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- [1] G. Mathaei, L. Young and E.M.T. Jones, *Microwave Filters, Impedance-Matching Networks, and Coupling Structures*, Artech House, Norwood, MA, 1980.
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5.6.4 Optimal Stripline Rounded Bend

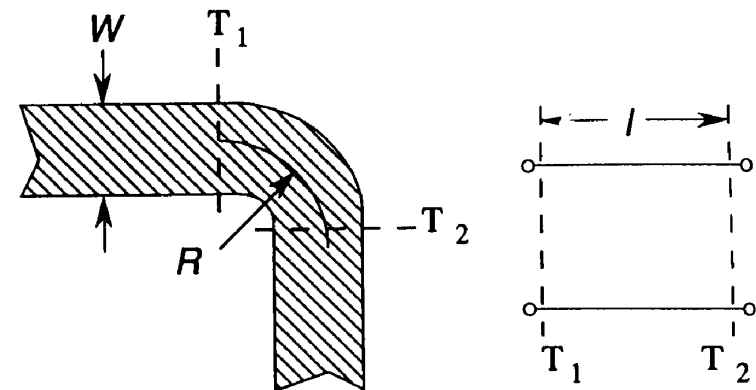


Figure 5.6.4.1: Optimal Stripline Rounded Bend

The rounded bend is an alternative to the mitred corner when a controlled impedance bend is required and space is not a constraint. For best controlled impedance, Howe [1] recommends

$$R > 3 \times w$$

The length of the section of arc can be approximated by its centerline length

$$l \cong \frac{\pi R}{2.0} \quad (5.6.4.1)$$

REFERENCES

- [1] Howe, Jr., Harlan, *Stripline Circuit Design*, Artech House, Norwood, MA, 1974.

5.6.5 Stripline Abrupt Width Change

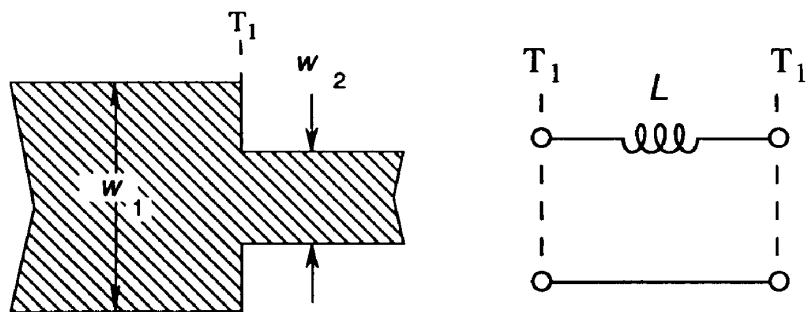


Figure 5.6.5.1: Stripline Abrupt Width Change

References [1, 2, 6] give the most commonly seen relation for the inductance of this discontinuity.

$$L = \frac{2.0 D}{\omega \lambda_g} \ln \left[\csc \left(\frac{\pi Z_{0,2}}{2.0 Z_{0,1}} \right) \right] \quad (5.6.5.1)$$

Which is valid for

$$b / \lambda_g \sim 0$$

$$w_1 \gg w_2$$

The impedances of isolated strips of width w_1 and w_2 are $Z_{0,1}$ and $Z_{0,2}$, respectively. In Nalbandian and Steenaert [5] an equation that has validity over a wider range of w_1 and w_2 is given:

$$L \cong \frac{Z_{0,1} 2.0 D K}{\lambda_g} \quad (5.6.5.2)$$

where

$$D' = \frac{b K(k)}{K(k')} \cong w_1 + \frac{2.0 b}{\pi} \ln 2.0 \quad (5.6.5.3)$$

$$D = \frac{b K(k)}{K(k')} \cong w_2 + \frac{2.0 b}{\pi} \ln 2.0 \quad (5.6.5.4)$$

$$k = \tanh \left(\frac{\pi w}{2.0 b} \right) \quad (5.6.5.5)$$

$$K = \ln \left[\left(\frac{1.0 - \alpha^2}{4.0 \alpha} \right) \left(\frac{1.0 + \alpha}{1.0 - \alpha} \right)^{\frac{\alpha + 1.0}{2.0 \alpha}} \right] + \frac{2.0}{A} \quad (5.6.5.6)$$

$$A = \left(\frac{1.0 + \alpha}{1.0 - \alpha} \right)^{2.0 \alpha} \left[\frac{1.0 + \sqrt{1.0 - D^2 / \lambda^2}}{1.0 - \sqrt{1.0 - D^2 / \lambda^2}} \right] - \frac{1.0 + 3.0 \alpha^2}{1.0 - \alpha^2} \quad (5.6.5.7)$$

$$\alpha = \frac{D'}{D} < 1.0 \quad (5.6.5.8)$$

REFERENCES

- [1] Altschuler, H.M., and A.A. Oliner, "Discontinuities in the Center Conductor of Symmetric Strip Transmission Line," *IRE Transactions on Microwave Theory and Techniques*, Vol. MTT-8, May 1960, pp. 328-339, and "Addendum to 'Discontinuities in the Center Conductor of Symmetric Strip Transmission Line'," Vol. MTT-10, No. 2, March 1962, p. 143.
- [2] Bahl, I.J., and Ramesh Garg, "A Designer's Guide to Stripline Circuits," *Microwaves*, January 1978, pp. 90-96.
- [3] Bhat, Bharathi, and Shiban K. Koul, *Stripline-Like Transmission Lines for Microwave Circuits*, John Wiley & Sons, New York, 1989.

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5.6.6 Compensated Stripline Step Discontinuities

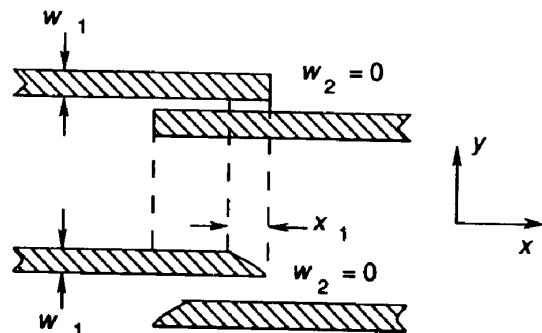


Figure 5.6.6.1: Compensated Step Discontinuities

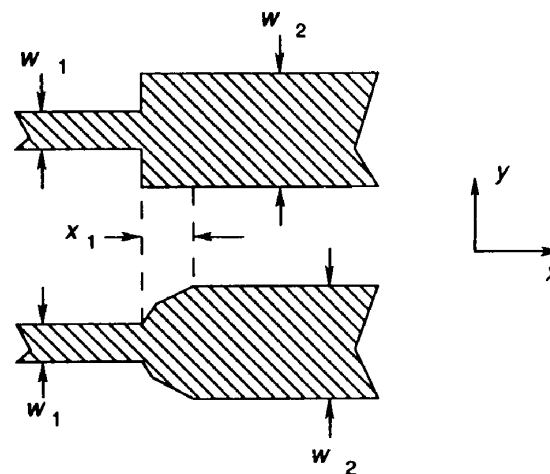


Figure 5.6.6.2: Compensated Step Discontinuities

The presence of a step discontinuity in any TEM line can be compensated for by the addition of a taper near the discontinuity as shown in the figures above (upper figures are uncompensated, lower are compensated). The cross-section is varied continuously to exchange the transmission line's parallel plate capacitance gradually for fringing capacitance, thus keeping the impedance constant until the step is complete.

For stripline, the line taper begins x_1 units back from the nominal position of the step:

$$x_1 = \frac{c}{2} \frac{b}{c} \cosh^{-1} \left(\frac{2w}{c} + 1.0 \right) \quad (5.6.6.1)$$

and the taper profile is given by:

$$y(x) = w - \frac{b}{2} \left\{ \cosh \left[\cosh^{-1} \left(\frac{2w}{c} + 1.0 \right) - \frac{2x}{b} \right] - 1.0 \right\} \quad (5.6.6.2)$$

where

$w = w_1 - w_2$ = the difference between the two trace widths

b = ground plane spacing

$$c = \frac{C_{fe}'}{\epsilon} \quad (5.6.6.3)$$

The calculation can be completed with the equations in the stripline section of this book.

REFERENCES

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5.6.7 Stripline Open End

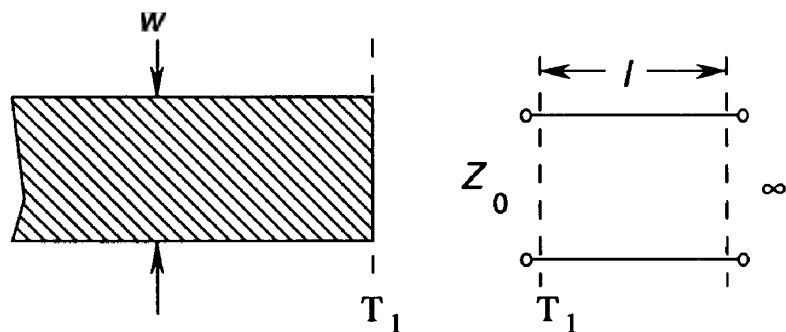


Figure 5.6.7.1: Stripline Open End Effect

This structure occurs in designs, such as filters that have $\lambda/4.0$ open-circuited lines. The open-circuit fringing capacitance increases the electrical length of the line, which must be corrected for proper design of the filter. This fringing effect can be modeled by either a length of transmission line or by a lumped capacitance.

For the infinitely wide strip the additional length is calculated with [1, 5]

$$\Delta l_{\infty} = \frac{b \ln 2.0}{\pi} \quad (5.6.7.1)$$

For the more realistic, finite-width stripline the additional length is [1]

$$\Delta l = \frac{1.0}{\kappa} \cot^{-1} \left[\frac{4.0 \Delta l_{\infty} + 2.0 w}{\Delta l_{\infty} + 2.0 w} \cot (\kappa \Delta l_{\infty}) \right] \quad (5.6.7.2)$$

where

$$\kappa = \frac{2.0 \pi}{\lambda_g} \quad (5.6.7.3)$$

This equation was derived from (5.6.7.1) and measured data. For $\kappa \Delta l_{\infty} \ll 1.0$ (5.6.7.2) can be approximated with

$$\Delta l \equiv \Delta l_{\infty} \left(\frac{\Delta l_{\infty} + 2.0 w}{4.0 \Delta l_{\infty} + 2.0 w} \right) \quad (5.6.7.4)$$

REFERENCES

- [1] Altschuler, H.M., and A.A. Oliner, "Discontinuities in the Center Conductor of Symmetric Strip Transmission Line," *IRE Transactions on Microwave Theory and Techniques*, Vol. MTT-8, May 1960, pp. 328–339 and "Addendum to 'Discontinuities in the Center Conductor of Symmetric Strip Transmission Line,'" Vol. MTT-10, No. 2, March 1962, p. 143.
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5.6.8 Stripline T Junction

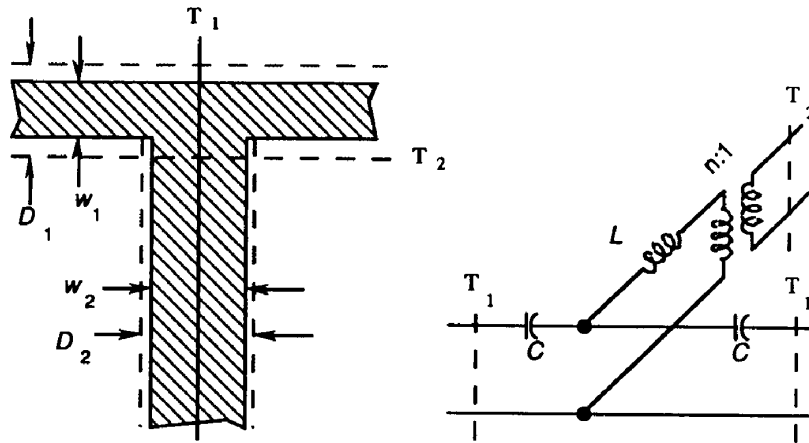


Figure 5.6.8.1: Stripline T Junction

The stripline T junction is described in [1] and [5]. These models and experimental data are compared in [4]. Models are obtained by equating the waveguide T model of [5] to an equivalent parallel plate guide model and then to stripline.

The equivalent circuit of the stripline T junction is calculated [6] with

$$n = n' \sqrt{\frac{D_2}{D}} \quad (5.6.8.1)$$

$$C = \frac{\lambda_g}{\omega D_2 Z_0 (0.785 \eta)^2} \quad (5.6.8.2)$$

for $D_2/D < 0.5$

$$L = \frac{-Z_0 X_a}{2.0 \omega} + \frac{Z_0}{2.0 \eta \omega} \left\{ \frac{B_t}{2.0} + \frac{2.0 D}{\lambda_g} \left[\ln 2.0 + \frac{\pi D}{6.0 D} + \frac{3.0}{2.0} \left(\frac{D}{\lambda_g} \right)^2 \right] \right\} \quad (5.6.8.3)$$

and for $D_2/D > 0.5$

$$L = \frac{-Z_0 X_a}{2.0 \omega} + \frac{2.0 Z_0 D}{\lambda_g \eta^{12} \omega} \left[\ln \left(\frac{1.43 D}{D_2} \right) + 2.0 \left(\frac{D}{\lambda_g} \right)^2 \right]$$

(5.6.8.4)

where

$$n' = \frac{\sin(\pi D_2 / \lambda_g)}{\pi D_2 / \lambda_g} \quad (5.6.8.5)$$

$$B_t = \frac{2.0 D_1}{\lambda_g} \left\{ \ln \left[\csc \left(\frac{\pi D_2}{2.0 D_1} \right) \right] + \frac{1.0}{2.0} \left(\frac{D_1}{\lambda_g} \right)^2 \cos^4 \left(\frac{\pi D_2}{2.0 D_1} \right) \right\} \quad (5.6.8.6)$$

$$\frac{D_1}{D_2} = \frac{Z_2}{Z_1} \quad (5.6.8.7)$$

$$D_1 = \frac{b K(k)}{K(k')} \equiv w + \frac{2.0 b}{\pi} \ln 2.0 \quad (5.6.8.8)$$

$$k = \tanh \left(\frac{\pi w}{2.0 b} \right) \quad (5.6.8.9)$$

REFERENCES

- [1] Altschuler, H.M., and A.A. Oliner, "Discontinuities in the Center Conductor of Symmetric Strip Transmission Line," *IRE Transactions on Microwave Theory and Techniques*, Vol. MTT-8, May 1960, pp. 328-339 and "Addendum to 'Discontinuities in the Center Conductor of Symmetric Strip Transmission Line,'" Vol. MTT-10, No. 2, March 1962, p. 143.
- [2] Bahl, I.J., and Ramesh Garg, "A Designer's Guide to Stripline Circuits," *Microwaves*, January 1978, pp. 90-96.
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5.6.9 Stripline Gap in Center Conductor

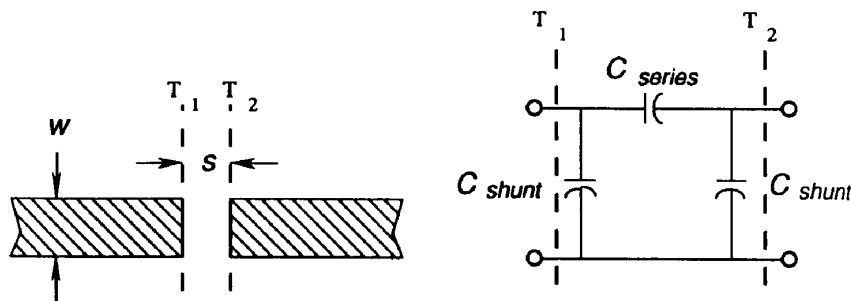


Figure 5.6.9.1: Stripline Gap in Center Conductor

The stripline gap can be used for impedance matching and as a series capacitive element. Two models are possible, depending on the choice of reference plane. The capacitive π model has reference planes at each end of the gap and the values are [4] calculated with

$$C_{series} = \frac{1.0}{2.0 Z_0 \omega} \left[\frac{1.0 + (2.0 Z_0 B_b + Z_0 B_a) \cot \left(\frac{\pi s}{\lambda_g} \right)}{\cot \left(\frac{\pi s}{\lambda_g} \right) - (2.0 Z_0 B_b + Z_0 B_a)} \right] - \frac{B_a}{2.0 \omega} \quad (5.6.9.1)$$

$$C_{shunt} = \frac{Y_0 + B_a \cot \left(\frac{\pi s}{\lambda_g} \right)}{\omega \cot \left(\frac{\pi s}{\lambda_g} \right) - \omega B_a / Y_0} \quad (5.6.9.2)$$

for $s \gg w$

$$C_{shunt} \approx \frac{\tan \left(\frac{2.0 b \ln 2.0}{\lambda_g} \right)}{Z_0 \omega} \quad (5.6.9.3)$$

where B_a and B_b are the center-line model elements calculated with

$$B_a = \frac{-2.0 b \ln \left[\cosh \left(\frac{\pi s}{2.0 b} \right) \right]}{\lambda_g} \quad (5.6.9.4)$$

$$B_b = \frac{b \ln \left[\coth \left(\frac{\pi s}{2.0 b} \right) \right]}{\lambda_g} \quad (5.6.9.5)$$

Oliner, in fact, indicates that the approximation for C_{shunt} should be used for $s > w$.

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- [1] Altschuler, H.M., and A.A. Oliner, "Discontinuities in the Center Conductor of Symmetric Strip Transmission Line," *IRE Transactions on Microwave Theory and Techniques*, Vol. MTT-8, May 1960, pp. 328-339 and "Addendum to 'Discontinuities in the Center Conductor of Symmetric Strip Transmission Line'," Vol. MTT-10, No. 2, March 1962, p. 143.
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5.6.10 Stripline Round Hole in Center Conductor

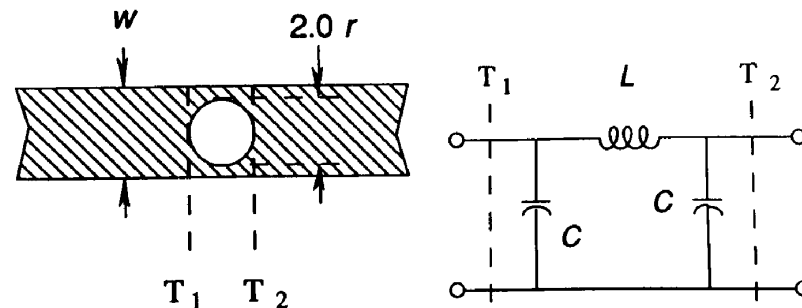


Figure 5.6.10.1: Stripline Round Hole

The model for a hole in stripline is given in Oliner [4]

$$C = \frac{1.0 + (B_a Z_0) \cot \left(\frac{\pi d}{\lambda_g} \right)}{\omega Z_0 \left[\cot \left(\frac{\pi d}{\lambda_g} \right) - B_a Z_0 \right]} \quad (5.6.10.1)$$

$$L = \frac{-1.0}{\omega B_B Z_0} \quad (5.6.10.2)$$

where B_a and B_b are the model elements of the center-line model

$$B_a = \frac{1.0}{4.0 B_B Z_0^2} \quad (5.6.10.3)$$

$$B_b = -\frac{3.0}{4.0 \pi Z_0} \frac{\lambda_g b D}{d^3} \quad (5.6.10.4)$$

REFERENCES

- [1] Altschuler, H.M., and A.A. Oliner, "Discontinuities in the Center Conductor of Symmetric Strip Transmission Line," *IRE Transactions on Microwave Theory and Techniques*, Vol. MTT-8, May 1960, pp. 328-339 and "Addendum to 'Discontinuities in the Center Conductor of Symmetric Strip Transmission Line'," Vol. MTT-10, No. 2, March 1962, p. 143.
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5.6.11 Stripline Slot in Center Conductor

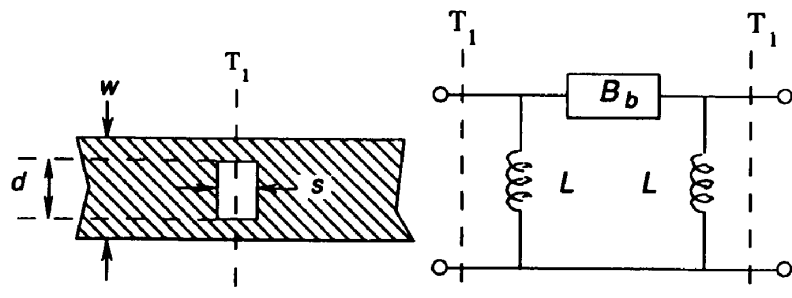


Figure 5.6.11.1: Stripline Slot

The model parameters for the stripline slot can be calculated with

$$L = \frac{1.0 + P}{\frac{-2.0 b \omega}{\lambda_g} \ln \left(\cosh \frac{\omega \tau}{b} \right)} \quad (5.6.11.1)$$

$$B_b' = \frac{b}{\lambda_g} \ln \left(\sinh \frac{w \tau}{b} \right) - \frac{B_a'}{2.0} - \frac{b \lambda_g}{w^2} \left(1.0 + \frac{2.0 b}{\pi w} \ln 2.0 \right) Q \quad (5.6.11.2)$$

where

$$Q \equiv \frac{1.0}{4.0 \tau} \left[\frac{-1.0}{\ln \beta} - \left(\frac{1.0 - \beta^2}{\ln \beta} \right)^2 \sum_{n=1.0}^N f(2.0 n \tau) \frac{X_n^2}{n} \right] \quad (5.6.11.3)$$

$$P \equiv \frac{1.0}{4.0} \left[\left(\frac{1.0 - \beta^2}{\ln \beta} \right)^2 \sum_{n=1.0}^{5.0} g(2.0 n \tau) \frac{X_n^2}{n} + \tau^2 Q \right] \quad (5.6.11.4)$$

$$N = \text{int} \left[(0.7 / \tau - 1.0) + 0.5 \right] \quad (5.6.11.5)$$

$$\tau = \frac{\pi s}{2.0 w} \quad (5.6.11.6)$$

$$\beta = \cos \frac{\pi d}{2.0 w} \quad (5.6.11.7)$$

and

$$X_1 = 1.0 \quad (5.6.11.8)$$

$$X_2 = -1.0 + 3.0 \beta^2$$

$$X_3 = 1.0 - 8.0 \beta^2 + 10.0 \beta^4$$

$$X_4 = -1.0 + 15.0 \beta^2 - 45.0 \beta^4 + 35.0 \beta^6$$

$$X_5 = 1.0 - 24.0 \beta^2 + 126.0 \beta^4 - 224.0 \beta^6 + 126.0 \beta^8$$

Curve fits to the functions $f(x)$ and $g(x)$ of Figure 8 in [1] were derived for calculation purposes

$$f(x) \equiv -0.924663 + 0.244265 x + 0.0130211 / x + 1.62211 e^{-x} \quad (5.6.11.9)$$

$$g(x) \equiv \exp (0.70411622 - 0.26293172 x + 0.01308667 x^2) \quad (5.6.11.10)$$

It is difficult to determine the fit accuracy due to the small size of the original figures. It is believed that $f(x)$ and $g(x)$ are accurate to better than 5% and 1%, respectively, over the required ranges. Model equations are valid for

$$\tau \geq 0.15$$

$$d / w \geq 0.25$$

REFERENCES

- [1] Altschuler, H.M., and A.A. Oliner, "Discontinuities in the Center Conductor of Symmetric Strip Transmission Line," *IRE Transactions on Microwave Theory and Techniques*, Vol. 8, May 1960, pp. 328–339 and "Addendum to "Discontinuities in the Center Conductor of Symmetric Strip Transmission Line," Vol. MTT-10, No. 2, March 1962, p. 143.

5.6.12 Stripline Multistrip Junction

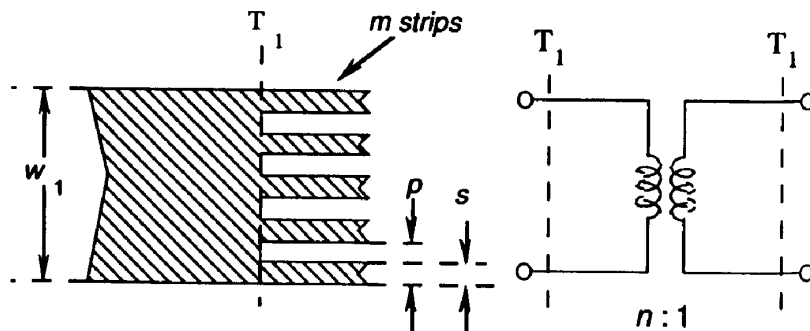


Figure 5.6.12.1: Multi-strip Junction

Multiple slots are described in [1]. Design equations are

$$n = \sqrt{\frac{Z_{0,m}}{Z_0}} \quad (5.6.12.1)$$

where

$$\frac{Z_{0,m}}{Z_0} = \frac{Z_{0,\infty} + A}{1.0 + A} \quad (5.6.12.2)$$

$$A = \frac{K(k')}{K(k)} \frac{p}{m b} \quad (5.6.12.3)$$

$$k = \tanh \left(\frac{\pi s}{2.0 b} \right) \quad (5.6.12.4)$$

$$k' = \sqrt{1.0 - k^2} \quad (5.6.12.5)$$

$$\frac{Z_{0,\infty}}{Z_0} = 1.0 + \frac{2.0 p}{\pi b} \ln \left[\csc \left(\frac{\pi s}{2.0 p} \right) \right] \quad (5.6.12.6)$$

These equations are valid for number of strips $m > 3$ with increasing accuracy with increasing m .

REFERENCES

- [1] Altschuler, H.M., and A.A. Oliner, "Discontinuities in the Center Conductor of Symmetric Strip Transmission Line," *IRE Transactions on Microwave Theory and Techniques*, Vol. MTT-8, May 1960, pp. 328–339 and "Addendum to 'Discontinuities in the Center Conductor of Symmetric Strip Transmission Line,'" Vol. MTT-10, No. 2, March 1962, p. 143.

5.6.13 Stripline Slot in Ground Plane

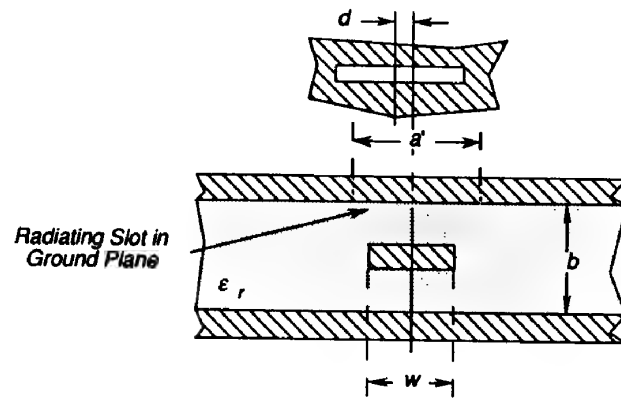


Figure 5.6.13.1: Stripline Slot in Ground Plane

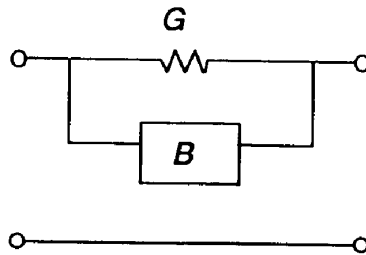


Figure 5.6.13.2: Equivalent Circuit

For the slot in the ground plane of a stripline, Oliner [2] gives the equivalent circuit

$$\frac{G}{Y_0} = \frac{16.0}{3.0 \pi} \frac{K(k')}{K(k)} \frac{a'}{\lambda} \left[1.0 - 0.374 \left(\frac{a'}{\lambda} \right)^2 + 0.130 \left(\frac{a'}{\lambda} \right)^4 \right] \quad (5.6.13.1)$$

Breithaupt [1] returned to Oliner's original equations and included offset (off-center) slots by adding a factor $(I_0/I)^2$

$$\frac{G}{Y_0} = \frac{16.0}{3.0 \pi} \left(\frac{I_0}{I} \right)^2 \frac{K(k')}{K(k)} \frac{a'}{\lambda} \left[1.0 - 0.374 \left(\frac{a'}{\lambda} \right)^2 + 0.130 \left(\frac{a'}{\lambda} \right)^4 \right] \quad (5.6.13.2)$$

where

$$I = \int_{a'/2 - d}^{-a'/2 - d} \frac{\cos(\pi(x+d)/a') dx}{\sqrt{1.0 + k'^2 \sinh^2(\pi x/b)}} \quad (5.6.13.3)$$

$$I_0 = I \text{ with } d = 0.0 \quad (5.6.13.4)$$

The distance between the centerline of the stripline and the aperture is d .

REFERENCES

- [1] Breithaupt, Robert W., "Conductance Data for Offset Series Slots in Stripline," *IEEE Transactions on Microwave Theory and Techniques*, Vol. MTT-16, No. 11, November 1968, pp. 969-970. (Modifies Oliner to include off-center slots.)
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5.7 SUSPENDED MICROSTRIP LINE

5.7.1 Abrupt Width Change and Open End

Although the references [1, 2] do not give any closed-form relations, it is heartening to note that work has begun in the characterization of discontinuities of lines other than stripline and microstrip line.

REFERENCES

- [1] Achkar, J., *et al.*, "Frequency-Dependent Characteristics of an Asymmetric Suspended Stripline and its Open End Discontinuity for Millimeter Wave Applications," *19th European Microwave Conference Proceedings*, September 4–7, 1990, London, England, pp. 767–772.
- [2] Hong, Jia Sheng, *et al.*, "Exact Computation of Generalised Scattering Matrix of Suspended Microstrip Step Discontinuity," *Electronics Letters*, Vol. 25, No. 5, March 2, 1989, pp. 335–336.

5.7.2 Asymmetric Gap

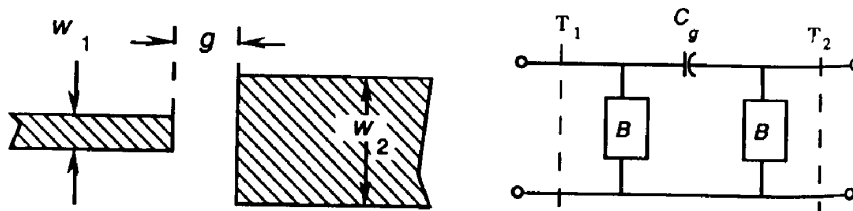


Figure 5.7.2.1: Suspended Substrate Asymmetric Gap

Koster and Jansen [1] give design curves for the circuit parameters of this configuration. Values of the discontinuity are given for

$$0.15 \leq g/h \leq 3.0$$

$$0.1 \leq w_2/w_1 \leq 10.0$$

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5.8 FINLINE DISCONTINUITIES

Finline discontinuities are not covered in this book, however the references for fin-line in Chapter 3 and the following references may provide a good starting point.

Discontinuity	References
Double steps	[2, 5]
Gaps in ground	[2, 4]
Open circuits	[1, 2]
Short circuits	[1, 2]
Steps	[2, 5, 8, 10]
Taper	[9, 10]
Transition between unilateral and bilateral fin-line	[3]
Transverse slits	[6–8]

REFERENCES

- [1] Burton, Miles, and Wolfgang J.R. Hoefer, "An Improved Model for Short- and Open-Circuited Series Stubs in Fin Lines," *1984 IEEE MTT-S Symposium Digest*, pp. 3330–332.
- [2] El Hennaway, Hadia, and Klaus Schünemann, "New Structures for Impedance Transformation in Fin-Lines," *1982 IEEE MTT-S Symposium Digest*, pp. 198–200.
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- [4] El Hennaway, H., *et al.*, "Impedance Transformation in Fin-Line," *IEEE Proceedings*, Part H, Vol. 129, No. 6, April 1983, pp. 342–350.
- [6] Meier, Paul J., "Integrated Fin-Line Millimeter Components," *1981 IEEE MTT-S Symposium Digest*, pp. 195–197.
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- [10] Saad, A.M.K., "Analysis of Fin-Line Tapers and Transitions," *IEEE Proceedings*, Part H, Vol. 130, No. 3, April 1983, pp. 230–235.

5.9 OTHER DISCONTINUITIES

5.9.1 Plated-Through Hole or Via

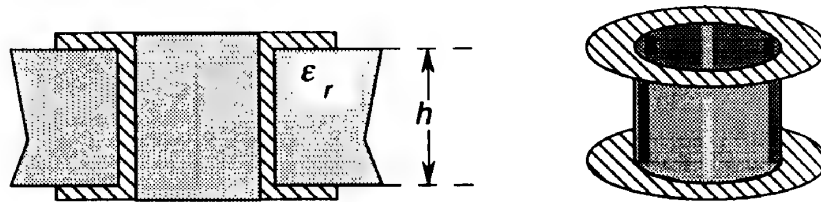


Figure 5.9.1.1: Plated-Through Hole

A via is used whenever we wish a signal to change layers in a two-sided or multilayer circuit. A *blind* or *buried* via occurs when the via is completely internal to the substrate and never reaches the surface. Vias may be created by drilling and plating, by an eyelet, or a wire.

There are three different cases commonly encountered for the plated-through hole or via. First, the via is used in series with the signal line when changing layers as shown in Figure 5.9.1.2. Second, the via can be used to connect surface mount components to a ground plane on another layer (Figure 5.9.1.3). Finally, vias are used to tie together ground planes on different layers of a PCB (Figure 5.9.1.4). When the via is in series with the signal path, we would like it to have the same characteristic impedance as the transmission lines it is interconnecting. When it is used for grounding, we want it to have as low an impedance as possible.

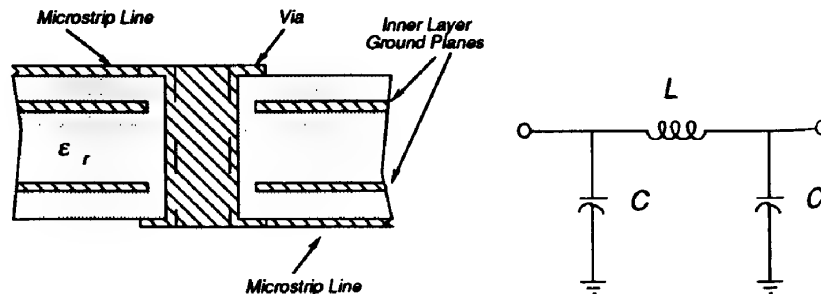


Figure 5.9.1.2: Via Connecting Microstrip Lines

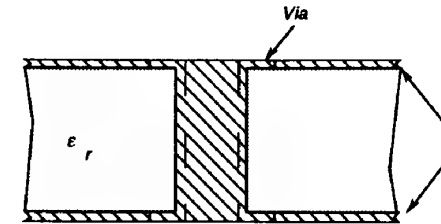


Figure 5.9.1.3: Via to Ground Plane

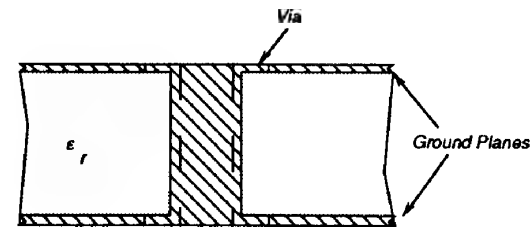


Figure 5.9.1.4: Via Connecting Ground Planes

The simplest model for vias is simply a round wire or cylindrical wire inductor. Equations of Chapter 6 can be used to calculate these inductances [8, 11, and others]. As frequency increases, the model may need to include resistance and inductance caused by current restriction at the transition from the planar line to the via, radiation losses, and stray capacitances to nearby ground planes and across the via.

The effect of the restriction of current at the transition from planar line to via is analyzed in Abbas and Weeks [1]. Design data is presented in a group of tables. Brewitt-Taylor [2] calculates the radiation resistance and a mutual impedance between two vias in the same substrate. As might be suspected, the via is found to radiate more than an equivalent length microstrip trace. The radiation resistance can be approximated with

$$R = \frac{k_0^3 h^3 Z_0 (1.0 - 1.0 / \epsilon_r)}{2.0 \epsilon_r} \quad (5.9.1.1)$$

where

$$k_0 = \frac{\omega}{c_0} \quad (\text{m}^{-1}) = \text{wave number}^1 \text{ in free space} \quad (5.9.1.2)$$

¹The wave number is the number of wavelengths per unit length.

Kasten, *et al.* [6] found model parameters for the via configuration of Figure 5.9.1.2 with a 0.0625 in long via and inner ground planes spaced 0.0056 in from the microstrip lines. The measured values of C and L were approximately 0.5 pF and 3.4 nH. This lumped model was valid to approximately 2 GHz.

In Finch and Alexopoulos a planar waveguide model is used to analyze metallic and dielectric posts using a new rapidly convergent formulation. Quine, *et al.* model semiconductor vias. The coupling between vias is also examined.

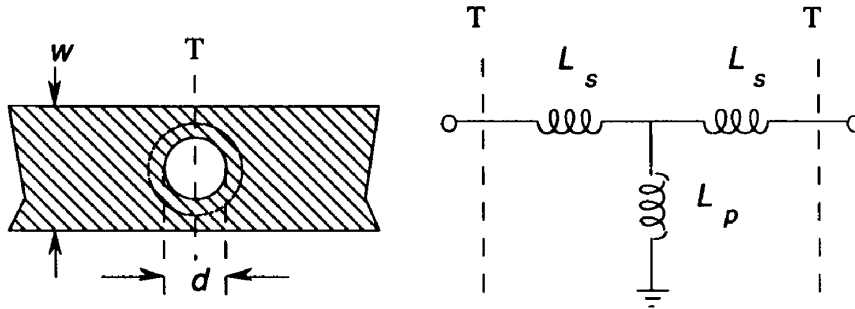


Figure 5.9.1.5: Alternative Via Model

A model for a via to ground in a microstrip line is shown in Figure 5.9.1.5 above. For $d \ll w$, Hoffman [5] states that the via is modeled well by L_p the inductance of a cylindrical wire (see Chapter 6). As d increases, the effect of L_s becomes important. Hoffman presents graphs of the equivalent circuit element values for this structure which are included below. The figures give the value of the model elements for a range of w/h and d/w and $\epsilon_r = 9.8$.

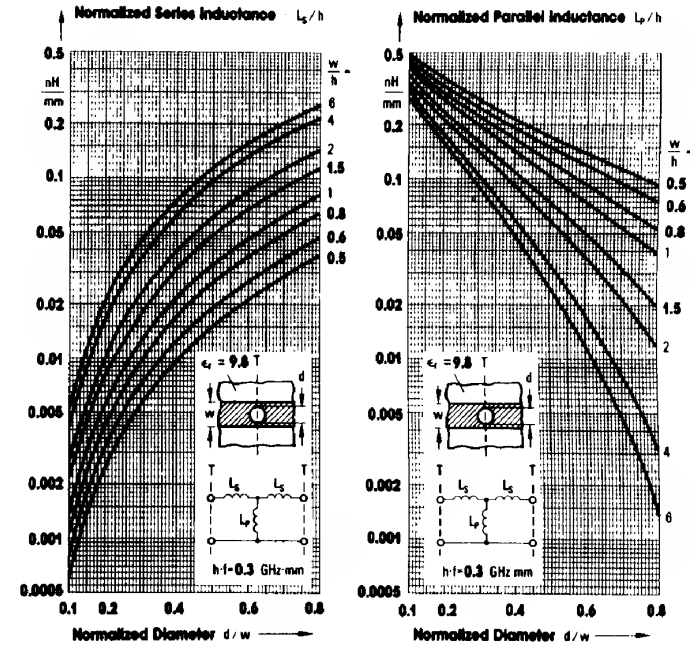


Figure 5.9.1.6: [5, Figure 10.20²]

²Figures from Hoffman, Reinmut K., *Handbook of Microwave Integrated Circuits*, Artech House, Norwood, MA, 1987. Reprinted by permission of the publisher.

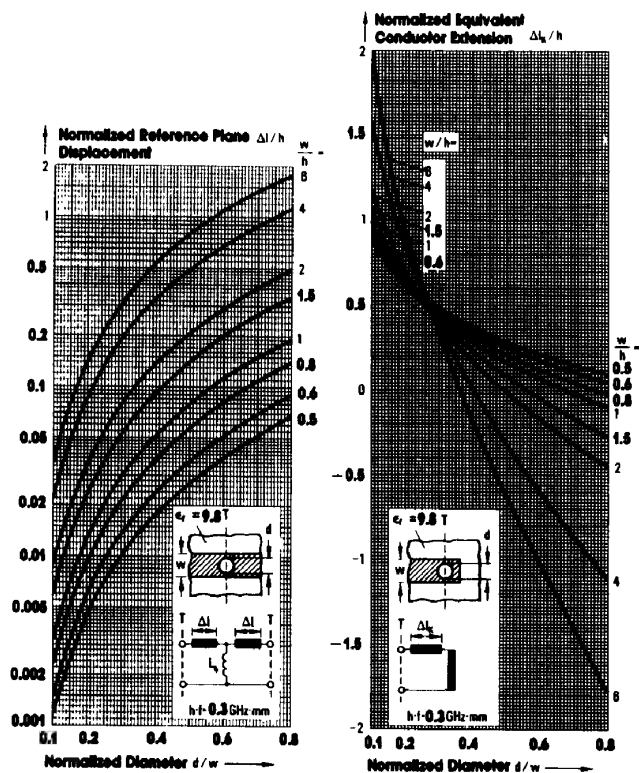


Figure 5.9.1.7: [5, Figure 10.21]

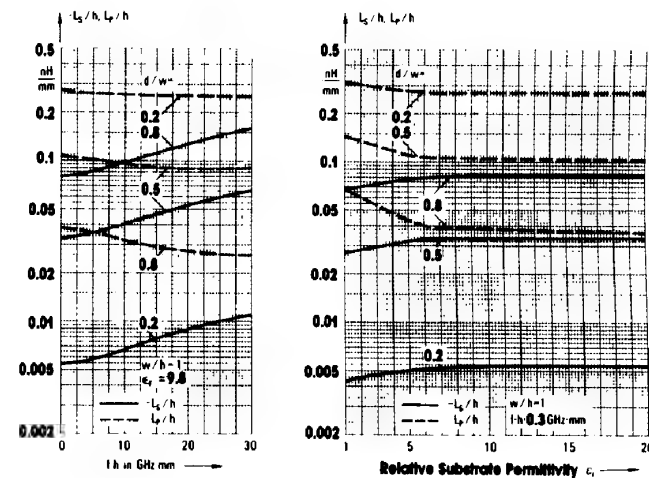


Figure 5.9.1.8: [5, Figure 10.22]

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- [3] Djordjevic, A.R., and T.K. Sarkar, "Computation of Inductance of Simple Vias Between Two Striplines Above a Ground Plane," *IEEE Transactions on Microwave Theory and Techniques*, Vol. MTT-33, No. 3, March 1985, pp. 265–269. (Models via as a flat wire connecting the two striplines. No closed-form equations.)

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5.9.2 Optimal Plated-Through Hole or Via

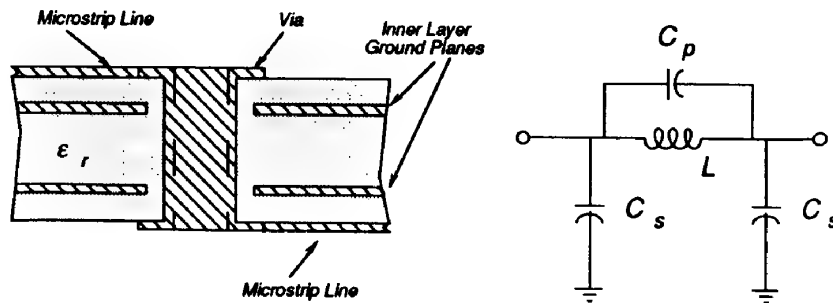


Figure 5.9.2.1: Optimal Via

As described in Wang, *et al.* [1], it is possible to design a via for a desired impedance. The discontinuity is made up of a series inductive tube shunted by a capacitance. Some of the capacitance is due to the open ends on the outer layers. Inner layer ground planes create a capacitance to ground along the inductor. It may be possible to minimize any discontinuity (over some bandwidth) by equating these to the distributed inductance and capacitance of a transmission line with the desired characteristic impedance.

Wang gives design charts which allow one to synthesize arbitrary impedance (matched) vias passing through a ground plane from wires on either side.

REFERENCES

- [1] Wang, Taoyun, *et al.*, "Quasi-static Analysis of a Microstrip Via Through a Hole in a Ground Plane," *IEEE Transactions on Microwave Theory and Techniques*, Vol. MTT-36, No. 6, June 1988, pp. 1007–1013.
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5.9.3 Wraparound Ground

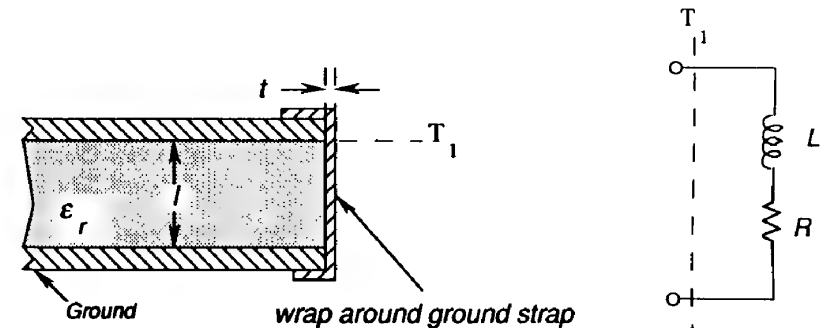


Figure 5.9.3.1: Wraparound Ground

On some circuits the ground plane is extended around the edges of the board. This can be done with a strap made of a thin piece of metal or by plating a slot in the dielectric. Components may also be strapped to the ground plane with this technique through a rectangular opening in the dielectric. This discontinuity may be modeled with the equations for the flat inductor found in Chapter 6.

Chapter 6

Inductors

In this chapter closed-form equations for inductor design are presented. We have already seen a number of techniques for creating inductive reactances. The incremental inductance of any of the structures in Chapter 3 may be calculated from Z_0 using the equations of Chapter 2. The techniques described in Chapter 2 can make inductive elements from open- and short-circuited lines. Many of the discontinuities in Chapter 5 are used as inductors. In this chapter we describe the type of structure that is optimized for its inductance. These equations enable the calculation of the inductance from physical dimensions.

Although the accuracy of these equations is quite good in calculating the inductance, the impedance of the inductor is made up of the resistance and stray capacitance as well. Since these strays are not included the equations can only provide a starting point for design. Typically several iterations of measurement and adjustment are required before the desired Q and resonance-free inductance are achieved over the frequency range. Once the dimensions are determined, the inductance is found to be repeatable.

This chapter includes many useful configurations; consult the references, particularly Grover [2] for others.

6.1 ABCD MATRICES

Series inductor:

$$\begin{bmatrix} 1.0 & j \omega L \\ 0.0 & 1.0 \end{bmatrix}$$

Shunt inductor:

$$\begin{bmatrix} 1.0 & 0.0 \\ \frac{-j}{\omega L} & 1.0 \end{bmatrix}$$

Ideal transformer:

$$\begin{bmatrix} \frac{n_1}{n_2} & 0.0 \\ 0.0 & \frac{n_2}{n_1} \end{bmatrix}$$

where n_1 and n_2 are number of primary and secondary turns, respectively.

6.2 WIRE INDUCTORS

6.2.1 Round Wire Inductor

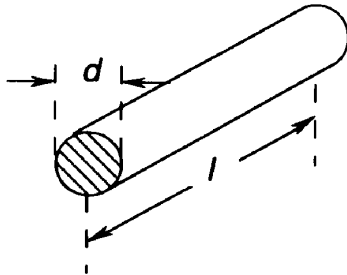


Figure 6.2.1.1: Straight Round Wire

The self-inductance of a straight wire at high frequency due to Grover [2] is:

$$L = 0.002 l \left[\ln \left(\frac{4.0 l}{d} \right) - 1.00 + \frac{d}{2.0 l} + \frac{\mu_r T(x)}{4.0} \right] \quad (\mu\text{H}) \quad (6.2.1.1)$$

where

$$r = d / 2.0 \quad (\text{cm})$$

$$l = \text{wire length} \quad (\text{cm})$$

$$f = \text{frequency} \quad (\text{Hz})$$

The function T corrects the inductance for ac effects. Values of T are tabulated in Grover [2, Table 52, p. 266] and in [5, Table 8, p. 282] where the variable δ is $T / 4$. Grover's tables are more comprehensive and precise. For convenience, we have developed an

approximation for T by curve fitting the data. It is accurate to a few percent over the range $0.0 \leq x \leq 100.0$. For large x , $T(x)$ approaches 0.0.

$$T(x) \equiv \sqrt{\frac{0.873011 + 0.00186128 x}{1.0 + -0.278381 x + 0.127964 x^2}} \quad (6.2.1.2)$$

where

$$x = 2.0 \pi r \sqrt{\frac{2.0 \mu f}{\sigma}} \quad (6.2.1.3)$$

The impedance of curved bond wires is found in Chapter 3.

REFERENCES

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6.2.2 Flat Wire Inductor

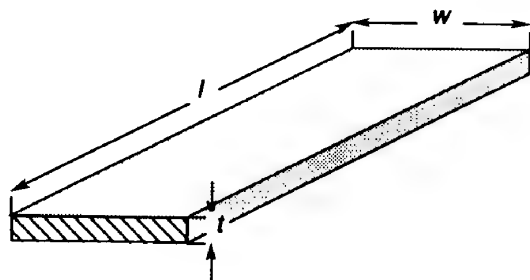


Figure 6.2.2.1: Flat Wire Inductor

Flat or ribbon wires are used as bond wires. Flat wire provides a low-inductance, low-resistance technique for interconnections. A similar option is mesh. A mesh provides the added advantage of being perforated at regular distances. This can be used to provide a more repeatable bond length. The inductance of circuit traces can also be modeled with this structure.

The equation below is often quoted in the literature as coming from [7]:

$$L = 2.0 \times 10^{-3} \times l \left[\ln \left(\frac{l}{w+t} \right) + 1.193 + 0.2235 \left(\frac{w+t}{l} \right) \right] \quad (\text{nH}) \quad (6.2.2.1)$$

Terman, in fact, gives:

$$L = 2.0 \times 10^{-3} \times l \left[\ln \left(\frac{2.0 l}{w+t} \right) + 0.5 + 0.2235 \left(\frac{w+t}{l} \right) \right] \quad (\text{nH}) \quad (6.2.2.2)$$

where l , w , t are in cm. This is the same equation with the factor of 2.0 removed from the numerator of the $\ln()$ term and added to the constant 0.5 giving 1.193 ($0.5 + \ln 2.0$). This is mentioned because the above equations are found in the literature in many different but related forms.

The factor 0.2235 that appears above in fact varies between 0.22313 for $t = 0$ and 0.22352 for $t = w$. A frequency-corrected equation for inductance can be found in Greenhouse [3]

$$L = 0.002 l \left[\ln \left(\frac{2.0 l}{w+t} \right) + 0.25049 + \left(\frac{w+t}{3.0 l} \right) + \frac{\mu T(x)}{4} \right] \quad (6.2.2.3)$$

Again, $T(x)$ can be found with the closed-form relation (6.2.1.2). The resistance is found with:

$$R = \frac{K_c R_s l}{2.0 (w+t)} \quad (\Omega) \quad (6.2.2.4)$$

The surface resistivity R_s is given in Chapter 2. The factor K_c is a current crowding factor found in Caulton [1, Figure 2] (original reference is Terman [4]). The curve has been fit with a closed-form equation:

$$K_c = 1.048312 + 0.21768001 \ln \left(\frac{w}{t} \right) + 0.77176189 e^{-w/t} \quad (6.2.2.5)$$

The fit is valid to within $\pm 0.5\%$ over the range

$$1.0 \leq w/t \leq 100.0$$

REFERENCES

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- [8] U.S. Department of Commerce, Circular of the National Bureau of Standards C74, *Radio Instruments and Measurements*, "Calculation of Inductance," January 1, 1937.

- [9] U.S. Department of Commerce and Labor, *Miscellaneous Pamphlets on Inductance*, U.S. Bureau of Standards, 1920. (A tremendously comprehensive collection of inductance equations unavailable anywhere else.)

6.2.3 Cylinder or Round Tubular Inductor

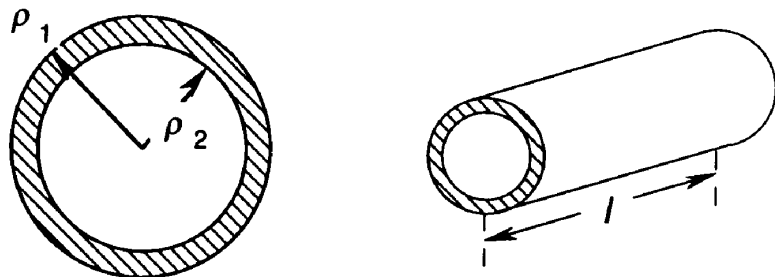


Figure 6.2.3.1: Tubular Inductor

The tubular inductor closely approximates the structure of a plated-through hole or via. The ratio of the ac resistance to dc resistance of a tubular conductor is more constant than for an equivalent solid conductor. This is useful where precise frequency response is to be maintained. Also see Chapter 5 for more information on the via as a discontinuity.

The equation of Grover [2] is

$$L = 0.002 l \left[\ln \left(\frac{2.0 l}{\rho_1} \right) + \ln \zeta - 1.0 \right] \quad (\mu\text{H}) \quad (6.2.3.1)$$

where

l, ρ_1, ρ_2 are in cm

The value of $\ln \zeta$ is in Grover [2, Table 4]. A polynomial has been fit for convenience:

$$\begin{aligned} \ln \zeta = & 0.25009128 - 0.0017049618 \left(\frac{\rho_2}{\rho_1} \right) - 0.51598981 \left(\frac{\rho_2}{\rho_1} \right)^2 \\ & + 0.37420782 \left(\frac{\rho_2}{\rho_1} \right)^3 - 0.10669571 \left(\frac{\rho_2}{\rho_1} \right)^4 \end{aligned} \quad (6.2.3.2)$$

The polynomial fit is within $\pm 0.3\%$ of the tabulated data for

$$0.0 \leq \frac{\rho_2}{\rho_1} \leq 1.0$$

REFERENCES

- [1] Djordjević, Antonije R., *et al.*, "Inductance of Perfectly Conducting Foils Including Spiral Inductors," *IEEE Transactions on Microwave Theory and Techniques*, Vol. MTT-38, No. 10, October 1990, pp. 1407–1414.
- [2] Grover, F.W., *Inductance Calculations*, Van Nostrand, Princeton, NJ, 1946; reprinted by Dover Publications, NY, 1962. (A concise book based on his work at NBS.)
- [3] U.S. Department of Commerce and Labor, *Miscellaneous Pamphlets on Inductance*, U.S. Bureau of Standards, 1920. (A tremendously comprehensive collection of inductance equations unavailable anywhere else.)

6.3 SOLENOIDAL OR HELICAL INDUCTORS

6.3.1 Round or Square Helical Coil Inductor

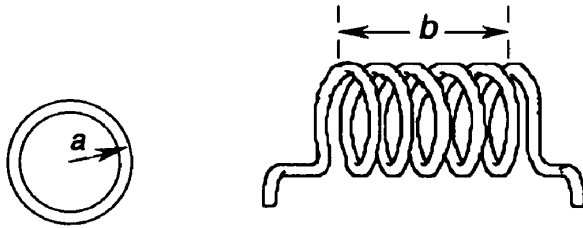


Figure 6.3.1.1: Round Helical Coil

Helical coils are conveniently wound on drill bits or other metal rod of specified diameter. If two wires are wound together on the form and one removed, a fixed distance between turns is easily achieved. When insulated wire is used, the insulation often determines the wire spacing. Spacing can also be checked in manufacturing using a spark plug style thickness gauge. Small PTFE spacers or solder mask on the PCB prevents shorting to any traces or ground plane near the coil and ensures repeatable positioning.

$$L = \frac{\mu_0 n^2 \pi a^2}{b} K_N \quad (6.3.1.1)$$

where

$$K_N = \frac{(d b / a^2) [F(k) - E(k)] + (4.0 d / b) E(k) - (8.0 a / b)}{3.0 \pi} \quad (6.3.1.2)$$

$$k = \frac{2.0 a}{d} \quad (6.3.1.3)$$

$F(k)$ = complete elliptic integral of the first kind

$E(k)$ = complete elliptic integral of the second kind

The radius, a , is measured from the helix center to the center of the wire. Nagaoka's constant K_N is an adjustment factor that corrects for the effects of the coil's length. Wheeler gives an approximation formula to avoid the use of complete elliptic integral

functions in the above; however as Miller [5] points out these are quite simple to calculate to any desired accuracy. See Chapter 12 for more information.

If a polynomial expression is desired, Lundin [2] presents a formula accurate to 1 ppm. For $2.0 a \leq b$ (long coil):

$$L = \frac{\mu_0 n^2 \pi a^2}{b} \left[f_1 \left(\frac{4.0 a^2}{b^2} \right) - \frac{4.0}{3.0 \pi} \frac{2.0 a}{b} \right] \quad (6.3.1.4)$$

and for $2.0 a > b$ (short coil):

$$L = \mu_0 n^2 a \left\{ \left[\ln \left(\frac{8.0 a}{b} \right) - 0.5 \right] f_1 \left(\frac{b^2}{4.0 a^2} \right) + f_2 \left(\frac{b^2}{4.0 a^2} \right) \right\} \quad (6.3.1.5)$$

where for $0 \leq x \leq 1.0$

$$f_1(x) = \frac{1.0 + 0.383901 x + 0.017108 x^2}{1.0 + 0.258952 x} \quad (6.3.1.6)$$

$$f_2(x) = 0.093842 x + 0.002029 x^2 - 0.000801 x^3 \quad (6.3.1.7)$$

Grover [1] notes that the maximum inductance for a given length of wire occurs at

$$\frac{b}{2.0 a} = 0.408$$

The maximum is broad so the exact ratio is not critical.

For the square helical inductor the equations of Wheeler [8] are

$$L = \frac{4.0 \mu_0 n^2 a}{\pi} \left\{ \ln \left[1.0 + \frac{\pi a}{b} + \frac{1.0}{3.6443 + 2.0 (b/a) + 0.5098 (b/a)^2} \right] \right\} \quad (6.3.1.8)$$

This relation is accurate to within 0.001 L .

For coils wound on polygonal forms with an even number of sides, N , the NBS [7] recommends replacing a with a_0 where

$$a_0 = r \cos^2 \frac{\pi}{2.0 N} \quad (6.3.1.9)$$

The variable r is the radius of a circle that circumscribes the polygon.

6.3.1.1 Inductor Stray Capacitance

An empirical formula for the stray capacitance of this inductor is available in Medhurst

[4]. This equation is

$$C = 2.0 a H \quad (6.3.1.1.1)$$

where H is tabulated in [4]. A CAE equation has been fit to this data:

$$H = 0.299267 + 0.1012 \left(\frac{b}{2.0 a} \right) + \frac{0.0705822}{\left(\frac{b}{2.0 a} \right)} + \frac{3.568348e-6}{\left(\frac{b}{2.0 a} \right)^2} \quad (6.3.1.1.2)$$

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6.4 PLANAR OR PANCAKE INDUCTORS

6.4.1 Planar Circular Spiral Inductors

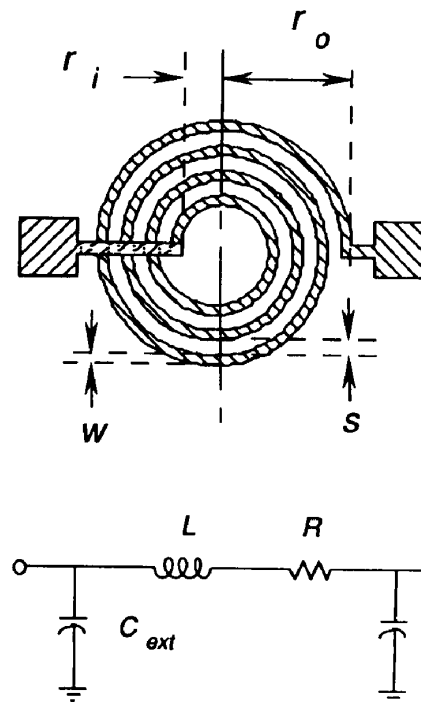


Figure 6.4.1.1: Planar Circular Spiral and Model

The spiral inductor is a technique of forming a planar inductor in a small space. The shape is described by an increasing radius with angle. The most common types of spirals have the radius functions given below¹ (the spiral of Figure 6.4.1.1 is an Archimedes spiral). These are the equations used to algorithmically generate the spiral inductor patterns.

$$r = r_i + k \theta \quad \text{Archimedes spiral} \quad (6.4.1.1)$$

¹Olivei [8] determined that the hyperbolic spiral has the lowest skin effect loss and the highest inductance per turn.

$$r = r_i e^{h\theta} \quad \text{Logarithmic spiral} \quad (6.4.1.2)$$

$$r = \frac{r_i}{1.0 - k r_i \theta} \quad \text{Hyperbolic spiral} \quad (6.4.1.3)$$

The equivalent circuit of the spiral inductor is also shown in Figure 6.4.1.1. The resistance is calculated with the skin effect resistance equations of the equivalent rectangular conductor.

$$R = \frac{K' \pi n a R_s}{w} \quad (\Omega) \quad (6.4.1.4)$$

where

$$R_s = \text{sheet resistance of the trace} \quad (\Omega / \text{square})$$

$$K' = 1.5, \text{ empirically}$$

Finite resistance of the coil will limit the Q . Circular spirals have higher Q 's (up to 100) than rectangular spirals; however, they also have lower inductance for an equivalent area. It has been found that $r_i / r_o = 1/5$ optimizes the Q .

The stray capacitance is distributed between turns and to any ground plane beneath the inductor. Burkett indicates that for spiral inductors with a lower ground plane the capacitance is 2 to 4 times the calculated parallel-plate capacitance due to fringing effects. This is probably because the fringing capacitance from the large outermost turn, C_{ext} , dominates. Wolff and Kapusta [16] successfully modeled the capacitance of a spiral inductor with a ground plane using coupled microstripline equations. This neglects the effects of the curved lines and coupling of nonadjacent turns. See Sections 5.5.4 and 6.6.2 for information on these effects.

The inductance is made up of three parts:

$$L = L_{\text{external}} + L_{\text{internal}} + L_{\text{mutual}} - L_{\text{image}} \quad (6.4.1.5)$$

The internal inductance is the self-inductance of the spiral created by magnetic flux internal to the conductor and is negligible. External inductance is the self-inductance created by flux external to the conductor. The self-inductance is increased by mutual inductance to adjacent turns (see Section 6.6.2). If a ground plane is present beneath the coil, the total inductance is reduced by an image of the coil that appears in the ground plane.

A number of relations for the inductance are found in the literature. Remke and Burdick [12] compare these formulas and their ranges of applicability. It is found that many of these relations are not valid for the range of parameters in which they are commonly used. A new formulation is presented

$$L_T = 2.0 \sum_{k=1.0}^{n-1.0} \sum_{j=k+1.0}^n \mu \sqrt{a b} \left[\left(\frac{2.0}{k_1} - 1.0 \right) K(k_1) - \frac{2.0}{k_1} E(k_1) \right] \\ + \sum_{k=1.0}^n \mu (2.0 c - w_1) \left[\left(1.0 - \frac{k_2^2}{2.0} \right) K(k_2) - E(k_2) \right] \quad (6.4.1.6)$$

where

$$k_1 = \sqrt{\frac{4.0 a b}{(a + b)^2}} \quad (6.4.1.7)$$

$$k_2 = \sqrt{\frac{4.0 c (c - w_1)}{(2.0 c - w_1)^2}} \quad (6.4.1.8)$$

$$a = \text{mean radius} = \frac{r_o + r_i}{2.0} \quad (6.4.1.9)$$

$$b = \text{thickness} = t$$

$$c = r_o - r_i \quad (6.4.1.10)$$

$$n = \text{number of turns}$$

$$w_1 = w / 2.0 \quad (6.4.1.11)$$

The functions $K(k)$ and $E(k)$ are the complete elliptic integrals of the first and second kind. See Chapter 12 for more information. This inductance equation is valid for $f \ll f_{\text{resonance}}$. As the resonance is approached the effect of the capacitances in the model need to be considered.

For the case of a spiral inductor having a lower ground plane the equation is

$$L_T = 2.0 \sum_{k=1.0}^{n-1.0} \sum_{j=k+1.0}^n \mu \sqrt{a b} \left[\left(\frac{2.0}{k_1} - 1.0 \right) K(k_1) - \frac{2.0}{k_1} E(k_1) \right]$$

$$\begin{aligned}
& + \sum_{k=1.0}^n \mu (2.0 c - w_1) \left[\left(1.0 - \frac{k_2^2}{2.0}\right) K(k_2) - E(k_2) \right] \\
& - \sum_{i=1.0}^n \sum_{j=1.0}^n \mu \sqrt{a_R b_R} \left[\left(\frac{2.0}{k_3} - 1.0\right) K(k_3) - \frac{2.0}{k_3} E(k_3) \right]
\end{aligned}
\tag{6.4.1.12}$$

where the variables are as defined above and

$$a_R = r_i + (i - 0.5)(w + s) \tag{6.4.1.13}$$

$$b_R = r_i + (j - 0.5)(w + s) \tag{6.4.1.14}$$

$$k_3 = \sqrt{\frac{4.0 a_R b_R}{(2.0 h)^2 + (a_R + b_R)^2}} \tag{6.4.1.15}$$

h = substrate thickness

Other Inductance Relations

Other equations found in the literature are listed here for reference as they are still commonly encountered and it may be desirable to adapt them for proprietary models. Before working with these, consult [12, 16].

Terman [14] gives the following relation, *Stefan's formula*, for the inductance of a flat wire or pancake spiral inductor:

$$\begin{aligned}
L &= 0.03193 a n^2 \left[\left(1.0 + \frac{h^2}{32.0 a^2} + \frac{c^2}{96.0 a^2}\right) \ln \left(\frac{8.0 a}{d}\right) - y_1 + \frac{y_3 c^2}{16.0 a^2} \right] \\
&+ 0.03193 a n \left(\ln \frac{D}{d} + 0.155 \right) \quad (\mu\text{H})
\end{aligned}
\tag{6.4.1.16}$$

where the dimensions (in inches) are

$$D = w + s \tag{6.4.1.17}$$

$$d = t + w \tag{6.4.1.18}$$

and y_1 and y_3 are in Table 16 of [14]. For a thin trace relative to c (the usual case for planar circuits), $y_1 = 0.5$ and $y_3 = 0.597$. For other cases the following equations were fit for CAE use:

$$y_1 = 3.74219 - 2.21817 \left(\frac{b}{c}\right) + 0.51724 \left(\frac{b}{c}\right)^2 - 3.24182 e^{-b/c} \tag{6.4.1.19}$$

$$\begin{aligned}
y_3 &= 0.59631105 + 0.038756584 \left(\frac{b}{c}\right) + 0.31720149 \left(\frac{b}{c}\right)^2 - 0.17592173 \left(\frac{b}{c}\right)^3 \\
&+ 0.039713327 \left(\frac{b}{c}\right)^4
\end{aligned}
\tag{6.4.1.20}$$

The accuracy of the curve fits (y_1, y_3) is to within $\pm 0.08\%$ and $\pm 1.5\%$, respectively. The accuracy of the inductance equation at dc is stated to be within 1%; however, this is not borne out in practice with planar structures.

A commonly quoted, simpler relation [4] due to Wheeler is:

$$L = \frac{n^2 a^2}{8.0 a + 11.0 c} \quad (\text{nH}) \tag{6.4.1.21}$$

where all of the dimensions are in mils. An improved equation that is more appropriate for use with planar spiral inductors was derived by Burkett [1]

$$L = \frac{0.8 n^2 a^2}{6.0 a + 10.0 c} \quad (\text{nH}) \tag{6.4.1.22}$$

Burkett's inductors were fabricated with a ground plane.

The connection from the center of the inductor to the outside can be made with a via and a trace on the opposite side of the substrate, an air bridge, or a bond wire. These may need to be included in the model as well.

Spiral inductors that have been ion beam milled having 0.5 mil lines and spaces can be purchased as SMD components. Olivei [8] overlaid the inductor patterns with a ferrite film to get large value inductances.

Easier to lay out versions of the circular spiral inductor are shown below [5]. This layout allows circular sections to be used rather than true spirals.

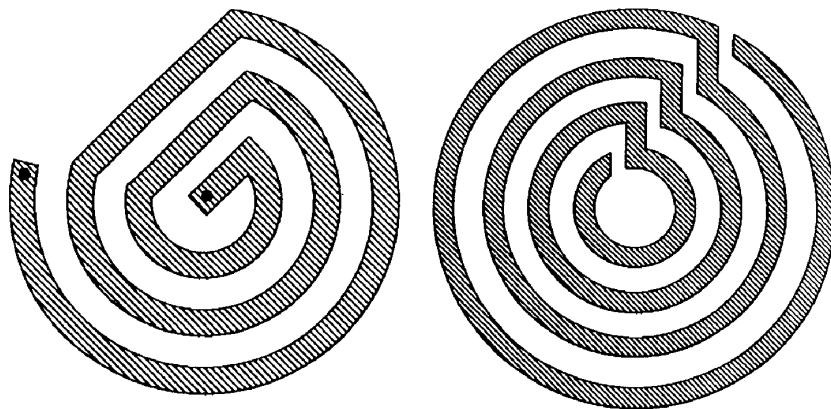


Figure 6.4.1.2: Alternate Circular Spiral Inductors

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6.4.2 Planar Square Spiral Inductor

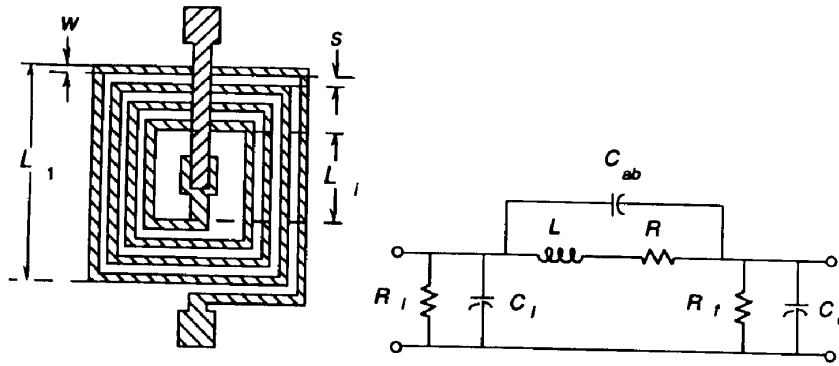


Figure 6.4.2.1: Planar Square Spiral

The square spiral inductor is an easier to lay out version of the circular spiral inductor. Again, connection from the inner turn to the outside is made with a via and a trace on the other side, an air bridge, or a bond wire. Modeling is somewhat complicated over the circular spiral by the addition of the bend discontinuities and coupling of the corners.

Shepherd [12] models this component with a cascade of coupled microstrip lines and abrupt bend discontinuities. The air bridge was modeled as parallel-plate capacitors coupling into the cascade of lines. Wolff and Kibuuka [14] show that the effect of the coupled corners is important, reducing the total inductance and modifying the frequency dependency. Design curves of the coupled corner discontinuity are available in Hill and Tripathi [7]. See Chapter 5 for CAE equations. The effect of a back-side ground plane (shunt capacitance) has been reduced by making the entire inductor a series of air bridges. The air under the bridges reduces the shunt capacitance.

It might be suggested that a concerted effort has not yet been made to simulate this structure rigorously and derive CAE equations that agree well with experimental data. Rather, most authors have taken an approach of adapting existing equations for wire inductors. Experimental data are sparse and agreement under one set of conditions has not guaranteed applicability for others. In fairness, this structure is complex—it has significant features in all three dimensions, can be used with several different techniques of connecting to the center, a ground plane may or may not be present and a multiple-turn inductor is not an extension of a single-turn inductor. The following categories can be identified as important to a model's broadband accuracy:

- Technique of connecting to center: via and backside trace, bond wire, air or dielectric bridge
- Dimensions w , s , substrate ϵ_r , overlay ϵ_r
- Number of turns

- Backside ground plane (and h) or no ground plane
- Presence of magnetic overlay or substrate
- Length of sections (from L_1) and frequency of operation
- Metallization losses and dielectric losses ($\tan \delta$)

If simulation below resonance is all that is desired, available static inductance equations may suffice; however, if the inductor is to be used as a resonator or the performance needs to be predicted beyond resonance some experimental work should be expected. A number of the available EM solving programs have shown good results with modeling this structure. With these caveats we will proceed to present the best available information for the designer.

In the model of Figure 6.4.2.1, C_{ab} represents the air bridge capacitance. Jansen *et al.* [9] found that this was calculable from the parallel-plate capacitance multiplied by a factor of 1.7 independent of turns. The resistor R is calculable from the skin effect resistance of the rectangular conductor; R_i and R_o represent dielectric and radiation losses.

In Greenhouse [5], an accurate technique is described after Grover for calculating the inductance of the square spiral inductor. This method calculates the inductance of the individual sections and includes both the positive and negative mutual inductances between parallel sections. Accuracies ranging from 6% to 20% of the measured value were obtained. Krafcsik and Dawson [9] added frequency dependency of the mutual inductances to the analysis of Greenhouse by taking into account the phase difference between the coupling inductors. As is the case with the circular spiral, the total inductance is

$$L = L_{\text{external}} + L_{\text{internal}} + L_{\text{mutual}} - L_{\text{image}} \quad (6.4.2.1)$$

The internal inductance is the self-inductance of the spiral created by magnetic flux internal to the conductor and is negligible. External inductance is the self-inductance created by flux external to the conductor. The self-inductance is increased by mutual inductance to adjacent turns. If a ground plane is present beneath the coil, the total inductance is reduced by an image of the coil that appears in the ground plane.

Greenhouse breaks the inductor into straight segments and sums the terms above for each of the segments. The internal inductance, L_{internal} , is calculated with (6.2.2.2). Mutual inductance is calculated for all possible pairs of parallel conductors with

$$L_{\text{mutual}} = 2.0 I_n \ln \left(\frac{l_n}{GMD_n} + \sqrt{1.0 + \frac{l_n^2}{GMD_n^2}} - \sqrt{1.0 + \frac{GMD_n^2}{l_n^2}} + \frac{GMD_n}{l_n} \right) \quad (6.4.2.2)$$

where the length of the n th set of parallel conductors is l_n and the geometric mean distance² between them in terms of centerline distance d is the variable GMD_n :

²Grover [6], p. 20.

$$GMD_n = \exp \left\{ \ln d - \left[\frac{1.0}{12.0 \left(\frac{d}{w} \right)^2} + \frac{1.0}{60.0 \left(\frac{d}{w} \right)^4} + \frac{1.0}{168.0 \left(\frac{d}{w} \right)^6} + \frac{1.0}{360.0 \left(\frac{d}{w} \right)^8} + \frac{1.0}{660.0 \left(\frac{d}{w} \right)^{10}} + \dots \right] \right\} \quad (6.4.2.3)$$

Mutual inductances add when the current flow is in the same direction and subtract when they are opposite. The mutual inductance to the image inductor appearing in any ground plane is

$$L_{\text{image}} = 2.0 l \left[\ln \left(\frac{L}{D} + \sqrt{1.0 + \frac{L^2}{D^2}} \right) - \sqrt{1.0 + \frac{D^2}{L^2}} + \frac{D}{L} \right] \quad (\text{nH}) \quad (6.4.2.4)$$

The mutual inductances can be modified to account for the difference in phase of the current in the segments following the suggestion of Krafcsik and Dawson

$$L_{\text{mutualKD}} = L_{\text{mutual}} \cos \theta \quad (6.4.2.5)$$

where

θ = phase between currents in the two segments

This can be calculated using linear distance (as if the coil were straightened out) between the two segments and the value of ϵ_{eff} calculated for microstrip (ground plane present) or coplanar strips (no ground plane). These equations are given in Chapter 3.

A computer program is necessary to evaluate these relations for any but the simplest of spiral inductors. Once this is accomplished, evaluation and design is quick and straightforward.

Other Square Spiral Inductor Equations

Liao [11] gives a simpler equation for the inductance

$$L = 8.5 \sqrt{A} n^{5/3} \quad (\text{nH}) \quad (6.4.2.6)$$

where

A = area occupied by inductor in cm^2

n = number of turns

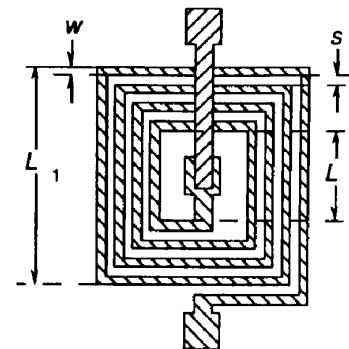


Figure 6.4.2.2: Rectangular Planar Spiral Inductors

NBS Circular C74 [13] gives a relation for the flat rectangular coil shown in Figure 6.4.2.2:

$$\begin{aligned} L_0 = & 0.009210 n^2 (a + a_1) \left[\log \left(\frac{2.0 a a_1}{w + n s} \right) - \frac{a}{a + a_1} \log (a + g) \right. \\ & \left. - \frac{a_1}{a + a_1} \log (a_1 + g) \right] + 0.004 (a + a_1) n^2 \left[2.0 \frac{g}{a + a_1} - \frac{1}{2} \right. \\ & \left. + 0.447 \frac{w + n s}{a + a_1} \right] - 0.004 n (a + a_1) (A_1 + B_1) \end{aligned} \quad (6.4.2.7)$$

The variables a_0 and a_0' are the dimensions of the outside of the n turn coil measured between center lines. The other variables are

$$g = \sqrt{a^2 + a_1^2} \quad (6.4.2.8)$$

$$a = a_0 - (n - 1.0) s \quad (6.4.2.9)$$

$$a_1 = a_0' - (n - 1.0) s \quad (6.4.2.10)$$

$$A_1 = \ln \frac{v + 1.0}{v + \tau} \quad (6.4.2.11)$$

$$B_1 = -2.0 \left[\frac{\delta_{12} (n - 1.0)}{n} + \frac{\delta_{13} (n - 2.0)}{n} + \frac{\delta_{14} (n - 3.0)}{n} + \dots + \frac{\delta_{1n}}{n} \right] \quad (6.4.2.12)$$

δ_{1n} = functions of $v = \frac{w}{s}$ and $\tau = \frac{t}{s}$ tabulated in Table 15 of [14]

(6.4.2.13)

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6.4.3 Single-Turn Circular Planar Inductor

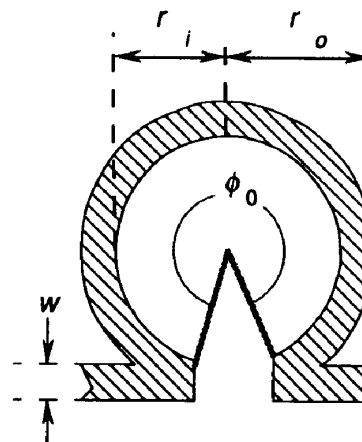


Figure 6.4.3.1: Single-Turn Planar Inductor

Pettenpaul [2] calculates the inductance of this structure with

$$L = \frac{\mu_0}{2.0 \pi (w + t)^2} \int_0^{\phi_0 / \sqrt{2}} (\phi_0 \sqrt{2.0} - 2.0 \phi) \cos(\sqrt{2.0} \phi) F(\phi) d\phi \quad (6.4.3.1)$$

where

$$F(\phi) = G(r_o, r_o) - G(r_i, r_o) - G(r_o, r_i) + G(r_i, r_i) \quad (6.4.3.2)$$

$$G(r_1, r_2) = \frac{R^3}{3.0} + \frac{2.0 r_1 r_2}{3.0} R \cos(\sqrt{2.0} \phi) + \frac{2.0 r_1^3 \cos(\sqrt{2.0} \phi)}{3.0} \operatorname{asinh} \frac{r_2 - r_1 \cos(\sqrt{2.0} \phi)}{r_1 |\sin(\sqrt{2.0} \phi)|} \quad (6.4.3.3)$$

$$R = \sqrt{r_1 + r_2 - 2.0 r_1 r_2 \cos(\sqrt{2.0} \phi)} \quad (6.4.3.4)$$

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6.4.4 Single-Turn Rectangular Planar Inductor

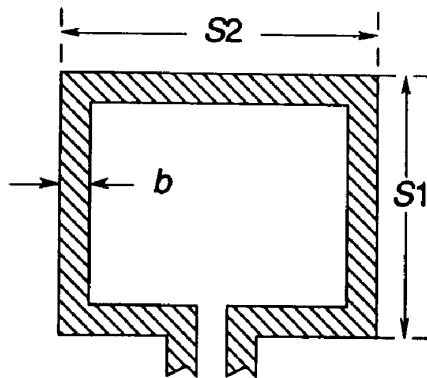


Figure 6.4.4.1: Single-Turn Rectangular Planar Inductor

$$L = 0.02339 \left[(s_1 + s_2) \log \left(\frac{2.0 s_1 s_2}{b + c} \right) - s_1 \log(s_1 + g) - s_2 \log(s_2 + g) \right] + 0.01016 \left[2.0 g - \frac{s_1 + s_2}{2.0} + 0.447(b + c) \right] \quad (\mu\text{H}) \quad (6.4.4.1)$$

where all dimensions are in inches, and

$b = w$ = the trace width

$c = t$ = the trace thickness

$$g = \sqrt{s_1^2 + s_2^2}$$

For a square, $s_1 = s_2 = s$ ($g = \sqrt{2} s$), the above simplifies to:

$$L = 0.04678 s \left\{ \log \left[\frac{2.0 s^2}{(s + \sqrt{2.0} s)(w + t)} \right] \right\} + 0.01016 [2\sqrt{2} s - s + 0.447(w + t)] \quad (\mu\text{H}) \quad (6.4.4.2)$$

where all dimensions are in inches.

REFERENCES

- [1] Djordjević, Antonije R., *et al.*, "Inductance of Perfectly Conducting Foils Including Spiral Inductors," *IEEE Transactions on Microwave Theory and Techniques*, Vol. MTT-38, No. 10, October 1990, pp. 1407-1414.
- [2] Terman, F.E., *Radio Engineer's Handbook*, McGraw-Hill, New York, 1945.

6.4.5 Octagonal Spiral Inductor

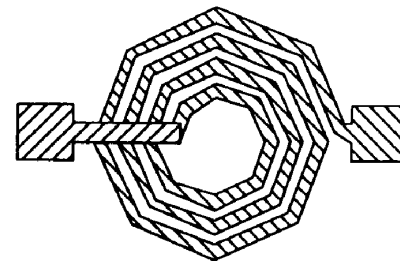


Figure 6.4.5.1: Octagonal Spiral Inductor

The octagonal spiral inductor is an alternative structure to the circular spiral and the square spiral inductor shapes. Although the spiral is an easy shape to describe mathematically, the CAD layout tool may not allow straightforward mathematical

programming of shapes. Most create metallization based on rectangular and circular primitives. This shape can be built up from rectangular shapes fairly easy.

Mondal [1], fit results of measurements to the inductance equations of Wheeler (1928) modifying the coefficients to the result

$$L = \frac{a^{1.54} n^{1.93}}{17.23 c^{0.55}} \quad (\text{nH}) \quad (6.4.5.1)$$

where a and c are in mils. The stated accuracy is 3.1% rms. Mondal also gives an interpolation equation for inductors with half turns.

$$L_{n+1/2} = \frac{(n+1/2)^2}{2} \left[\frac{L_n}{n^2} + \frac{L_{n+1}}{(n+1)^2} \right] \quad (6.4.5.2)$$

where

L_n = the inductance calculated with n turns

L_{n+1} = the inductance calculated with $n+1$ turns

$L_{n+1/2}$ = the inductance of an $n+1/2$ turn inductor

$n \in \mathbb{I}$ (integers)

The techniques of Greenhouse described in the above section on the square spiral inductor may also be adapted to this structure and other planar polygon inductors.

REFERENCES

- [1] Mondal, Jyoti P., "Octagonal Spiral Inductor Measurement and Modelling for MMIC Applications," *International Journal of Electronics*, Vol. 68, No. 1, 1990, pp. 113-125.
- [2] Stubbs, Malcolm, "Simulation Tool Accurately Models MMIC Passive Elements," *Microwaves & RF*, January 1988, pp. 75-79.

6.5 TOROIDAL INDUCTORS

Inductance of toroids is generally calculated from manufacturer's data with the relation:

$$L = \left(\frac{N}{100} \right)^2 \times (\text{inductance} / 100 \text{ turns}) \quad (6.5.1)$$

Where the (inductance / 100 turns) term is a parameter of the chosen coil form given by the manufacturer for a given material and toroid size and N is the number of turns. Curves of the Q vs. frequency are also given for many particular cases.

Toroidal and other inductors that use high-permeability materials to increase their inductance will introduce distortion. The distortion is generated by core material nonlinearities. If this is a problem the effect can be minimized by:

- keeping signal levels low
- removing dc before it reaches the inductors
- using air cores where possible

For a complete listing of inductor core shapes, sizes, materials, and their parameters see manufacturers' catalogs (Micrometals, Nytronics, and others). Some common materials' parameters are summarized in Table 6.5.1.

The wire gauge used in winding the inductor should be chosen as large as possible for lowest coil resistance and still fit on the form. The wire should also be able to sustain the required current flow without the wire or its insulation melting. Tables of wire fusing current for various gauges and insulation melting temperatures are available in *Reference Data for Radio Engineers*.

Mix	Material	Relative Permeability	Temperature (ppm/°C)	Frequency Range (MHz)	Color Code
0	Phenolic	1.0	0	50-250	Tan
1	Carbonyl C	20.0	280	0.15-2	Blue
2	Carbonyl E	10.0	95	0.25-10	Red
3	Carbonyl HP	35	370	0.02-1.0	Gray
3F	HP/Ferrite	80.0	700	0.01-1.0	Gray/Orange
6	Carbonyl SF	8.5	35	2.0-30	Yellow
7	Carbonyl TH	9.0	30	1.0-20	White
8	Carbonyl GQ4	35.0	255	0.02-1.0	Orange
10	Carbonyl W	6.0	150	10-100	Black
12	Synthetic Oxide	4.0	170	20-200	Green/White
15	Carbonyl GS6	25.0	190	0.10-3.0	Red/White
17	Carbonyl	4.0	50	20-200	Lavender
22	Synthetic Oxide	4.0	410	20-200	Green/Orange
26		75	822		Yellow/White
28		22	415		Gray/Green
33		33	635		Gray/Yellow
40		60	950		Green/Yellow
50	Ferrite	125	1000	0.01-1.0	Orange

Table 6.5.1: Some Common Core Materials [3]

6.5.1 Toroidal Inductor on Form with Square Cross-Section

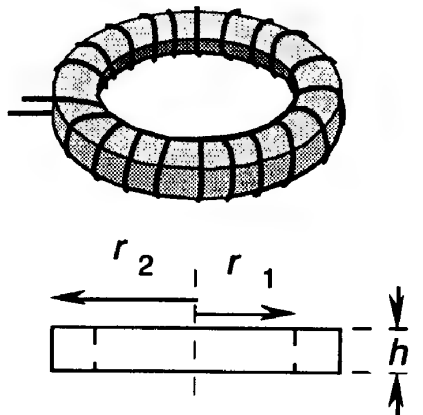


Figure 6.5.1.1: Toroidal Inductor with Square Cross-section

For a coil form whose μ_r is 1.0 [2]:

$$L = 0.004606 n^2 h \log_{10} \left(\frac{r_2}{r_1} \right) \quad (\mu\text{H}) \quad (6.5.1.1)$$

where

n = number of turns

6.5.2 Toroidal Inductor on Form with Circular Cross-Section

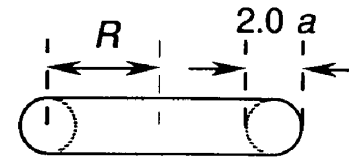


Figure 6.5.2.1: Toroidal Inductor with Circular Cross-Section

The toroid of circular cross section can be calculated with [3]

$$L = 0.01257 n^2 (R - \sqrt{R^2 - a^2}) \quad (6.5.2.1)$$

REFERENCES

- [1] DeMaw, M.F., "Magnetic Cores in RF Circuits," *RF Design*, April 1980, pp. 12-16.
- [2] Micrometals, Inc., *Inductor Catalog*.
- [3] U.S. Department of Commerce, National Bureau of Standards, Circular C74, *Radio Instruments and Measurements*, "Calculation of Inductance," January 1, 1937.

6.6 MUTUAL INDUCTANCES

In previous sections we have calculated the inductance of various conductors in isolation. When multiple conductors are involved, the inductance of each is affected by the presence of the others. In the case of two conductors carrying the same current, the inductance is either increased or decreased by an amount known as the **mutual inductance**. The direction of inductance change is determined by the relative directions of the current flows.

The total inductance of a two wire circuit carrying the same current is therefore

$$L = L_1 + L_2 \pm 2 M_{12} \quad (6.6.1)$$

6.6.1 Parallel Wires

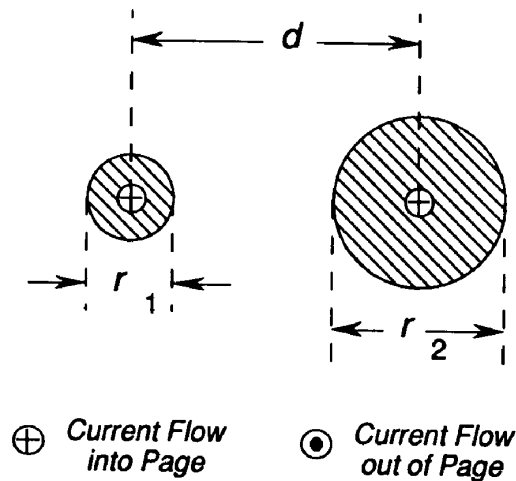


Figure 6.6.1.1: Mutual Inductance of Parallel Wires

For equal size ($r_1 = r_2$), parallel, wires with current flowing in the same direction the mutual inductance is

$$M_{12} = 0.002 l \left[\ln \left(\frac{l}{d} + \sqrt{1.0 + \frac{l^2}{d^2}} \right) - \sqrt{1.0 + \frac{d^2}{l^2}} + \frac{d}{l} \right] \quad (\mu\text{H}) \quad (6.6.1.1)$$

where d, l are in cm. The total inductance is

$$L_{\text{total}} = L_1 + L_2 + 2 M_{12} \quad (6.6.1.2)$$

where l is the length and L_1 and L_2 are calculated with (6.2.1.1).

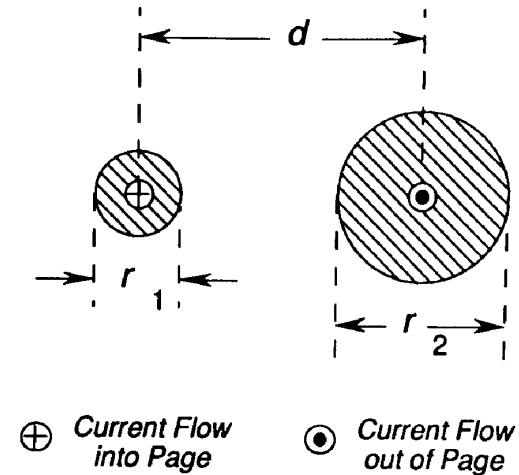


Figure 6.6.1.2: Return Circuit of Parallel Wires

For a return circuit of parallel wires the total inductance is

$$L_{\text{total}} = L_1 + L_2 - 2 M_{12} \quad (6.6.1.3)$$

where l is the length, L_1 and L_2 are calculated with (6.2.1.1), and M_{12} is calculated with (6.6.1.1). This can be simplified to

$$L_{\text{total}} = 0.002 l \left(\ln \frac{d^2}{r_1 r_2} + 0.5 \right) \quad (6.6.1.4)$$

REFERENCE

- [1] Grover, F.W., *Inductance Calculations*, Van Nostrand, Princeton, NJ, 1946; reprinted by Dover Publications, NY, 1962. (A concise book based on work at NBS.)

6.6.2 Parallel Rectangular Wires

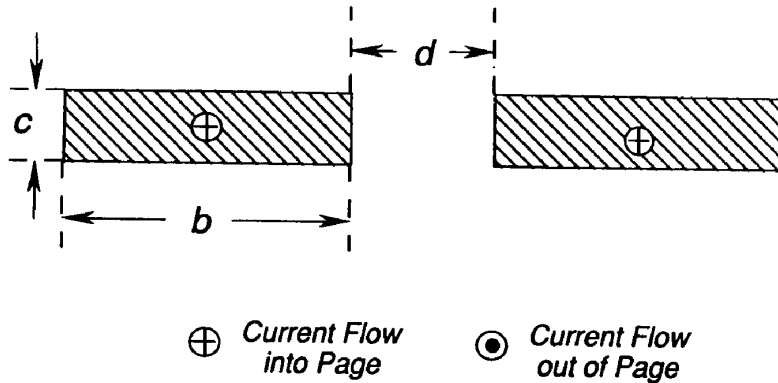


Figure 6.6.2.1: Parallel Rectangular Conductors

For parallel rectangles with current flowing in the same direction the mutual inductance is

$$M_{12} = 0.002 l \left(\ln \frac{2.0 l}{d} - \ln k - 1.0 + \frac{d}{l} - \frac{d^2}{4.0 l^2} \right) \quad (6.6.2.1)$$

For values of k , see Tables 1 and 2 of Grover [1]. Curve fits to this data have been derived for several values of b/c . For $b/c = 0.0$:

$$\ln k = 0.019361421 - 0.080643435 (c/d)^3 - 0.018564917 e^{c/d} \quad (6.6.2.2)$$

For $b/c = 0.1$:

$$\ln k = 0.020327157 - 0.077011572 (c/d)^3 - 0.019343819 e^{c/d} \quad (6.6.2.3)$$

For $b/c = 1.0$:

$$\ln k = 0.0010533337 - 0.0080261331 (c/d)^3 - 0.00094568688 e^{c/d} \quad (6.6.2.4)$$

These are accurate to approximately 4% over the range. The total inductance is

$$L_{total} = L_1 + L_2 + 2 M_{12} \quad (6.6.2.5)$$

where l is the length and L_1 and L_2 are calculated with (6.2.2.1).

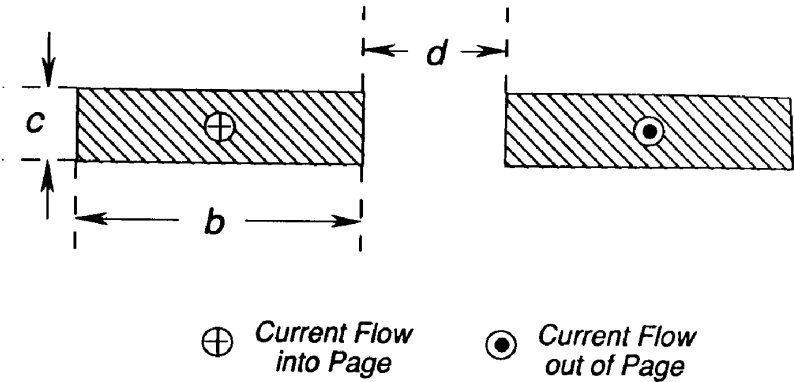


Figure 6.6.2.2: Return Circuit of Parallel Rectangular Inductors

The total inductance of the return circuit is

$$L_{total} = L_1 + L_2 - 2 M_{12} \quad (6.6.2.6)$$

where M_{12} is calculated with (6.6.2.1) and L_1 and L_2 are calculated with (6.2.2.1). The total inductance may also be written as

$$L_{total} = 0.004 l \left[\ln \frac{d}{b+c} + 1.5 + \ln k - \ln e \right] \quad (6.6.2.7)$$

where $\ln e$, and $\ln k$ are given in Grover [1] Tables 1, 2, and 3. Equations (6.6.2.2–4) are curve fits to $\ln k$ data. A curve fit to $\ln e$ is:

$$\ln e \cong \frac{7.3458933 \times 10^{-5} + 0.03476855 x - 0.01272244 x^2}{1.0 + 4.1576571 x + 22.249414 x^2 - 14.940776 x^3} \quad (6.6.2.8)$$

where

$$x = b/c \text{ or } c/b \quad (6.6.2.9)$$

This equation is an accurate fit to the data of Grover to within approximately 1%.

The meander line inductor can be analyzed with the equations of this section and Section 6.2.2. Meander lines are used as resistive elements, inductors, and delay lines. Note that adjacent conductors have equal and opposite current flow which reduces the total inductance over the equivalent length of straight conductor.

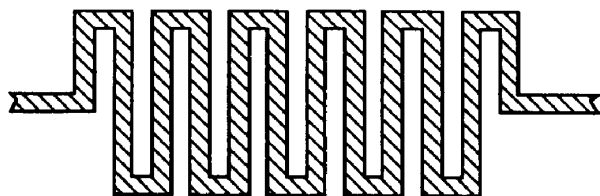


Figure 6.6.2.3: Meander Line

REFERENCE

- [1] Grover, F.W., *Inductance Calculations*, Van Nostrand, Princeton, NJ, 1946; reprinted by Dover Publications, NY, 1962. (A concise book based on work at NBS.)
- [2] Harokopus, William P., Jr., "Electromagnetic Coupling and Radiation Loss Considerations in Microstrip (M)MIC Design," *IEEE Transactions on Microwave Theory and Techniques*, Vol. MTT-39, No. 3, March 1991, pp. 413–421.
- [3] Sato, Risaburo, "A Design Method for Meander-Line Networks Using Equivalent Circuit Transformations," *IEEE Transactions on Microwave Theory and Techniques*, Vol. MTT-19, No. 5, May 1971, pp. 431–442.

Chapter 7

Capacitors

The capacitances of the structures in Chapters 3, 4, and 5 can be calculated using the equations of those chapters and (2.3.4). Chapter 2 describes capacitors made from short- and open-circuited transmission lines. The structures of this chapter are optimized for maximum capacitive reactance.

7.1 ABCD MATRICES

Series capacitor:

$$\begin{bmatrix} 1.0 & \frac{-j}{\omega C} \\ 0.0 & 1.0 \end{bmatrix}$$

Shunt capacitor:

$$\begin{bmatrix} 1.0 & 0.0 \\ j\omega C & 1.0 \end{bmatrix}$$

7.2 PARALLEL-PLATE CAPACITOR

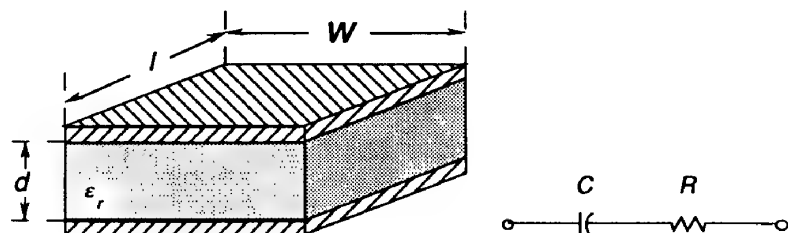


Figure 7.2.1: Parallel-Plate Capacitor

The equation for the parallel-plate capacitor is:

$$C = \frac{\epsilon_0 \epsilon_r l w}{d} \quad (F) \quad (7.2.1)$$

The equation assumes that there are no significant fringing fields or

$$d \ll l$$

$$d \ll w$$

$$\epsilon_r \gg \epsilon_0$$

The finite Q of the capacitor is expressed in terms of its Dissipation factor:

$$DF = \tan \delta + \omega C R_e \quad (7.2.2)$$

where

R_e = is the lead or electrode resistance

$\tan \delta$ = the loss tangent of the dielectric

For finite l and w and for unequal top and bottom plate areas, see [3] and [4].

REFERENCES

- [1] Cohn, Seymour, "Thickness Corrections for Capacitive Obstacles and Strip Conductors," *IRE Transactions on Microwave Theory and Techniques*, Vol. MTT-8, No. 11, November 1960, pp. 638–644.
- [2] Herbert, J.M., *Ceramic Dielectrics and Capacitors*, Gordon and Breach Science Publishers, New York, 1983.
- [3] Langton, N.H., "The Parallel-Plate Capacitor with Symmetrical Placed Unequal Plates," *Journal of Electrostatics*, Vol. 9, 1981, pp. 289–305.
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- [5] Mondal, Jyoti P., "An Experimental Verification of a Simple Distributed Model of MIM Capacitors for MMIC Applications," *IEEE Transactions on Microwave Theory and Techniques*, Vol. MTT-35, No. 4, April 1987, pp. 403–408.
- [6] Natarajan, Sundaram, "Measurement of Capacitances and Their Loss Factors," *IEEE Transactions on Instrumentation and Measurement*, Vol. MTT-38, No. 6, December 1989, pp. 1083–1087.

7.3 MULTILAYER DIELECTRIC PARALLEL-PLATE CAPACITOR

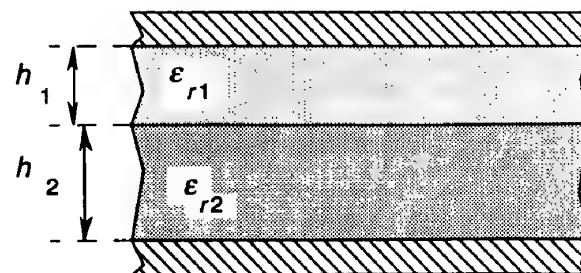


Figure 7.3.1: Parallel-Plate Capacitor with Layered Dielectric

The equations for a capacitor with layered dielectric are useful when calculating the effects of air or water intrusion into the dielectric. Use (7.2.1) and replace the dielectric constant, ϵ_r , with

$$\epsilon_{r,layered} = \left(\sum_{n=1}^{\text{all layers}} \frac{h_n}{h_t \epsilon_{r,n}} \right)^{-1} \quad (7.3.1)$$

where

$$h_t = \text{total thickness} = \sum_{j=1}^{\text{all layers}} h_j$$

$\epsilon_{r,n}$ = relative dielectric constant of the n th layer

7.4 THIN-FILM CAPACITORS



Figure 7.4.1: Thin-Film Capacitor

This structure may be calculated with the usual equations for a parallel-plate capacitor. Hurt [1] uses an equation that includes the effects of fringing capacitance:

$$C = \frac{\epsilon_r \epsilon_0 (l + \Delta_f) (w + \Delta_f)}{h} \quad (7.4.1)$$

where l and w are the dimensions of the top plate, which are effectively enlarged by an amount Δ_f due to fringing fields:

$$\Delta_f = \frac{4.0 h \ln(2.0)}{\pi} \quad (7.4.2)$$

REFERENCES

- [1] Hurt, John C., "A Computer-Aided Design System for Hybrid Circuits," *IEEE Transactions on Components, Hybrids, and Manufacturing Technology*, Vol. CHMT-3, No. 4, December 1980, pp. 525-535.
- [2] McLean, David A., "Dielectric Materials and Capacitor Miniaturization," *IEEE Transactions on Parts, Materials and Packaging*, Vol. PMP-3, No. 4, December 1967, pp. 163-169.
- [3] Pettenpaul, Ewald, *et al.*, "CAD Models of Lumped Elements on GaAs up to 18 GHz," *IEEE Transactions on Microwave Theory and Techniques*, Vol. MTT-36, No. 2, February 1988, pp. 294-304.

7.5 MULTILAYER CAPACITOR

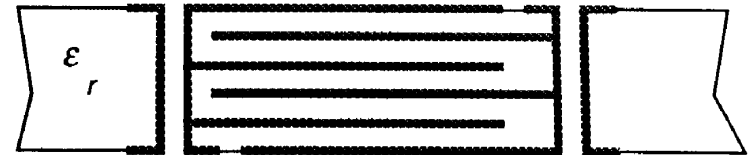


Figure 7.4.1: PCB Capacitor

The design equations for multilayer capacitors are an extension of the parallel-plate calculations above.

$$C = \frac{0.2249 \epsilon_r A (n - 1.0)}{d} \quad (\text{pF}) \quad (7.5.1)$$

where

A = area of plates in square inches

n = number of conductor layers

d = plate spacing

The equation is valid for d small.

REFERENCE

- [1] U.S. Department of Commerce, National Bureau of Standards, Circular C74, *Radio Instruments and Measurements*, "Calculation of Inductance," January 1, 1937.

7.6 INTERDIGITAL CAPACITOR

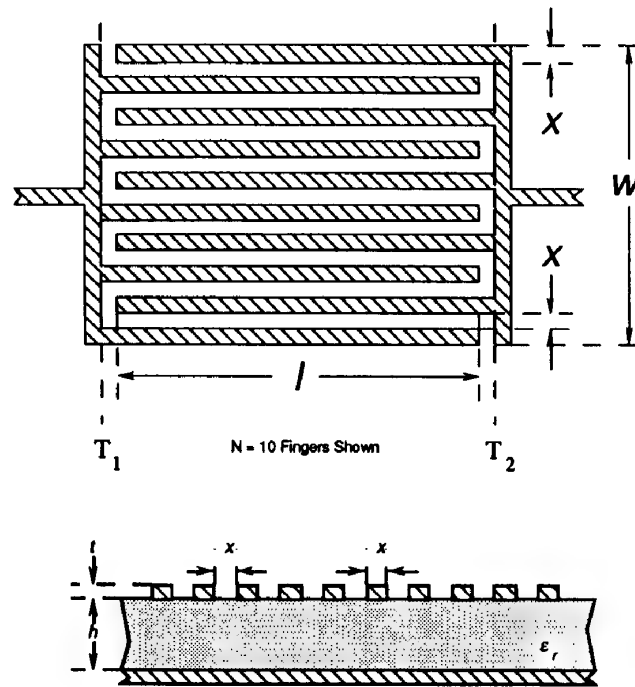


Figure 7.6.1: Interdigital Capacitor

The most often quoted relation for this structure is attributed to Alley [2]

$$C_2 = \frac{\epsilon_r + 1.0}{w} l [(N - 3.0)A_1 + A_2] \quad (\text{pF/in}) \quad (7.6.1)$$

$$R_{\text{series}} = \frac{4.0 l R_s}{3.0 w} \quad (7.6.2)$$

in pF/in of width, w , and N is the number of fingers. Equations are valid for

$$h > w/N$$

CAE equations have been fit to the curves of A_1 and A_2 in Alley:

$$A_1 = \left[0.3349057 - 0.15287116 \left(\frac{l}{x} \right) \right]^2 \quad (\text{pF/in}) \quad (7.6.3)$$

$$A_2 = \left[0.50133101 - 0.22820444 \left(\frac{l}{x} \right) \right]^2 \quad (\text{pF/in}) \quad (7.6.4)$$

These equations are valid to better than $\pm 0.35\%$ for

$$3.0 \leq l/x \leq \infty$$

The model element C_1 is calculated from the sum of the equivalent capacitance of the input (output) terminal and the shunt capacitance of the array of lines. The shunt capacitance can be calculated with the software presented in Smith [11]. The presence of fringing capacitance to the ground plane lowers inductance (due to the mutual inductance of the image); however, the Q is also decreased.

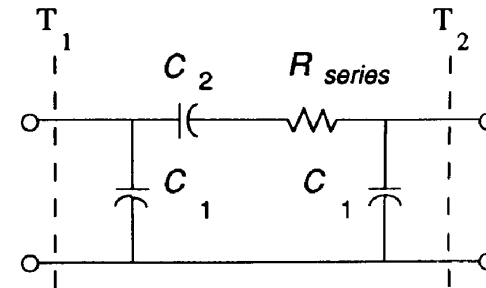


Figure 7.6.2: Interdigital Capacitor Model

Other models have been described [4, 12, 13] that are based on the parameters of arrays of lines and the discontinuities or numerical solutions. In particular, finger inductance, T junctions, open ends, and right-angle bends create significant effects.

Other interdigital capacitors are described in Crampagne *et al.* [3] (stripline interdigital capacitor) and Fouladian [5] (multiple generalized edge coupled striplines).

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- [1] Aitchison, Colin S., *et al.*, "Lumped-Circuit Elements at Microwave Frequencies," *IEEE Transactions on Microwave Theory and Techniques*, Vol. MTT-18, No. 12, December 1971, pp. 928-937.
- [2] Alley, Gary D., "Interdigital Capacitors and Their Application to Lumped-Element Microwave Integrated Circuits," *IEEE Transactions on Microwave Theory and Techniques*, Vol. MTT-18, No. 12, December 1970, pp. 1028-1033.
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- [4] Esfandiari, Reza, "Design of Interdigitated Capacitors and Their Application to Gallium Arsenide Monolithic Filters," *IEEE Transactions on Microwave Theory and Techniques*, Vol. MTT-31, No. 1, January 1983, pp. 57-64.
- [5] Fouladian, Jamshid, and Protap Pramanick, "Generalized Analytical Equations for Strip Transmission Lines," *Microwave Journal*, November 1989, pp. 165-168.
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- [11] Smith, J. I., "The Even- and Odd-Mode Capacitance Parameters for Coupled Lines in Suspended Substrate," *IEEE Transactions on Microwave Theory and Techniques*, Vol. MTT-19, No. 5, May 1971, pp. 424-430.
- [12] Wolff, Ingo, and Godfrey Kibuuka, "Computer Models for MMIC Capacitors and Inductors," *14th European Microwave Conference Proceedings*, 1984, Liège, Belgium, pp. 853-858.
- [13] Yanyang, Shi, and Jiang Zongquiang, "A Hybrid Approach to Analyse a Complex Microwave Planar Network," 1990, pp. 1509-1512.

7.7 CONCENTRIC CYLINDERS

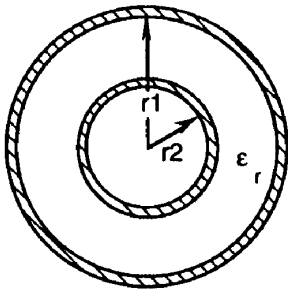


Figure 7.7.1: Capacitance of Concentric Cylinders

This equation can be used to calculate the capacitance of two concentric cylinders and to estimate the stray capacitance of spiral inductors.

$$C = 1.112 \epsilon_r \frac{r_1 r_2}{r_1 - r_2} \text{ (pF / cm)} \tag{7.7.1}$$

Results are in pF/cm of length. The stated accuracy is 0.1%.

REFERENCE

[1] U.S. Department of Commerce, National Bureau of Standards, Circular C74, *Radio Instruments and Measurements*, "Calculation of Inductance," January 1, 1937.

7.8 CAPACITOR DIELECTRIC CONSTANTS AND LOSS TANGENTS

Some of the dielectrics found in capacitors are tabulated in Table 7.8.1.

Dielectric	ϵ_r
Air	1.0006
Aluminum	8.0
Aluminum Oxide	7.0
Ceramic	35.0–6000
Glass	6.11
Kapton	3.2
Mica	6.8
Mylar	3.0
Polycarbonate	2.7
Polyethylene	3.3
Polypropylene	2.5
Polystyrene	2.5
Polysulfone	2.7
Quartz	3.8
Sapphire	9.4 (x, y) and 11.6 (z)
Tantalum	27.6
Tantalum Oxide	11.0
Teflon	2.1
Vacuum	1.00

Table 7.8.1: Capacitor Dielectric Constants

Chapter 8

Resistors

8.1 ABCD MATRICES

The ABCD matrix of a shunt resistor is:

$$\begin{bmatrix} 1.0 & 0.0 \\ \frac{1.0}{R} & 1.0 \end{bmatrix}$$

The ABCD matrix of a series resistor is:

$$\begin{bmatrix} 1.0 & R \\ 0.0 & 1.0 \end{bmatrix}$$

8.2 RECTANGULAR PLANAR RESISTOR

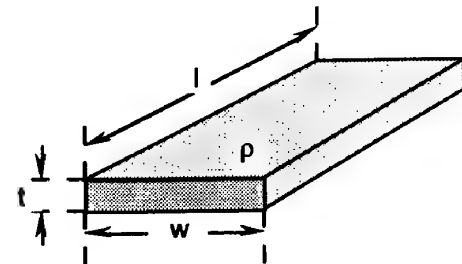


Figure 8.2.1: Rectangular Planar Resistors

This is the relation for resistance on which most equations are based. Chip resistors are formed by a rectangular film on a ceramic substrate with metal end caps. Trimming will introduce strays [1]. At high frequencies the skin effect effectively reduces the amount of the conductor used to carry the current. See Chapter 2 for more information.

$$R = \frac{\rho l}{w t} \quad (\Omega) \quad (8.2.1)$$

where the resistivity, ρ , is in Ω -units of l , w , and t . The inverse of resistivity is the conductivity σ . See Chapter 9 for a tabulation of common conductor resistivities.

For other cross-sectional geometries the resistance is calculated by replacing $w t$ in the above equation with the area. Calculation of the high frequency resistance can be calculated by multiplying the length of the surface by the skin depth and using this for the area. This approximation is not valid for surfaces with tight bends relative to the skin effect.

The above formula must be modified in many situations due to the effects of termination of the resistive film to the metallization at the ends and other process variables. Hybrid and resistor manufacturers will have empirical modifications of the (8.2.1). See [2] for an example of what has been done at Tektronix, Inc.

REFERENCES

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8.3 LEADED RESISTORS

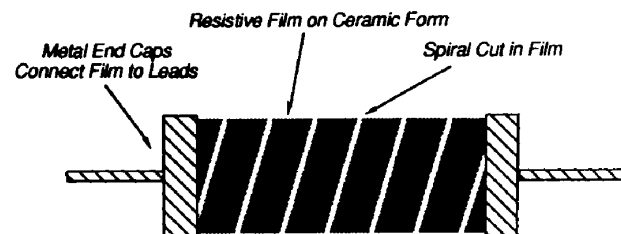


Figure 8.3.1: Metal Film Resistor

The metal film resistor is formed by coating a ceramic form with a resistive film. The film is then physically cut to form a helix of film. The ends are then capped to make connections to wire leads. The length, thickness, and resistivity of the film determine the total resistance.

This type of resistor is useful into the low hundreds of MHz range and possible further in less stringent applications. The capacitance of the end caps and the spiral gap are represented by a capacitor shunting the resistor. The leads and the spiral cut add an inductance. A family of resistors can be characterized by a resistance value below which the dominant strays are inductive and above which the dominant strays are capacitive.

For high frequency designs this type of component has two main problems. First, the lead inductance is a function of the PCB layout and the way the component is loaded in the board. The second problem is that there are different ways to build the resistor that give the same resistance but different strays. A different resistivity material can be substituted and the helix lengthened or shortened to correct for the change. This can wreak havoc with designs that depend on the strays being repeatable.

8.4 COAXIAL RESISTORS

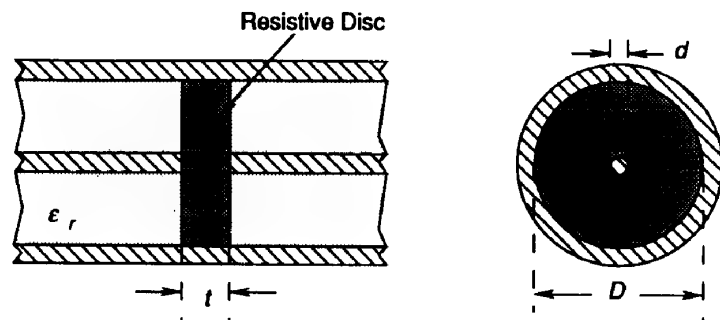


Figure 8.4.1: Shunt Coaxial Resistor

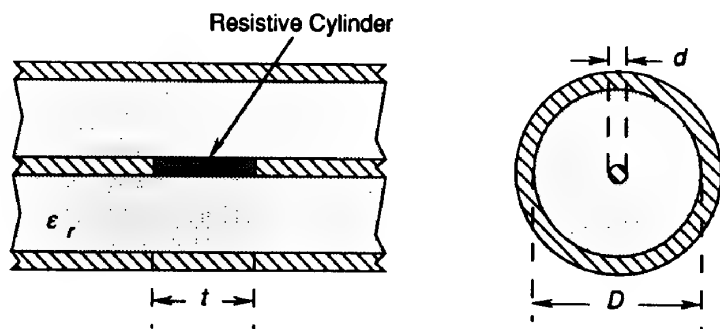


Figure 8.4.2: Coaxial Series Resistor

As described in Section 8.1 the resistance is calculated with (where d is the diameter of the resistive disk);

$$R = \frac{4.0 \rho t}{\pi d^2} \quad (\Omega) \quad (8.4.1)$$

If the diameter of the bend is small relative to the skin depth the high frequency resistance can be calculated with

$$R_{ac} = \frac{4.0 \rho t}{\pi \delta (2.0 d - 1.0)} \quad (\Omega) \quad (8.4.2)$$

The shunt and series coaxial resistors are distributed resistances used for high frequency attenuators and terminations. The shunt resistor is plated on a disk of dielectric. Resistive films are plated thin relative to the skin depth over the operating frequency range. This ensures a constant resistance vs. frequency. The lengths are kept electrically short to prevent discontinuity reactances.

By the appropriate combination of elements, a T or π pad attenuator can be built. A termination can be either a series or shunt element ending in a short.

8.5 MEANDER LINE RESISTORS

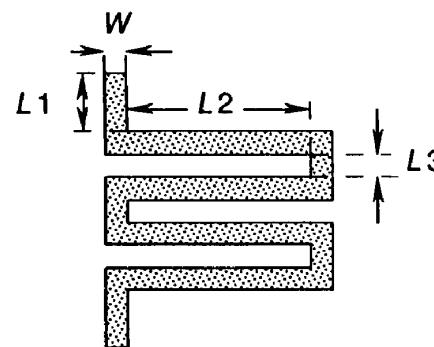


Figure 8.5.1: Meander Line Resistor

Meander line resistors are used to increase the amount of resistance that can be achieved in a given space. A similar structure can result from resistance trimming material or removing jumpers across resistive loops. This structure is not suitable for high frequency design because of the capacitive and inductive reactances of the bends and capacitances of the parallel coupled lines.

The resistance of meander lines can be calculated with a modification of the standard planar resistor equation. The resistance of the sections (L_1 , L_2 , L_3) is calculated as above and then the corner resistances are calculated [5]:

$$R_{corner} = 0.441 \frac{\rho l}{w t} \quad (\Omega) \quad (8.5.1)$$

The resistance of the corners is thus lower than an equivalent straight section. This effect can be eliminated by a slightly different construction which uses metal at the corners leaving only straight resistive sections (Figure 8.5.2 below).

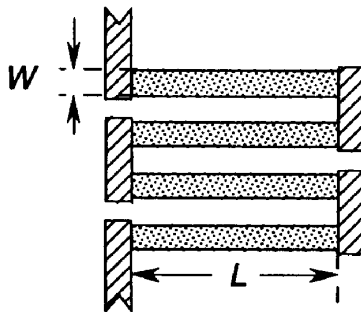


Figure 8.5.2: Shorted Meander Line Resistor

In this configuration metal foil shorts the corners, minimizing their effect on the resistance equations.

REFERENCES

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Chapter 9

Printed Circuit Fabrication

9.1 PCB CONSTRUCTION

Inherent in any PCB design are a number of electrical performance trade-offs. For boards having no controlled impedance lines, the choice of the materials and layering is usually left up to the PCB fabrication house. For controlled impedance design it is important to understand the trade-offs implicit in the layup. In this chapter the components used to build a printed-circuit board are described together with the parameters relevant to transmission line design. In particular, the effect of dimensional tolerances and material properties on electrical performance can be calculated with the information of this chapter using the relations of the previous chapters. In order to design a successful printed-circuit board transmission line it is also necessary to have a good understanding of the process. With this information, and an experienced printed-circuit board fabrication house, a design can be both electrically correct and cost effective.

The basic building blocks of a PCB are illustrated in Figure 9.1.1.

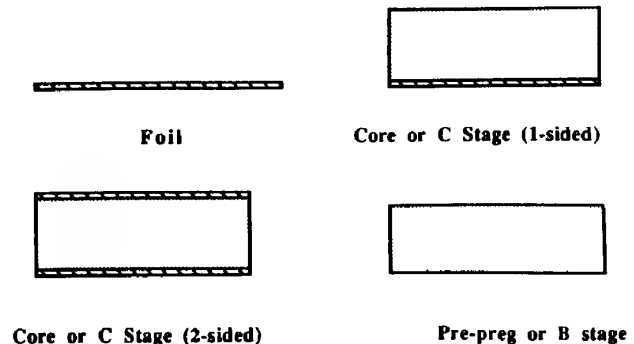


Figure 9.1.1: Components of a Printed-Circuit Board

Metal foil is available as separate sheets or preapplied to the dielectric material. Metal foil is applied under heat and pressure to the dielectric. The dielectric flows around *dendrites* that were formed during creation of the metal foil and create a strong physical bond. To create traces and component pads, a chemically resistant mask is applied to the metal. The undesired metal is then selectively removed by etching. Through-holes are plated by additive techniques which build up metal on the dielectric inside the holes, creating plated-through holes (also known as vias).

The dielectric layers can be solid pieces of material or a combination of dielectric and a binder such as epoxy.¹ The dielectric is often a woven fabric or, less commonly, chips or loose fibers. A *pre-preg* layer is a dielectric layer whose epoxy has not been fully cured. This allows it to reflow easily when adhering layers together. This flow of epoxy is important to bond the layers and to prevent shorts during plating of any plated-through holes.

When the patterns for each metal layer have been created, a sandwich is made up of the desired components. This is then pressed together under heat and vacuum to bond the layers together and create the final circuit board. This stack of metal and dielectric layers is known as a *layup*.

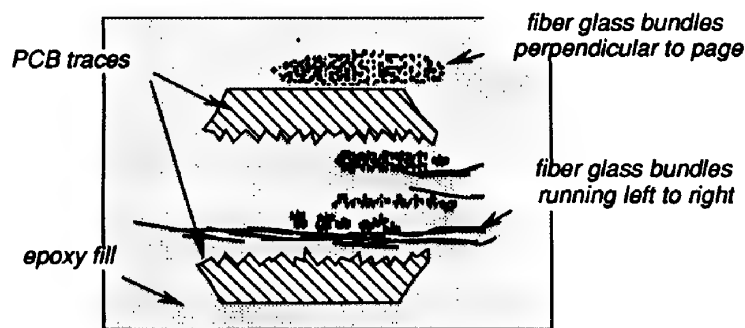


Figure 9.1.2: Cross-Section of PCB Inner Layer

In Figure 9.1.2, the cross-section of part of a laminated sandwich of materials is diagrammed. Part of a cross-section of an inner layer of an FR-4 printed-circuit board is shown (additional layers would continue vertically on the page). Notice the

¹For example, G-10 and FR-4 are dielectrics constructed from a fiberglass fabric and epoxy.

roughened surfaces of the copper traces that grip the fiber glass and epoxy. There are only a few layers of fiber glass cloth between the metallization layers, indicating that the layer is thin (on the order of ten mils). In this case, we see that a great deal of the space is filled with epoxy, rather than a homogenous mix of glass fibers and epoxy. Because it is the epoxy which actually flows to fill in surface irregularities in any layup, inner layer pc traces will generally be surrounded by epoxy.

9.1.1 Designing a Layup

Designing a layup begins with a number of physical constraints.

The thickness of the finished board is determined by system requirements (connectors, card guides, etc.).

The layers of the layup must be made symmetric about the center line. Different thicknesses of dielectric or metallization on either side of the pc board center line can cause board warpage. Board warpage is also minimized by equalizing the density of traces and balancing ground planes.

Inner layer dielectric thickness must be at least 1.5 times as large as the sum of the opposing metallization thicknesses. These inner layers must also have sufficient epoxy to fill the bumps created by the PC traces when the layers are compressed.

Finally, only a finite number of material thicknesses and dielectric constants are available off-the-shelf, so we are limited to selecting from these. The vendor may not have all of these in stock, so it is good practice to call the vendor who will actually build the board for a list of the materials on hand. The layup is then designed and line widths are calculated with these materials from this list. There is no point in carefully designing a layup with catalog materials only to discover that the vendor does not have what you wanted on hand.

For electrical purposes the starting point of any controlled impedance design should be C stage or core material. C stage has less variation because it is pre-cured and has copper already applied. The alternatives—built up B stage or pre-preg—introduce additional error sources and variations. These create a greater variation in impedance and propagation velocity.

In FR-4 up to 44 layers are possible at some vendors, 14 is a more typical maximum for production runs. Microwave laminations of 50 layers (Rogers RO2800) and 14 layers (RTI Duroid) have been reported in Markstein [1].

It is important to work with the PCB fabrication shop to develop a layup that meets electrical requirements and can be built.

9.2 SUBSTRATE DIELECTRICS

9.2.1 Properties of Commercial Dielectrics

The two parameters of a dielectric material which are critical in controlled impedance designs are ϵ_r and $\tan \delta$ (also called loss tan, or tan delta). These determine the line dimensions for a particular Z_0 and the dielectric losses, respectively. Table 9.2.1.1 summarizes these parameters for many of the popular commercial dielectrics.

Company, Dielectric	ϵ_r	$\tan \delta$	Description	Appearance
Generic				
G-10	4.3±0.05	0.008	Epoxy/ Glass	woven
Norplex/Oak				
FR-4 ED 130	4.3	0.020	Epoxy/ Glass	woven
FR-4 EM 145	4.3	0.020	Epoxy/ Glass	woven
FR-4 G 50	4.3	0.020	Epoxy/ Glass	woven
Cyanate Ester CE 245	3.5	0.005		
Polyimide G 30	4.2	0.015		
Paper/Phenolic NP 492	4.8	0.045		
Paper/Phenolic NP 930	4.7	0.025		
Mica				
FR-4 EG 150	4.5	0.020		
Polyimide PG 418	4.2	0.005		
Mitsubishi				
Plastics				
E002	5.10	0.022 (1 MHz) 0.020 (1 GHz)		
K002	3.6	0.0035 (1 MHz) 0.0035 (1 GHz)		
K012	3.5	0.0021 (1 MHz)		
Polyclad				
FR-4 PCL-FR-204	4.6	0.20		
Polyimide PCL-GI-702	4.6	0.010		
Hi-Tek				
Cyanate Ester AroCy M-40S	3.63	0.002		
Arlon				
DiClad 810	10.2	0.0027 (10 GHz)		
Cu Clad 217	2.17±0.04	0.0009 (1 MHz)	PTFE/Glass	woven, bendable

Cu Clad 233	2.33±0.02	0.0014		
Cu Clad 250	2.45, 2.50,	0.0008 (1 MHz)		
	2.55±0.4	0.0018 (10 GHz)		
Rogers Corp.				
Duroid 5500	2.5±0.04		Ceramic/ PTFE	smooth, uniform
Duroid 5870	2.33±0.02	0.0005 (1 MHz)	PTFE/Random	
		0.0012 (10 GHz)	Glass	
Duroid 5880	2.20±0.02	0.0004 (1 MHz)	PTFE/Random	
		0.0009 (10 GHz)	Glass	
Duroid 6002	2.94±0.04			
Duroid 6006	6.00±0.2	0.0025 (10 GHz)		
Duroid 6010.2	10.2±0.25		Ceramic/PTFE	
Duroid 6010.5	10.5±0.25	0.0028, max (10 GHz)		
RO2800	2.88±0.06			
Polyimide Fleximid	3.4	0.003		
TMM-3	3.24	0.0018 (10 GHz)		
TMM-4	4.5	0.0018 (10 GHz)		
TMM-6	6.5	0.0018 (10 GHz)		
TMM-10	9.8	0.0017 (10 GHz)		
TMM-13	12.85	0.0019 (10 GHz)		
Keene Corp.				
522-50	2.5±0.05	0.0010 (1 MHz)		
522-45	2.45±0.05	0.0010 (1 MHz)		
522-48	2.48±0.05	0.0010 (1 MHz)		
522-55	2.55±0.05	0.0010 (1 MHz)		
527-45	2.45±0.04	0.0019 (10 GHz)		
527-50	2.50±0.05	0.0019 (10 GHz)		
527-55	2.55±0.04	0.0019 (10 GHz)		
527-68	2.68±0.04	0.0019 (10 GHz)		
870-33	2.33±0.04	0.0012 (10 GHz)		
880-20	2.20±0.04	0.00085 (10 GHz)		
810-20	10.2±0.25	0.002 (10 GHz)		
810-50	10.5±0.25	0.002 (10 GHz)		
Epsilam 6	6.00	0.0018 (10 GHz)		
Epsilam 10	10.2±.25	0.002 (10 GHz)		

Crane Polyflon		
CuFlon	2.1	0.0001 0.00045 (1 GHz) 0.00045 (18 GHz)
Others		
Polysulfone	3.5-3.9 (1 MHz)	0.0056-0.009
	3.0 (10 GHz)	(1MHz) 0.005 (10 GHz)

Table 9.2.1.1: Properties of Dielectrics

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9.2.2 Effect of Resin/Glass Mix on ϵ_r of FR-4

In FR-4 and G-10 boards, a fabric of woven glass fibers is impregnated with epoxy and cured to form the substrate. Depending on the amount of glass used in a given PCB layer, the dielectric constant will be different. The effect of the amount of epoxy in the dielectric on the dielectric constant is shown in Figure 9.2.2.1 [1-3]. Table 9.2.2.1 lists the parameters for some commonly available materials. The overall trend

is the introduction of greater amounts of glass for thicker layers, which increases the dielectric constant.

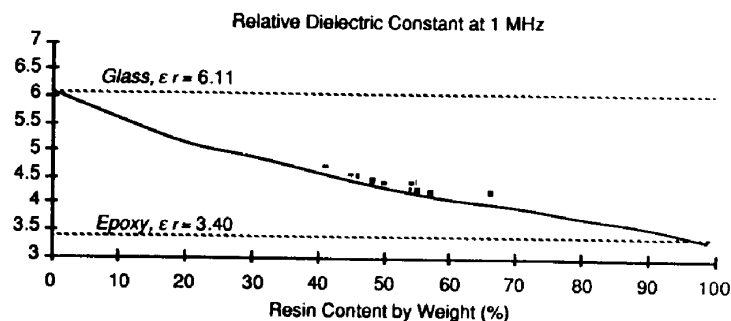


Figure 9.2.2.1: ϵ_r vs. Resin Content

The amount of resin and glass in a particular layer is dependent on a number of variables under the control of the vendor. A sheet of material will have 2 to 4% more epoxy if it is a B stage or pre-preg than if it is a C stage. Some of this extra epoxy will be pressed out during lamination. A variation of approximately $\pm 2\%$ in epoxy content can be expected, which produces a variation in ϵ_r of about ± 0.06 . This mix will also vary from vendor to vendor.

Material	Resin Content (%)	Nominal Thickness (mils)	ϵ_r
1080	57	2.50	4.22
2113	48	3.75	4.46
106/1080	60	4.42	4.20
2 ply 1080	57	4.90	4.22
2116	50	4.91	4.41
106/2113	55	5.42	4.27
1080/2113	55	6.50	4.27
106/2116	55	6.55	4.27
1080/2116	52	6.80	4.35
7628	41	7.13	4.68
2 plies 2113	48	7.30	4.46
106/7628	46	7.80	4.52
2113/2116	46	8.25	4.52
2 plies 2116	50	9.35	4.40
1080/2116/1080	54	9.90	4.30
2116/1080/2116	46	10.00	4.50
1080/7628	46	10.00	4.52
1080/2116/1080	54	10.10	4.30
106/7628/106	50	11.30	4.40
2116/7628	45	11.90	4.55
1080/7628/1080	48	12.20	4.46
2 plies 7628	41	14.30	4.68
3 plies 2116	50	14.30	4.40
2113/7628/2113	41	14.50	4.55
7628/106/7628	43	16.10	4.62
7628/1080/7628	44	16.20	4.58
2 plies 2113/2 plies 2116	46	16.20	4.58
2116/7628/2116	47	17.00	4.49
7628/2113/7628	43	17.50	4.62
2 plies 1080/2 plies 7628	46	19.50	4.52
2 plies 1080/2 plies 7628	46	19.50	4.52
2 plies 1080/2628	44	20.00	4.70
3 plies 7628	41	20.50	4.68
2 plies 2113/2 plies 7628	44	20.00	4.70
2 plies 2116/2 plies 7628	45	23.20	4.55
2 plies 1080/3 plies 7628	43	25.00	4.75
4 plies 7628	40	28.00	4.90

Table 9.2.2.1: FR-4 Characteristics

REFERENCES

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9.2.3 Ceramic Dielectrics

Ceramics	ϵ_r	$\tan \delta$
Alumina (Al_2O_3 85%)	8.3	0.0058
Alumina (Al_2O_3 96%)	10	0.0002
Alumina (Al_2O_3 99.5%)	9.70	0.0002 (1 MHz) 0.0003 (10 GHz)
Aluminum Nitride (AlN)	8.8	0.0005–0.001
Beramic Z (BeO 97.9%)	6.65 (1 MHz)	0.0006
Beryllia (BeO 97%)	6.90	0.0002 (1 MHz) 0.0003 (10 GHz)
Boron Nitride (BN)	4.0	0.0004
Fused Silica	3.78	0.0001 (1 MHz) 0.00025 (25 GHz)
Magnesium Titanate	16	0.0002
Quartz, Fused Silica	3.78	0.0004 (10 GHz)
Rutile (TiO_2 99.5%)	85	0.0002
Silicon Carbide (SiC)	40.0	0.5

Table 9.2.3.1: Ceramic Dielectric Properties

REFERENCES

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9.2.4 Flexible PCB Dielectrics

Printed-circuit boards can be made from thin sheets of dielectric and metallization which are flexible. This allows the board to be folded to conform to odd packaging

demands and to replace cabling. The table below lists the important dielectric properties of these materials. Flexible dielectrics are commonly available in thicknesses of 0.5, 1, 2, 3, and 5 mils.

Dielectric	ϵ_r	$\tan \delta$
Dacron/Epoxy	3.2	0.015
Kapton® (Polyimide)	3.5 (1 kHz)	0.0025 (1 kHz)
Mylar® (Polyester)	3.45 (1 kHz)	0.0028 (1 kHz)
	3.35 (1 MHz)	0.008 (100 kHz)
	3.32 (10 MHz)	0.0115 (10 MHz)
Nomex® (Aramid fibers)	1.6–3.7 (higher for thicker)	0.004–0.007 (higher for thicker)
	(60 Hz)	(60 Hz)

Table 9.2.4.1: Flexible PCB Dielectric Properties

REFERENCE

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9.2.5 Semiconductor Dielectrics

Dielectric	ϵ_r	$\tan \delta$	σ (S / m)
Gallium Arsenide (GaAs)	13.10	0.0016 (10 GHz)	8×10^{-3}
Germanium (Ge)	16.0		
InSb	15.9		
Sapphire (Al_2O_3)	10	0.0001 (10 GHz)	5.5×10^{-4}
Silicon (Si)	11.7	0.005	4.39×10^{-4}
SiO	6	0.007	
SiO ₂	3.9	0.007	
Ta ₂ O ₅	25	0.003	

Figure 9.2.5.1: Semiconductor Dielectric Properties

9.2.6 Properties of Other Dielectric Materials

Occasionally it is desired to build a transmission line structure into the mechanical structure of the system it is in. For example, a controlled impedance line might be molded into a case.

The plastics vendor literature provides a good start to finding the properties of the material. Dielectric constants and power factors are usually given at 60 Hz and some higher frequency (1 kHz or 1 MHz). Plastics with common trade names come in a variety of forms with slightly different electrical properties—try to determine the exact type you are using. The most variation will be found in materials that are loaded or filled with a second material to varying degrees in order to achieve different physical parameters such as rigidity.

It should not be assumed that the low frequency parameters will hold at microwave or even RF frequencies. High frequency data or sample testing will be necessary.

REFERENCE

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9.2.7 Dielectric Dimensions and Tolerances

To determine the effects of mechanical tolerances on the design we need to know the vendor specifications for these materials. Although these will vary depending on the material quality and the vendor, the numbers in this section can be used for estimation purposes. Together with the equations of the previous chapters the effect of variations on the electrical parameters Z_0 , ϵ_r , v_p , and electrical length can be calculated.

9.2.7.1 Finished PCB Dielectric Thicknesses

As discussed in Section 9.1, the printed-circuit board is built out of a number of layers sandwiched together. A large number thicknesses can be achieved from the available component materials; however, system physical requirements will usually dictate the exact value of the overall thickness. The following overall thicknesses are the most commonly used.

Backplanes	0.125 in $\pm 10\%$
Daughter Cards	0.03125 in, 0.0625 in, 0.093 in

Table 9.2.7.1: Finished PCB Dielectric Thicknesses

Standard microwave dielectric thicknesses (overall) vary, depending on material. Some typical values are:

0.010 ± 0.0007 in
0.031 ± 0.001 in
0.062 ± 0.002 in
0.093 ± 0.003 in
0.125 ± 0.004 in

Table 9.2.7.1.2: Microwave Dielectric Thicknesses

9.2.7.2 Finished B-Stage Dielectric Tolerances

Dielectric Thickness (in)	Tolerance (±in)
0.005 to 0.008	0.0015
0.009 to 0.016	0.002
0.017 to 0.030	0.003
0.031 to 0.060	0.005

Table 9.2.7.2.1: Dielectric Thickness Tolerances

9.2.7.3 C-Stage Dielectric Tolerances

Tolerances are a function of the dielectric thickness and are given in Table 9.2.7.3.1. Class 4 applies to microwave dielectrics; classes 2 and 3 are two grades of glass reinforced dielectrics (FR-4, G-10).

dielectric thickness (in)	Class 2 (± in)	Class 3 (± in)	Class 4 (± in)
0.0010 to 0.0045	0.00075	0.00075	
0.0045 to 0.0065	0.0010	0.00075	
0.0066 to 0.0120	0.0015	0.0010	
0.0121 to 0.0200	0.0020	0.0015	
0.0201 to 0.0299	0.0025	0.0020	
0.0300 to 0.040	0.0040	0.0030	0.0020
0.041 to 0.065	0.0050	0.0030	0.0020
0.066 to 0.100	0.0070	0.0040	0.0030
0.101 to 0.140	0.0090	0.0050	0.0035
0.141 to 0.250	0.0120	0.0060	0.0040

Table 9.2.7.3.1: Dielectric Thickness Tolerances

REFERENCES

- [1] Hornig, Carl F., "Multilayer Backplanes are Designed for Efficiency," *Electronic Packaging and Production*, June 1985, pp. 162-164.
- [2] Markstein, Howard W., "Multilayer Board Fabrication Sheds Previous Limitations," *Electronics Packaging and Production*, July 1989, pp. 24-28.

9.3 METALLIZATION

9.3.1 Standard Conductive Layers and Tolerances

C stage materials usually begin the lamination process with copper material already attached by the vendor. On some layers, sheets of copper are added. These conductive layers of the PCB are made from thin sheets of electrodeposited or rolled copper. The sheets are purchased in rolls or in flat sheets prepunched with alignment holes.

Metal backings can also be provided which are quite thick (e.g., 0.125"). This extra thickness is useful for improved grounding, mechanical rigidity, and thermal management.

9.3.2 Metal Thickness

Inner metallization layers are bare copper. Outer layers are usually solder-plated copper although they may be solder mask over bare copper. Metallization thickness is specified in ounces of copper per square foot.

Copper Weight	Inner Layer	Outer Layer (after plating)
1/2 oz	0.0007±50 µin	
1 oz	0.0014±50 µin	0.0028±190 µin
2 oz	0.0028±50 µin	0.0042±190 µin
3 oz	0.0042±50 µin	0.0056±190 µin

Table 9.3.1.1.1: Metallization Thicknesses

It is preferable to use a single weight of copper on inner layers (core layers) due to possible etching restrictions (some vertical etchers cannot etch each side at a different rate). Outer copper layers can be applied separately and do not have this restriction.

9.3.3 Trace Width Tolerances

Table 9.3.1.2 tabulates line width tolerances as a function of copperweight. Greater metal thickness requires a longer etch time and creates a corresponding increase in uncertainty of line widths and spacings. You should specify the thinnest metal that meets your other requirements (losses, power handling).

Another large influence on the tolerances of trace widths and spacings is the density of metal over the board's surface. It is best to keep the density constant over the board's surface. This can be done by filling empty areas of the board with dummy traces or cross hatching to keep ensure an even etching rate. These areas are known as *thieving*.

Copper Weight	I	II, III
0.5 oz	±1 mil	±0.75 mil
1 oz	2	1.5
2 oz	3	2
3 oz	4	3

Table 9.3.1.2: Trace Width and Spacing Tolerances

9.3.4 Minimum Trace Widths and Line Spacings

There are limitations to the thinnest line widths that can be achieved repeatably. Some vendors report that 2 to 4 mil lines and spaces are possible [1].

Copper Weight	Nominal Thickness (in)	Min. Spacing (in)	Min. Width (in)
1/4-1/2 oz.	0.00035	0.006	0.006
1/2-1 oz.	0.0007	0.009	0.010
1 oz.	0.0014	0.012	0.013
2 oz.	0.0028	0.015	0.010

Table 9.3.4.1: Minimum Trace Width and Spacings

REFERENCE

- [1] Markstein, Howard W., "Multilayer Board Fabrication Sheds Previous Limitations," *Electronics Packaging and Production*, July 1989, pp. 24-28.

9.3.5 Plating

Metal is usually plated over the etched circuit traces and into holes as summarized in Table 9.3.5.1. It is becoming more common to have solder mask over bare copper on outer layers.

This construction works particularly on SMD boards and boards with large areas of ground plane. In plated designs, the ground plane will collect solder during wave soldering or if solder masked the plating will reflow and crinkle the solder mask.

Plating is applied by electrodeposition or by an electroless (chemical) process.

Gold plating (≥99.7% Au) is typically applied as a greater than 30 µin layer over 100 µin of electroplated nickel.

Plating Type	Usage	Thickness
Copper	Good conductivity	1-2 mils
Gold	Wire bondable, corrosion resistant, low-loss	10-50 µin
Tin-Lead	Corrosion resistant, solderable	0.1-0.5 mils
Nickel	Barrier between Cu and Au	50-300 µin
Solder	Corrosion resistant, solderable	0.1-0.3 mil

Table 9.3.5.1: Plating Properties

9.3.6 Electrical Properties of Metals

The electrical properties of commonly used metals are given in Table 9.3.2. The relative resistivity numbers are provided for reference when using certain CAE programs. The resistivity of thin films of metal will be substantially higher than the values in the table.

Material	μ_r	ϵ_r	σ (S / m)	ρ / ρ_{Cu} (unitless)	ρ / ρ_{Al} (unitless)
Silver	0.99998	1	6.17×10^7	0.9335	0.6645
Copper	0.999991	1	5.76×10^7	1.0000	0.7118
Aluminum	1.00002	1	3.72×10^7	1.5484	1.1022
Sodium	1	1	2.1×10^7	2.7429	1.9524
Brass	1	1	2.56×10^7	2.2500	1.6016
Tin	1	1	0.87×10^7	6.6207	4.7126
Graphite	1	1	0.01×10^7	576.0	410.0
Gold	1	1	4.10×10^7	1.4049	1.000
Steel ¹	1,000		$0.14-0.77 \times 10^7$	7.481-41.143	5.324-29.286
Stainless Steel	1,000		0.11×10^7	52.364	37.273
Mumetal	100,000				
Solder	1		0.7×10^7	8.2286	5.8571
Nickel	600		1.45×10^7	3.9724	2.8276

Table 9.3.2: Electrical Properties of Common Metals

REFERENCE

- [1] *Reference Data for Radio Engineers*, Howard W. Sams & Co., Inc., 1956.

¹Steel comes in a variety of alloys with varying permeabilities. This is a representative value.

9.3.7 Metal Roughness

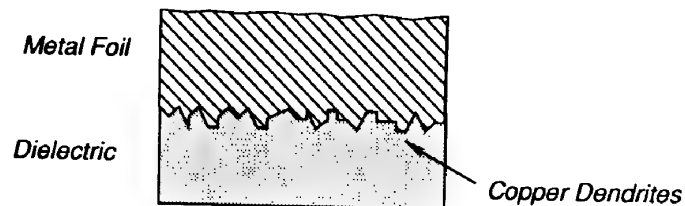


Figure 9.3.3.1: Metal Roughness

As discussed earlier, the surface roughness of substrate metallization entwines with the dielectric to form a mechanical bond. This same roughness introduces additional losses at high frequencies by increasing the path length when the skin depth approaches the surface roughness.

There are two different techniques used to create the thin layer of copper used on printed-circuit boards. **Electrodeposited copper** is made by electrically attracting copper from a bath onto a rotating drum. This gives a rough surface, which is desirable for good adhesion to the dielectric. The conductor side nearest the drum is smoother than the electrochemically applied surface. **Rolled copper** is cast, milled smooth, and then cold rolled to the desired thickness. This gives a very smooth surface that has lower electrical losses.

Material	Surface Roughness	Resistivity
Electrodeposited Copper		$<1.5940 \Omega\text{-g/m}^2$
1/2 oz	75 $\mu\text{in, rms}$	
1 oz	95 $\mu\text{in, rms}$	
2 oz	115 $\mu\text{in, rms}$	
Rolled Copper		$1.52361 \Omega\text{-g/m}^2$
1/2 oz	55 $\mu\text{in, rms}$	
1 oz	55 $\mu\text{in, rms}$	
2 oz	55 $\mu\text{in, rms}$	

Table 9.3.3.1: Metal Foil Roughness

Although rolled copper has less surface roughness, the conductivity is not anisotropic whereas with electrodeposited copper it is. The conductivity of rolled copper is about 20% lower against the direction of the roll [3].

Thick metal layers laminated to the dielectric have surface roughnesses of approximately 70 $\mu\text{in.}$, rms.

REFERENCES

- [1] Bonfield, Robin, "Choose Metals for Optimum PTFE Performance," *Microwaves & RF*, March 1990, pp. 121–124.
- [2] Gurley, Steve, *Flexible Circuits: Design and Applications*, Marcel Dekker, New York, 1984.
- [3] Hiroyuki Tanaka, and Fumiaki Okada, "Precise Measurements of Dissipation Factor in Microwave Printed Circuit Boards," *IEEE Transactions on Instrumentation and Measurement*, Vol. IM-38, No. 2, April 1989, pp. 509–514.
- [4] Swansen, William K., "Good Planning Ensures Clean Microwave Boards," *Microwaves & RF*, December 1985, pp. 79–86.

9.4 PLATED-THROUGH HOLES OR VIAS

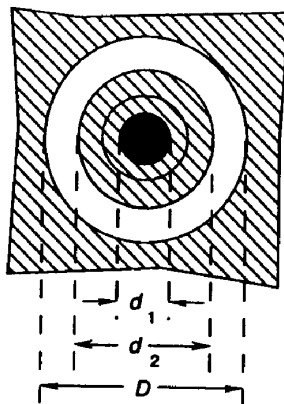


Figure 9.4.1: Plated-Through Hole Pad Dimensions

Plated-through holes can have diameters of approximately 0.032 to 0.189". A ring of metallization and a ring of clearance must be provided on each layer the hole passes through. The required dimensions are given in the table below for preferred plated-through hole sizes. Remember that the outer conductor can introduce a capacitive discontinuity to the pad if it is too close. Vias can pass completely through the board; be buried internally (*buried via*); or start on the outside and stop inside the board (*blind via*).

Plated-through holes should have a 5:1 or lesser ratio of board thickness to hole diameter (9:1 is possible) and should have diameters of at least 20 mil to avoid drilling and plating difficulties. Plated-through holes can be provided that connect to a thick metal backing layer. The plating thickness in the hole is approximately 1 mil.

Pins, rivets, and eyelets are possible alternatives to plated-through holes; however, they introduce an extra step in the assembly process.

Plated-Through Hole Diameter, d_1	Minimum Annular Ring Width, d_2	Metallization Clearance, D
(in)	(in)	(in)
0.036	0.080	0.150
0.058	0.100	0.170
0.082	0.130	0.200
0.036	0.060	0.100
0.058	0.080	0.130
0.082	0.110	0.160

Table 9.4.1: Plated-Through Hole Dimensions

Thermal relief vias (Figure 9.4.2) isolate a pad thermally from a solid ground or power plane. This allows the pad to be heated more easily for soldering and unsoldering component leads. This pattern introduces an inductive discontinuity at high frequencies.

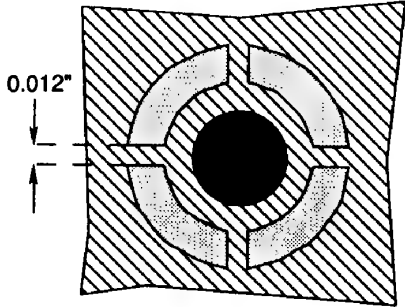


Figure 9.4.2: Thermal Relief Via

9.5 SOLDER MASK

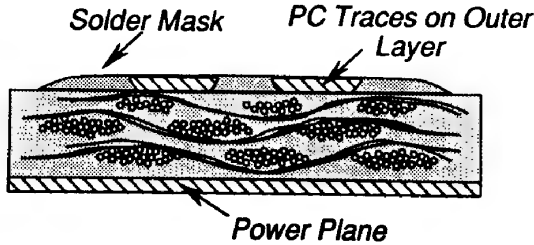


Figure 9.5.1: Solder Mask

Solder mask is applied over a finished printed-circuit board's traces to restrict the flow of solder to the desired areas during wave soldering. Solder mask also prevents the flow of small leakage currents between traces caused by impurities introduced during handling. Hard solder mask also prevents fine traces from being damaged during handling. Solder mask can be applied dry in sheets or wet by screening or spraying.

Electrically, our concern is that the addition of the solder mask reduces outer layer trace impedance. The exact amount of the reduction depends on the dielectric constants

of the solder mask and the substrate, and the w/h ratio. See Chapter 3 for information on calculating this effect.

Type	Description	ϵ_r	$\tan \delta$	thickness
Standard (PC401)	epoxy	3.3-4.0	0.03-0.05	1-2 mil
Dry Film		3.80		3-3-1/2 mil
Photoimaged	Screened	3.8		1 mil
Conformask™		4.0-4.2		2 mil

Table 9.5.1: Solder Mask Properties

The thickness of the solder mask is controlled by the screen mesh used in the screening process. Excessively thick masks can make it difficult to mount SMD components.

REFERENCES

- [1] Coombs, C.F., Jr., ed., *Printed Circuits Handbook*, McGraw-Hill, New York, 1979.
- [2] *Electronic Designer's Handbook*, McGraw-Hill, New York.
- [3] Grace, Frank E., "Tolerance Study of Printed-Circuit Process Steps for a Microwave Application," *IEEE Transaction on Parts, Materials, and Packaging*, Vol. PMP-1, June 1965, No. 1, pp. S-16-S-27.
- [4] Waite, John M., "Impedance Control Discussion," Unpublished, December 6, 1989.

Chapter 10

Cable Dielectrics

The dielectrics used in constructing cables are, in general, different from those used in printed-circuit boards and ICs. Table 10.1 presents constants that are useful when calculating the performance of coaxial cables, twisted pairs, and other cables. It is also helpful when a custom cable is to be designed. These constants can be used with the equations of Chapter 3 to make the required calculations.

Material	ϵ_r	$\tan \delta$
Air	1.0006	
Cellular FEP	1.50	0.0007
Cellular TFE	1.40	0.0002
Cellular Polyethylene	1.40–2.10	0.0003
Ethylene Propylene	2.24	0.00046
Expanded PTFE ¹ (Teflon)	1.93	
FEP	2.10	0.00007
KEL-F ²	2.37	0.027–0.0053
Kynar	6.4	
Nylon	4.60–3.5	0.04–0.03
Perforated TFE	1.50	0.0002
Polyimide	3–3.5	0.0001
PFA	2.06	
PVC	3–8	
Si Rubber	2.08	0.007–0.016
SiO ₂	3.78	0.0001
Solid Polyethylene	2.35	0.0003
Solid PTFE	2.10	0.0002
Spline PTFE	1.35	0.0002
Tefzel	2.6	0.005

Table 10.1: Cable Dielectrics

¹Polytetrafluoroethylene

²Polymonochlorotrifluoroethylene

Chapter 11

Wire Gauges

It is convenient to use standard wire sizes when building slab lines, twisted pair, and other transmission lines. Wire can be used as high-impedance lines (substituting for small PC traces, which would be difficult or impossible to fabricate), inductors, and as conductors in various multiple-line configurations. Bond wires are used to interconnect substrates and to connect substrates to their package leads.

In addition to designing for the proper impedance or inductance it is important to make use of sizes that are standard and use a gauge and insulation sufficient for the expected power dissipation. Power dissipation numbers are approximate—they depend heavily on the physical situation (*i.e.*, length, airflow, presence of other heated conductors). The following tables provide a starting point for making these design trade-offs.

11.1 FUSING CURRENT

Fusing currents of wire are tabulated in Table 11.2.1. For CAE use the fusing current (long wire in 25°C still air) can be approximated for copper wire with [1]:

$$I_{fusing} = 80.0 d^{1.5}$$

and for aluminum wire

$$I_{fusing} = 59.2 d^{1.5}$$

REFERENCE

- [1] Giacoletto, L.J., ed., *Electronics Designers' Handbook*, 2nd Edition, McGraw-Hill, New York.

11.2 WIRE GAUGE

Wires are commonly referenced by their gauge. Table 11.2.1 tabulates the gauges and corresponding wire diameters. For CAE, the wire gauge is related to the wire diameter by the following approximation [3]:

$$d \approx 0.324883 e^{-0.115947 g}$$

or

$$g \approx \frac{-\ln\left(\frac{d}{0.324883}\right)}{0.115947}$$

where g is the wire gauge and d is the wire diameter in inches. For zero and negative g the gauges follow the pattern: 0 = 0 gauge, -1 = 00 gauge, -2 = 000 gauge, etc.

AWG B & S Gauge	Diameter (inches)	Aluminum Wire Fusing Current (A)	Copper Wire Fusing Current (A)
6	0.1620	495	668
7	0.1443	416	561
8	0.1285	349	472
9	0.1144	293	396
10	0.1019	247	333
11	0.09074	207	280
12	0.08081	174	235
13	0.07196	146	197
14	0.06408	123	166
15	0.05707	103	140
16	0.0582	86.8	117
17	0.04526	72.9	98.4
18	0.04030	61.4	82.9
19	0.03589	51.6	69.7
20	0.03196	43.2	58.4
21	0.02846		
22	0.02535	30.5	41.2
23	0.02257		
24	0.02010	21.6	29.2
25	0.01790		
26	0.01594	15.2	20.5
27	0.01420		
28	0.01264	10.7	14.4
29	0.01126		
30	0.0103	7.58	10.2
31	0.008928		
32	0.007950	5.32	7.19
33	0.007080		
34	0.006305	3.79	5.12
35	0.005615		4.28
36	0.005000	2.68	3.62
37	0.004453		
38	0.003965	1.85	2.50
39	0.003531		
40	0.003145	1.31	1.77
41	0.0028		1.52
42	0.0025		1.28
43	0.0022		1.06
44	0.0020		0.916

Table 11.2.1: Wire Gauges and Fusing Currents

REFERENCES

- [1] Giacoletto, L.J., ed., *Electronics Designers' Handbook*, 2nd Edition, McGraw-Hill, New York.
- [2] *Reference Data For Radio Engineers*, Howard W. Sams, Indianapolis, Fifth Edition, 1982.

- [3] Williamsen, Mark, "Calculating Wire Design Factors," *Machine Design*, October 12, 1989, p. 99.

11.3 BOND WIRES

Bond Wire	Material	Fusing Current ¹	Usage	Typical Resisitivity
1 mil	99.99% Gold	0.74 A	thermocompression or thermosonic bonding	1.17 mΩ/mil
1.3 mil	99.99% Gold	0.95 A	thermocompression or thermosonic bonding	.54 mΩ/mil
.7 mil	99.99% Gold		thermocompression or thermosonic bonding	
1 mil	Al with 1% Si	0.47 A	ultrasonic bonding	1.50 mΩ/mil
1.3 mil	Al with 1% Si	0.66 A	ultrasonic bonding	1.00 mΩ/mil
5 × 1 mil	99.99% Gold	3.50 A	ultrasonic bonding	0.1 mΩ/mil

Round bond wires are available in diameters of 0.5, 0.7, 1.0, 1.25, 1.5, 2, and 3 mils. Ribbon is available in (thickness × width) 0.8 × 1.6, 1.0 × 2.0, 1.5 × 3.0, 2.0 × 4.0, 2.0 × 10.0, 2.0 × 20.0, 2.0 × 30.0, and as small as 0.25 × 1.5 mils.

REFERENCES

- [1] Boulanger, Denis, "Wirebonding to Soft Substrates," *Microwave Journal*, February 1990, pp. 159–164.
- [2] Carlson, Jerry, "From Bonded Wire to Bumpless Ribbon/TAB Fine Pitch Interconnects," *Hybrid Circuit Technology*, April 1990, pp. 26–28.
- [3] Ham, R.E., "Prediction of Bond Wire Temperature Using an Electronic Circuit Analogy," *Hybrid Circuit Technology*, April 1990, pp. 53–54. (Describes an accurate thermal model suitable for predicting failure of bond wires subject to current pulses.)
- [4] Tektronix, "Thick-Film Hybrid Design Guidelines", pp. 3–6.

¹Fusing current is dependent on the wire's length and any heatsink or airflow. These numbers are guidelines.

Chapter 12

Special Functions

The functions in this chapter are common in transmission line problems. Often a problem that seems formidable at first look becomes quite easy to solve when a library of the desired routines is available. These functions have been successfully implemented as spreadsheet functions (Microsoft Excel™) as well as C subroutines. Many computers will have some or all of these available as libraries.

There are many ways to implement a given function, each of which has different speed–accuracy trade-offs. Depending on the task, optimization for speed may be required. A useful technique is to curve-fit a simple equation to the more complicated equation over a limited range.

The final topic covered here is the representation of circuits with chain matrices.

12.1 STANDARD FUNCTIONS

To build the special functions, standard functions are used that are available on most computers. The C math library includes all of the functions found in the file <math.h>. A typical list includes:

acos()	sinh()	log()	frexp()
atan2()	tanh()	pow()	log10()
cosh()	asinh()	sqrt()	sin()
floor()	ceil()	atan()	tan()
ldexp()	exp()	cos()	
modf()	fmod()	fabs()	

The meaning of most of the above is obvious. Parameters to trigonometric functions and return values of the inverse trigonometric functions are in radians. The absolute value of doubles is found with `fabs()`; use `abs()` for integers. Similarly, `fmod()` returns the remainder of its operands; use the binary operator `%` for integers.

The function `ldexp()` returns the mantissa and power of 2 required to represent the passed parameter; `frexp()` is its inverse function. The natural logarithm, `ln()` is the function `log()`.

12.2 SPECIAL FUNCTIONS

12.2.1 Hyperbolic Functions

The hyperbolic functions used in transmission line problems are simply related to the exponential function.

$$e^{j\theta} = \cos \theta + j \sin \theta$$

$$e^x = \cosh x + j \sinh x$$

$$\sinh x = \frac{e^x - e^{-x}}{2} = \frac{\sin jx}{j}$$

$$\cosh x = \frac{e^x + e^{-x}}{2} = \cos jx$$

$$\tanh x = \frac{e^x - e^{-x}}{e^x + e^{-x}} = \frac{\tan jx}{j}$$

$$\sinh^{-1} x = \ln (x + \sqrt{x^2 + 1.0})$$

$$\cosh^{-1} x = \ln (x + \sqrt{x^2 - 1.0}), \quad x \geq 1.0$$

$$\tanh^{-1} x = \frac{1}{2} \ln \left(\frac{1.0 + x}{1.0 - x} \right), \quad x^2 < 1.0$$

12.2.2 Elliptic Integrals

Elliptic integrals occur when evaluating several types of transmission lines, among them helical inductors and coplanar waveguide. We have encountered the elliptic integrals of the first and second kind and the complete integrals of the first and second kind.

The elliptic integral of the first kind is of the form

$$F(k, \phi) = \int_0^\phi \frac{d\Phi}{\sqrt{1.0 - k^2 \sin^2 \Phi}} = \int_0^x \frac{d\xi}{\sqrt{(1.0 - k^2)(1.0 - k^2 \xi^2)}}$$

where

$$x = \sin \phi$$

$$k^2 < 1.0$$

and the elliptic integral of the second kind is of the form

$$E(k, \phi) = \int_0^\phi \sqrt{1.0 - k^2 \sin^2 \Phi} d\Phi = \int_0^x \frac{\sqrt{1.0 - k^2 \xi^2} d\xi}{\sqrt{1.0 - \xi^2}}$$

where

$$x = \sin \phi$$

$$k^2 < 1.0$$

When the parameter ϕ is set to $\pi/2$ we define the complete elliptic integrals of the first $[K(k)]$ and second kind $[E(k)]$. The parameter ϕ is implicit in these cases.

12.2.2.1 Complete Elliptic Integral of the First Kind

The complete elliptic integral of the first kind is encountered both unprimed and primed. The prime is a shorthand form for the parameter k .

$$K'(k) = K(k')$$

$$k' = \sqrt{1 - k^2}$$

Miller [6] gives the following recursive equation for use in calculating inductance of helical coils, which is equally applicable for use with other structures using $K(k)$.

$$K(k) = \frac{\pi}{2 a_N}$$

where

$$a_n = \frac{a_{n-1} + b_{n-1}}{2}$$

$$b_n = \sqrt{a_{n-1} b_{n-1}}$$

$$c_n = \frac{a_{n-1} - b_{n-1}}{2}$$

and

$$a_0 = 1.0$$

$$b_0 = \sqrt{1.0 - k^2}$$

$$c_0 = k$$

To calculate the complete elliptic integral of the first kind, start the calculation with the initial values a_0 , b_0 , and c_0 , and iterate until $c_N = 0$ to within the desired accuracy.

12.2.2.2 Complete Elliptic Integral of the Second Kind

The complete elliptic integral of the second kind, $E(k)$, is also found both primed and unprimed. Again, this is just a shorthand form for the parameters.

$$E'(k) = E(k')$$

$$k' = \sqrt{1 - k^2}$$

As before, a simple recursion is given by Miller [6]

$$E(k) = F(k) \left(1.0 - 0.5 \sum_{n=0}^N 2^n c_n^2 \right)$$

where

$$F(k) = \frac{\pi}{2 a_N}$$

$$a_n = \frac{a_{n-1} + b_{n-1}}{2}$$

$$b_n = \sqrt{a_{n-1} b_{n-1}}$$

$$c_n = \frac{a_{n-1} - b_{n-1}}{2}$$

and

$$a_0 = 1.0$$

$$b_0 = \sqrt{1.0 - k^2}$$

$$c_0 = k$$

Start the calculation with the initial values a_0 , b_0 , and c_0 , and iterate until $c_N = 0$ to within the desired accuracy.

12.2.2.3 Ratios of Complete Elliptic Integrals of the First Kind

A form that commonly occurs in transmission line problems is the ratio $K(k) / K(k')$. This ratio has been approximated by Hilberg [4] to very good accuracy using simple functions. For $1.0 \leq K(k) / K(k') \leq \infty$ and $1.0 / \sqrt{2.0} \leq k \leq 1.0$:

$$\frac{K(k)}{K(k')} \approx \frac{1.0}{2.0 \pi} \ln \left(2.0 \frac{\sqrt{1.0 + k} + \sqrt[4]{4.0 k}}{\sqrt{1.0 + k} - \sqrt[4]{4.0 k}} \right)$$

For $0.0 \leq K(k) / K(k') \leq 1.0$ and $0.0 \leq k \leq 1.0 / \sqrt{2.0}$

$$\frac{K(k)}{K(k')} \approx \frac{2.0 \pi}{\ln \left(2.0 \frac{\sqrt{1.0 + k'} + \sqrt[4]{4.0 k'}}{\sqrt{1.0 + k'} - \sqrt[4]{4.0 k'}} \right)}$$

These equations are accurate to within 4×10^{-12} .

REFERENCES

- [1] Abramowitz, Milton, and Irene A. Stegun, *Handbook of Mathematical Functions*, Dover Publications, New York, 1965. (If you are going to be using special functions a lot, this is a great all-in-one reference. Includes definitions, graphs, tables, approximation formulas, and more.)
- [2] Fettis, Henry E., and James C. Casling, *Tables of Elliptic Integrals of the First, Second and Third Kind*, Aerospace Research Laboratories, ARL 64-232, December 1964.
- [3] Hart, John F., et al., *Computer Approximations*, John Wiley & Sons, New York, 1968.
- [4] Hilberg, Wolfgang, "From Approximations to Exact Relations for Characteristic Impedances," *IEEE Transactions on Microwave Theory and Techniques*, Vol. MTT-17, No. 5, May 1969, pp. 259-265.

- [5] Luke, Yudell L., *The Special Functions and their Approximations*, Vol. II, Academic Press, NY, 1969. (Has a large number of tables of coefficients of Chebyshev polynomials for approximating functions ranging from $1/x$ to the special functions, and their integrals.)
- [6] Miller, H. Craig, "Inductance Formula for a Single-Layer Circular Coil," *Proceedings of the IEEE*, Vol. 75, No. 2, February 1987, pp. 256–257.
- [7] Press, William H., et al., *Numerical Recipes in C: The Art of Scientific Computing*, Cambridge University Press, New York, 1988. (This is the reference book for software which must solve any type of numerical problem. It is also available in BASIC, and FORTRAN versions.)

12.2.3 Bessel Functions

Bessel functions are solutions of the equation

$$x^2 \frac{d^2 y}{dx^2} + x \frac{dy}{dx} + (x^2 - \nu^2)y = 0 \quad (12.2.3.1)$$

There are several types of solutions:

$$J_\nu(x) = \text{Bessel function of the first kind of order } \nu \quad (12.2.3.2)$$

$$Y_\nu(x) = \text{Bessel function of the second kind of order } \nu \quad (12.2.3.3)$$

$$H_\nu(x) = \text{Bessel function of the third kind of order } \nu \quad (12.2.3.4)$$

The modified Bessel functions are solutions of

$$x^2 \frac{d^2 y}{dx^2} + x \frac{dy}{dx} - (x^2 + \nu^2)y = 0 \quad (12.2.3.5)$$

Again there are several solutions:

$$I_\nu(x) = \text{modified Bessel function of the first kind of order } \nu \quad (12.2.3.6)$$

$$K_\nu(x) = \text{modified Bessel function of the second kind of order } \nu \quad (12.2.3.7)$$

Three types of problems occur in transmission line use of Bessel functions:

- value of Bessel function is required
- zeroes of the Bessel function are required
- zeroes of combinations of Bessel functions

Some useful relations among the Bessel functions are:

$$J_{\nu+1}(x) Y_\nu(x) - J_\nu(x) Y_{\nu+1}(x) = \frac{2.0}{\pi x} \quad (12.2.3.8)$$

$$J_{\nu+1}(x) = \frac{2.0}{x} J_\nu(x) - J_{\nu-1}(x) \quad (12.2.3.9)$$

$$Y_{\nu+1}(x) = \frac{2.0}{x} Y_\nu(x) - Y_{\nu-1}(x) \quad (12.2.3.10)$$

12.2.3.1 Values of Bessel Functions

A number of numerical techniques are found in [1, 5] for these functions. Vaughan [6] has given a simple technique of calculating four of the Bessel functions for integer orders. C subroutine listings are given below (the original source is BASIC):

```
double
bessj(n, x)
int      n;
double   x;
{
    double   J1 = 1.0;
    double   J0 = 1.0;
    double   Q0 = 0.0;
    double   Q1 = 1.0;
    double   Q2 = 1.0;
    double   Q3 = 1.0;
    double   Q6 = 0.0;
    double   Q7;
    double   array[6];
    int      j;

    for(j = 0; j < 6; j++)
        array[j] = 0.0;

    Q7 = x * x / 4.0;

    /* Calculate J0 and J1 */
    do {
        Q6++;
        Q1 = -Q1;
        Q2 *= Q7 / (Q6 * Q6);
        Q3 *= Q7 / (Q6 * (Q6 + 1));

        J0 += Q1 * Q2;
        J1 += Q1 * Q3;
    } while(Q2 > 1.0e-05);

    J1 *= x/2.0;

    /* Do the rest of the array */
    array[0] = J0;
    array[1] = J1;
    for(j = 2; j <= n; j++) {
        array[j] = ((2.0 * (j - 1)) / x) * array[j - 1]
```

```

        - array[j - 2];
    }
    return array[n];
}

double
bessy(n, x)
int      n;
double   x;
{
    double   Y0 = 0.0;
    double   Y1 = 1.0;
    double   J1 = 1.0;
    double   J0 = 1.0;
    double   Q0 = 0.0;
    double   Q1 = 1.0;
    double   Q2 = 1.0;
    double   Q3 = 1.0;
    double   Q4, Q5;
    double   Q6 = 0.0;
    double   Q7;
    double   array[6];
    int      j;

    for(j = 0; j < 6; j++)
        array[j] = 0.0;

    Q7 = x * x / 4.0;

    /* Calculate Y0 and Y1 */
    do{
        Q6++;
        Q0 += 1.0 / Q6;
        Q1 = -Q1;
        Q2 = (Q2 * Q7) / (Q6 * Q6);
        Q3 = (Q3 * Q7) / (Q6 * (Q6 + 1.0));
        Q4 = Q3 * ((2.0 * Q0 + 1.0) / (Q6 + 1.0));

        Y0 += Q1 * Q2 * Q0;
        Y1 += Q1 * Q4;
        J0 += Q1 * Q2;
        J1 += Q1 * Q3;
    }while(Q2 > 1.0e-05);

    J1 *= x/2.0;
    Q5 = CONST1 + log(x/2.0);
    Y0 = CONST2 * (J0 * Q5 - Y0);
    Y1 = CONST2 * (J1 * Q5 - (1.0 / x) - (Y1 * x / 4.0));

    /* Do the rest of the array */
    array[0] = Y0;
    array[1] = Y1;
    for(j = 2; j <= n; j++){
        array[j] = ((2.0 * (j - 1)) / x) * array[j - 1] -
array[j - 2];
    }
}

```

```

    }
    return array[n];
}

double
bessi(n, x)
int      n;
double   x;
{
    double   I1 = 1.0;
    double   I0 = 1.0;
    double   Q0 = 0.0;
    double   Q1 = 1.0;
    double   Q2 = 1.0;
    double   Q3 = 1.0;
    double   Q6 = 0.0;
    double   Q7;
    double   array[6];
    int      j;

    for(j = 0; j < 6; j++)
        array[j] = 0.0;

    Q7 = x * x / 4.0;

    /* Calculate I0 and I1 */
    do{
        Q6++;
        Q2 *= Q7 / (Q6 * Q6);
        Q3 *= Q7 / (Q6 * (Q6 + 1.0));

        I0 += Q2;
        I1 += Q3;
    }while(Q2 > 1.0e-05);

    I1 *= x/2.0;

    /* Do the rest of the array */
    array[0] = I0;
    array[1] = I1;
    for(j = 2; j <= n; j++){
        array[j] = ((-2.0 * (j - 1)) / x) * array[j - 1]
+ array[j - 2];
    }
    return array[n];
}

double
bessk(n, x)
int      n;
double   x;
{
    double   K0 = 0.0;
    double   K1 = 1.0;
    double   I1 = 1.0;
    double   I0 = 1.0;

```



```

double    Q0 = 0.0;
double    Q1 = 1.0;
double    Q2 = 1.0;
double    Q3 = 1.0;
double    Q4, Q5;
double    Q6 = 0.0;
double    Q7;
int       j;
double    array[6];

for(j = 0; j < 6; j++)
    array[j] = 0.0;

Q7 = x * x / 4.0;

/* Calculate Y0 and Y1 */
do{
    Q6++;
    Q0 += 1 / Q6;
    Q2 += Q7 / (Q6 * Q6);
    Q3 += Q7 / (Q6 * (Q6 + 1));
    Q4 = Q3 * ((2 * Q0 + 1) / (Q6 + 1));

    K0 += Q2 * Q0;
    K1 += Q4;
    I0 += Q2;
    I1 += Q3;

}while(Q2 > 1.0e-05);

I1 *= x/2.0;
Q5 = CONST1 + log(x/2.0);
K0 -= I0 * Q5;
K1 = I1 * Q5 + (1.0 / x) - (K1 * x / 4.0);

/* Do the rest of the array */
array[0] = K0;
array[1] = K1;
for(j = 2; j <= n; j++){
    array[j] = ((2.0 * (j - 1)) / x) * array[j - 1]
        + array[j - 2];
}
return array[n];
}

```

Figure 12.2.3.1.1: Bessel Function Source Code

The equations are valid for

$$0.0 < x \leq 10.0$$

$$0 \leq \nu \leq 5 \quad (\text{integers})$$

Accuracy is to five places or five significant digits except for small x and large ν (4, 5).

12.2.3.2 Zeroes of Bessel Functions

The s th real zero of Bessel functions can be found by a root-finding program in conjunction with the relations of the previous section or by polynomial expansions that are available. The first zeroes of the Bessel functions of the first and second kind are given approximately by [1]:

$$j_{\nu,1} \equiv \nu + 1.8557571 \nu^{1/3} + 1.033150 \nu^{-1/3} - 0.00397 \nu^{-1} - 0.0908 \nu^{-5/3} + 0.043 \nu^{-7/3} \quad (12.2.3.2.1)$$

$$y_{\nu,1} \equiv \nu + 0.9315768 \nu^{1/3} + 0.260351 \nu^{-1/3} - 0.01198 \nu^{-1} - 0.0060 \nu^{-5/3} + 0.001 \nu^{-7/3} \quad (12.2.3.2.2)$$

12.2.3.3 Zeroes of Combination of Bessel Functions

Another case that occurs in waveguide problems is the solution of

$$J_{\nu}(x) Y_{\nu}(kx) - J_{\nu}(kx) Y_{\nu}(x) = 0 \quad (12.2.3.3.1)$$

which is also found rearranged as

$$\frac{J_{\nu}(x)}{J_{\nu}(kx)} = \frac{Y_{\nu}(x)}{Y_{\nu}(kx)}$$

Gunston [4] and Cochran [3] have found simple relations for the first zeroes of this equation.

$$z_{\nu,s} \equiv \sqrt{\frac{(s\pi)^2}{(k-1.0)^2} + \frac{4.0 \nu^2 - 1.0}{(k+1.0)^2}} \quad (12.2.3.3.2)$$

A similar problem is

$$J_{\nu}'(x) Y_{\nu}'(kx) - J_{\nu}'(kx) Y_{\nu}'(x) = 0 \quad (12.2.3.3.3)$$

again, this can be rearranged to the equivalent form

$$\frac{J_{\nu}'(x)}{J_{\nu}'(kx)} = \frac{Y_{\nu}'(x)}{Y_{\nu}'(kx)}$$

The primed quantities are derivatives with respect to x . The solution of the equation is found with

$$z'_{v,s} \equiv \sqrt{\frac{(s\pi)^2}{(k-1.0)^2} + \frac{4.0 v^2 + 3.0}{(k+1.0)^2}} \quad (12.2.3.3.4)$$

$$z'_{v,0} \equiv \frac{2.0 v}{(k+1.0)} \left[\frac{1.0 + (k-1.0)^2}{6.0 (k+1.0)^2} \right] \quad (12.2.3.3.5)$$

where the z_s are the s th zeroes of the corresponding combination of Bessel function equations. Accuracy of the above is dependent on the combination of parameters. See [3, 4] for more information. In general accuracy is better than a few percent and improves with increasing s and v and decreasing k . See [1, p. 374] for another solution of this problem.

12.2.3.4 Kelvin Functions

The Kelvin functions are the real and imaginary parts of Bessel functions with imaginary arguments.

$$\text{ber}_v x + j \text{bei}_v x = J_v(x e^{3\pi j/4}) \quad (12.2.3.4.1)$$

$$\text{ker}_v x + j \text{kei}_v x = e^{-v\pi j/2} K_v(x e^{\pi j/4}) \quad (12.2.3.4.2)$$

which are defined for

$$v = \text{real}$$

$$x = \text{real}, \geq 0$$

$$n = \text{integer}, \geq 0$$

See the references for more information on calculating these functions.

REFERENCES

- [1] Abramowitz, Milton, and Irene A. Stegun, *Handbook of Mathematical Functions*, Dover Publications, New York, 1965. (If you are going to be using special functions a lot, this is a great all-in-one reference. Includes definitions, graphs, and tables, approximation formulas, and more.)
- [2] Balanis, Constantine A., *Advanced Engineering Electromagnetics*, John Wiley & Sons, New York, 1989.
- [3] Cochran, J.A., "Further Formulas for Calculating Approximate Values of the Zeros of Certain Combinations of Bessel Functions," *IEEE Transactions on*

Microwave Theory and Techniques, Vol. MTT-11, No. 11, November 1963, pp. 546-547. (Extends Gunston, and corrects errors in accuracy.)

- [4] Gunston, M.A.R., "A Simple Formula for Calculating Approximate Values of the First Zeros of a Combination Bessel Function Equation," *IEEE Transactions on Microwave Theory and Techniques*, Vol. MTT-11, No. 1, January 1963, pp. 93-94.
- [5] Press, William H., et al., *Numerical Recipes in C: The Art of Scientific Computing*, Cambridge University Press, New York, 1988. (This is the reference book for software that must solve any type of numerical problem. It is also available in BASIC, and FORTRAN versions.)
- [6] Vaughan, J. Rodney M., "Routine to Supply Bessel Functions as Required in BASIC Programs," *IEEE Transactions on Microwave Theory and Techniques*, Vol. MTT-18, No. 8, August 1970, pp. 511-512.

12.2.4 Sine-Integral Function

The sine integral function was encountered in coaxial discontinuity relations.

$$Si(x) = \int_0^x \frac{\sin y}{y} dy \quad (12.2.4.1)$$

An approximation formula [1] for $0.0 \leq x \leq 1.0$ is

$$Si(x) \equiv x - \frac{x^3}{18.0} + \frac{x^5}{600.0} - \frac{x^7}{35,280.0} + \frac{x^9}{3,265,920.0} \quad (12.2.4.2a)$$

and for $1.0 \leq x \leq \infty$

$$Si(x) \equiv \frac{\pi}{2.0} - f(x) \cos(x) - g(x) \sin(x) \quad (12.2.4.2b)$$

where

$$f(x) = \frac{1.0}{x} \frac{x^8 + 38.027264 x^6 + 265.187033 x^4 + 335.677320 x^2 + 38.102495}{x^8 + 40.021433 x^6 + 322.624911 x^4 + 570.236280 x^2 + 157.105423} \quad (12.2.4.3)$$

$$g(x) = \frac{1.0}{x^2} \frac{x^8 + 42.2242855 x^6 + 302.757865 x^4 + 352.018498 x^2 + 21.821899}{x^8 + 48.196927 x^6 + 482.485984 x^4 + 1114.978885 x^2 + 449.690326} \quad (12.2.4.4)$$

These equations give results that are accurate to better than 1 ppm.

REFERENCES:

- [1] Abramowitz, Milton, and Irene A. Stegun, *Handbook of Mathematical Functions*, Dover Publications, New York, 1965. (If you are going to be using special functions a lot, this is a great all-in-one reference. Includes definitions, graphs, tables, approximation formulas, and more.)

12.2.5 Gamma Function

The gamma function is used alone and as part of other functions. The definition is

$$\Gamma(z) = \int_0^{\infty} t^{z-1.0} e^{-t} dt \quad (12.2.5.1)$$

For integer z , $\Gamma(z)$ equals $(z-1)!$. An approximation formula is

$$\log \Gamma(x) = (x-1) \log x - x + \frac{1}{2} \log(2\pi) + \frac{1}{2} \log\left(x + \frac{1}{6}\right) + E(x) \quad (12.2.5.2)$$

where for large x

$$E(x) = \left(\frac{1}{12x + \frac{4.6}{15}} \right)^2 \quad (12.2.5.3)$$

This approximation to $\Gamma(x)$ is accurate to within approximately 2% for all x .

REFERENCES

- [1] Lewin, L., "An Accurate Formula for the Gamma Function," *IEEE Transactions on Microwave Theory and Techniques*, Vol. MTT-22, No. 10, October 1974, p. 910. (Corrects the usual approximation formula to improve accuracy near $x = 0$.)

12.2.6 Hecken Impedance Function

A function used in calculating optimal impedance tapers is the Hencken impedance function defined as

$$G(B, \xi) = \frac{B}{\sinh B} \int_0^{\xi} I_0(B \sqrt{1.0 - \xi^2}) d\xi \quad (12.2.6.1)$$

where $I_0(z)$ is the zeroth-order modified Bessel function of the first kind whose calculation is described earlier. This function can be calculated using the recursion of [1]:

$$G(B, \xi) = \frac{B}{\sinh B} \sum_{k=0}^{\infty} a_k b_k \quad (12.2.6.2)$$

where

$$a_0 = 1.0 \quad (12.2.6.3)$$

$$a_k = \left(\frac{B^2}{4.0 k^2} \right) a_{k-1}, \quad k = 1, 2, \dots \quad (12.2.6.4)$$

$$b_0 = \xi \quad (12.2.6.5)$$

$$b_k = \frac{\xi(1.0 - \xi^2)^k + 2k b_{k-1}}{2.0 k + 1.0} \quad (12.2.6.6)$$

REFERENCES

- [1] Cloete, J.H., "Computation of the Hecken Impedance Function," *IEEE Transactions on Microwave Theory and Techniques*, Vol. MTT-25, No. 5, May 1977, p. 440.

12.3 TWO-PORT REPRESENTATIONS

RF and microwave circuits can be represented by a number of different matrices that describe the two-port network. The most commonly used in measurements is the s -parameter or scattering matrix. For CAE calculations it is convenient to have a matrix where the product of the component matrices produces the response of a circuit. These matrices are the ABCD and the T matrix. Matrix operations also allow us to calculate the result of two paralleled two-ports. Conversion among these parameters and low frequency parameters are also possible.

Some circuits have three or more ports. For example, the directional coupler is a four-port circuit.

12.3.1 Scattering or s -Parameters

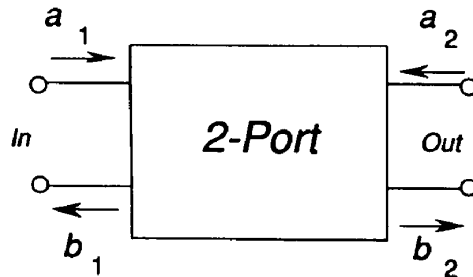


Figure 12.3.1.1: s -Parameter Definition

The s -parameters are defined using a two-port model as shown in Figure 12.3.1.1. The equations are

$$b_1 = s_{11} a_1 + s_{12} a_2 \quad (12.3.1.1)$$

$$b_2 = s_{21} a_1 + s_{22} a_2 \quad (12.3.1.2)$$

where¹

$$s_{11} = \left. \frac{b_1}{a_1} \right|_{a_2 = 0} \quad (12.3.1.3)$$

$$s_{12} = \left. \frac{b_2}{a_1} \right|_{a_2 = 0} \quad (12.3.1.4)$$

$$s_{21} = \left. \frac{b_2}{a_2} \right|_{a_1 = 0} \quad (12.3.1.5)$$

$$s_{22} = \left. \frac{b_1}{a_2} \right|_{a_1 = 0} \quad (12.3.1.6)$$

and

a_n = incident wave

b_n = reflected wave

Conversions to T -parameters and s -parameters are given below.

12.3.2 Transmission or T -Parameters

The T -parameters are defined with the same figure as the s -parameters.

$$a_1 = T_{11} b_2 + T_{12} a_2 \quad (12.3.2.1)$$

$$b_1 = T_{21} b_2 + T_{22} a_2 \quad (12.3.2.2)$$

These parameters can be derived in terms of the s -parameters and have the added advantage that two-port networks described by T -parameters can be chained by simply multiplying the matrices.

To transfer to and from s -parameters, the relations are

$$\begin{bmatrix} T_{11} & T_{12} \\ T_{21} & T_{22} \end{bmatrix} = \begin{bmatrix} \frac{1.0}{s_{21}} & \frac{-s_{22}}{s_{21}} \\ \frac{s_{11}}{s_{21}} & s_{12} - \frac{s_{11}s_{22}}{s_{21}} \end{bmatrix} \quad (12.3.2.3)$$

$$\begin{bmatrix} s_{11} & s_{12} \\ s_{21} & s_{22} \end{bmatrix} = \begin{bmatrix} \frac{T_{21}}{T_{11}} & T_{22} - \frac{T_{21}T_{12}}{T_{11}} \\ \frac{1.0}{T_{11}} & \frac{-T_{12}}{T_{11}} \end{bmatrix} \quad (12.3.2.4)$$

¹Note that $a_n = 0$ is equivalent to matching port n .

12.3.3 ABCD Matrices

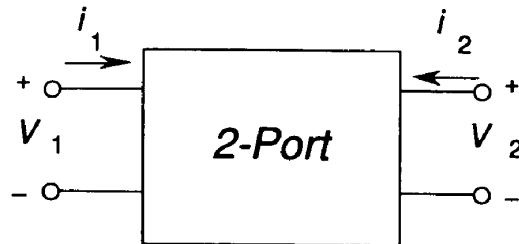


Figure 12.3.3.1: ABCD Parameter Definitions

The equations describing the $ABCD$ matrix are:

$$v_1 = A v_2 - B i_2 \quad (12.3.3.1)$$

$$i_1 = C v_2 - D i_2 \quad (12.3.3.2)$$

The relationship to the s -parameters is not as straightforward as for T -parameters, but the $ABCD$ matrix shares the ability to chain matrices to build up a circuit. To convert between s -parameters and $ABCD$ parameters

$$\begin{bmatrix} A & B \\ C & D \end{bmatrix} = \begin{bmatrix} \frac{(1 + s_{11})(1 - s_{22}) + s_{12} s_{21}}{2 s_{21}} & \frac{(1 + s_{11})(1 + s_{22}) + s_{12} s_{21}}{2 s_{21}} \\ \frac{(1 - s_{11})(1 - s_{22}) + s_{12} s_{21}}{2 s_{21}} & \frac{(1 - s_{11})(1 + s_{22}) + s_{12} s_{21}}{2 s_{21}} \end{bmatrix} \quad (12.3.3.3)$$

$$\begin{bmatrix} s_{11} & s_{12} \\ s_{21} & s_{22} \end{bmatrix} = \begin{bmatrix} \frac{A + B - C - D}{A + B + C + D} & \frac{2.0 (A D - B C)}{A + B + C + D} \\ \frac{2.0}{A + B + C + D} & \frac{-A + B - C + D}{A + B + C + D} \end{bmatrix} \quad (12.3.3.4)$$

REFERENCES

- [1] Choma, Jr., J., "Signal Flow Analysis of Feedback Networks," *IEEE Transactions on Circuits and Systems*, Vol. CS-37, No. 4, April 1990, pp. 455–463.
- [2] Gonzalez, Guillermo, *Microwave Transistor Amplifiers: Analysis and Design*, Prentice-Hall, NJ, 1984. (A very succinct and thorough book with good coverage of s -parameter design applied to amplifiers and oscillators.)
- [3] Liao, Samuel Y., *Microwave Circuit Analysis and Amplifier Design*, Prentice-Hall, NJ, 1987.
- [4] Teyssier, Edward M., "Two-Port to Three-Port S-Parameter Conversion," *MSN&CT*, September 1987, pp. 84–91.
- [5] Young, Leo, and H. Sobol, eds., *Advances in Microwaves*, Volume 8, "Computer Aided Design, Simulation and Optimization," by B.S. Perlman and V.G. Gelnovatch, Academic Press, New York, 1974.

12.3.4 Series Combinations of Two-Port Networks

To analyze a circuit made up of transmission lines and lumped and active elements, we multiply the individual *ABCD* or *T*-parameter matrices. The product matrix then represents the matrix of the combination. Any arbitrary set of *s*-parameters can be converted with the above equations into the required chain matrix form. The multiplication of two 2×2 matrices is

$$\begin{aligned}
 [A][B] &= \begin{bmatrix} a_{00} & a_{01} \\ a_{10} & a_{11} \end{bmatrix} \begin{bmatrix} b_{00} & b_{01} \\ b_{10} & b_{11} \end{bmatrix} \\
 &= \begin{bmatrix} a_{00}b_{00} + a_{01}b_{10} & a_{00}b_{01} + a_{01}b_{11} \\ a_{10}b_{00} + a_{11}b_{10} & a_{10}b_{01} + a_{11}b_{11} \end{bmatrix} \quad (12.3.4.1)
 \end{aligned}$$

The chain matrix can then be converted to *s*-parameters, impedance, gain, etc., as required.

12.3.5 Parallel Combinations of Two-Port Networks

In many cases it is desirable to connect two-port networks in other than simple cascades. To do this we can transform the desired two-ports using the following *s*-parameter matrix relations.

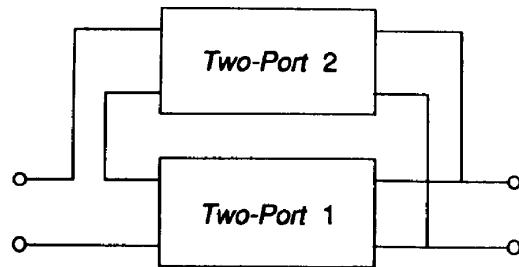


Figure 12.3.5.1: Series-Parallel Connection of Two-Port Networks

$$S_{sp} = \begin{bmatrix} -1.0 & 0.0 \\ 0.0 & 1.0 \end{bmatrix} \times f \left(\begin{bmatrix} -1.0 & 0.0 \\ 0.0 & 1.0 \end{bmatrix} \times S_1, \begin{bmatrix} -1.0 & 0.0 \\ 0.0 & 1.0 \end{bmatrix} \times S_2 \right) \quad (12.3.5.1)$$

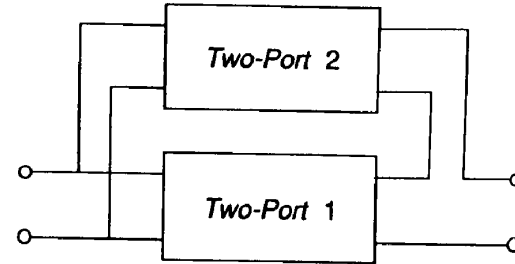


Figure 15.3.5.2: Parallel-Series Connection of Two-Port Networks

$$S_{ps} = \begin{bmatrix} 1.0 & 0.0 \\ 0.0 & -1.0 \end{bmatrix} \times f \left(\begin{bmatrix} 1.0 & 0.0 \\ 0.0 & -1.0 \end{bmatrix} \times S_1, \begin{bmatrix} 1.0 & 0.0 \\ 0.0 & -1.0 \end{bmatrix} \times S_2 \right) \quad (12.3.5.2)$$

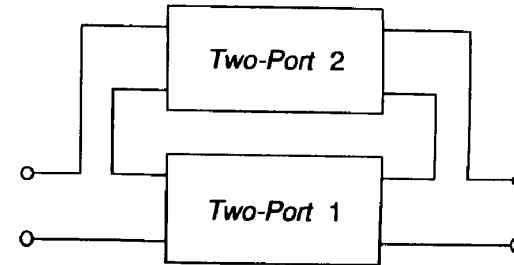


Figure 15.3.5.3: Series-Series Connection of Two-Ports

$$S_{ss} = \begin{bmatrix} -1.0 & 0.0 \\ 0.0 & -1.0 \end{bmatrix} \times f \left(\begin{bmatrix} -1.0 & 0.0 \\ 0.0 & -1.0 \end{bmatrix} \times S_1, \begin{bmatrix} -1.0 & 0.0 \\ 0.0 & -1.0 \end{bmatrix} \times S_2 \right) \quad (12.3.5.3)$$

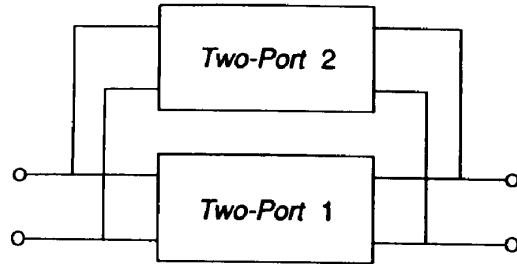


Figure 15.3.5.4.: Parallel-Parallel Connection of Two-Port Networks

$$S_{pp} = \begin{bmatrix} 1.0 & 0.0 \\ 0.0 & 1.0 \end{bmatrix} \times f \left(\begin{bmatrix} 1.0 & 0.0 \\ 0.0 & 1.0 \end{bmatrix} \times S_1, \begin{bmatrix} 1.0 & 0.0 \\ 0.0 & 1.0 \end{bmatrix} \times S_2 \right) \quad (12.3.5.4)$$

In the above relations

$$A = 3 \times I \quad (12.3.5.5)$$

$$B = S_1 - I \quad (12.3.5.6)$$

$$C = S_1 + I \quad (12.3.5.7)$$

$$f(S_1, S_2) = A^{-1} \times \left[B + 4.0 \times C \times S_2 \times (A - B \times S_2)^{-1} \times C \right] \quad (12.3.5.8)$$

$$I = \begin{bmatrix} 1.0 & 0.0 \\ 0.0 & 1.0 \end{bmatrix}, \text{ the identity matrix} \quad (12.3.5.9)$$

$$S_1 = \text{the } s\text{-parameter matrix of two-port 1} \quad (12.3.5.10)$$

$$S_2 = \text{the } s\text{-parameter matrix of two-port 2} \quad (12.3.5.11)$$

REFERENCES

- [1] Bodharamik, Phichani, *et al.*, "Two Scattering Matrix Programs for Active Circuit Analysis," *IEEE Transactions on Circuit Theory*, Vol. CT-18, No. 6, November 1971, pp. 610–619.
- [2] Vendelin, George D., *et al.*, *Microwave Circuit Design Using Linear and Nonlinear Techniques*, John Wiley & Sons, New York, 1990.

Glossary

ANISOTROPIC. A material having parameters (σ , ϵ_r , etc.) that vary with direction in the material. For example, woven materials have one ϵ_r with the grain and another perpendicular to it.

B STAGE. Partially cured dielectric, used to bond the layers together

COMMENSURATE. Equal length.

CONFORMAL MAPPING. Conformal mapping is a technique used for analytically solving EM field problems that are otherwise too complex. The technique maps the physical structure into another domain via a transforming equation that preserves the structure's capacitance and Z_0 . The transformed equations are then easily solvable. This is analogous to a Mercator projection, which transforms the surface of the earth (essentially a sphere) onto a map (two-dimensional).

C STAGE. Fully cured epoxy; C stage dielectric usually has the foil already attached to the dielectric.

Effective. The effective relative dielectric constant taking into account that a portion of the wave is propagating in air ($\epsilon_r = 1$) and a portion in the dielectric (ϵ_r).

FEA OR FINITE ELEMENT

ANALYSIS. A technique used to solve Maxwell's equations that divides the transmission line space into a grid of small elements. A simplified version of Maxwell's equations is then applied iteratively to these elements until the solution is achieved.

FRINGING CAPACITANCE. The capacitance of a line's edges; the increased capacitance beyond the ideal parallel-plate capacitance; capacitance due to edge fields that do not reach from one edge to the other.

GENERALIZED SCATTERING

MATRIX TECHNIQUE. A technique of calculating discontinuities by combining the s -parameters of the simple discontinuities making up the discontinuity. The simple

discontinuity s -parameters are calculated with mode-matching or other techniques.

GREEN'S FUNCTIONS. A technique used to solve EM problems. In this technique, the driving function is a Dirac delta. The response to the delta is the Green's function. The actual driving function is then written as the sum or integral of offset Dirac deltas. The Green's function can then be used to solve for the response to the actual driving function.

MAGNETIC. A material is considered magnetic when μ_r is greater than one.

MAXWELL'S EQUATIONS. The physical laws that describe the relationships between electric and magnetic fields. They can be expressed in a set of differential or integral equations.

NEWTON'S METHOD. A method of finding roots of an equation which evaluates the local slope of the function and uses this to estimate the next guess at the root. Successive iterations quickly home in on the root's value.

NEWTON-RAPHSON METHOD.
See *Newton's Method* above.

NONHOMOGENEOUS. A dielectric or other material is non-homogeneous when its physical parameters vary with position in the material. For example, microstrip line has dielectric below the center conductor and air above—this is non-homogeneous.

SPECTRAL DOMAIN METHOD. A technique used for EM analysis of structures with thin conductors. The Fourier transform is applied converting the space equations to the spectral domain.

SUPERCONDUCTOR. A material is a superconductor when $\sigma = \infty$.

TELEGRAPHER'S EQUATIONS. A set of differential equations that describes a one-dimensional transmission line.

TLM OR TRANSMISSION LINE MATRIX METHOD.
The space of a transmission line structure is divided into a mesh of transmission lines excited with a Dirac pulse. An equivalence between this and the solution to Maxwell's equations is established.

TRANSVERSE RESONANCE TECHNIQUE. The line or discontinuity is modeled as a resonant waveguide structure. The resonance is used to calculate

the propagation constants of the structure.

QUASI-STATIC. Quasi-static analyses are carried out with simplifications that are only strictly valid at dc; however, the ac (time-varying) solution is found to be similar.

WAVEGUIDE MODEL. The actual structure is mapped into an equivalent waveguide structure. This waveguide is then analyzed with existing techniques. Finally this is converted back to the original structure, which is the desired solution.

W_{EFFECTIVE}. The effective trace width taking into account nonzero trace thickness fringing fields; the width of a parallel-plate waveguide having the same thickness dielectric, impedance, and propagation velocity.

Z₀. The characteristic impedance.

Appendix A

List of Symbols, Acronyms, and Abbreviations

α	attenuation constant	δ	skin depth
α_c	attenuation due to conductor losses	e	base of natural logarithm, 2.718281828...
α_d	attenuation due to dielectric losses	$E(k, \phi)$	elliptic integral of the second kind
α_r	attenuation due to radiative losses	ϵ	$\epsilon = \epsilon_0 \epsilon_r$, permittivity
B	susceptance	ϵ_0	permittivity of free space, 8.854183×10^{-12} F / m
β	phase constant (rad / m)	ϵ_{eff}	relative effective permittivity
b	dielectric thickness, used with stripline	ϵ_r	relative permittivity, dielectric constant
$bei_{\nu} x$	Kelvin function	F	farad
$ber_{\nu} x$	Kelvin function	f	frequency
C	capacitance	$F(k, \phi)$	elliptic integral of the first kind
C'	capacitance per length, incremental capacitance	g	gap width
c, c_0	speed of light, 2.99792456×10^8 m / s	G	conductance, S
CPW	coplanar waveguide	G'	incremental conductance, conductance per unit length
CPWG	coplanar waveguide with ground plane	$G(B, \xi)$	Hencken function
$Ct(x, y)$	large radial cotangent function	Γ	reflection coefficient
D	outer (or larger) diameter	$\Gamma(z)$	gamma function
d	inner (or smaller) diameter	γ	propagation constant
		H	henry

$H_v(x)$	Bessel function of the third kind of order v , Hankel function	λ_0	wavelength in free space
h	height, the dielectric thickness	\ln	logarithm base e , also known as natural or Napierian log
η	wave impedance, Ω	\log	logarithm base 10, also known as common logarithm
η_0	characteristic impedance of free space, $120 \pi \Omega$	m	meter
j	$\sqrt{-1}$	θ	electrical length
$j_{v,s}$	the s th zero of the Bessel function of the first kind of order v	μ	permeability
$j'_{v,s}$	the s th zero of the derivative of the Bessel function of the first kind of order v	μ_0	permeability of free space, $4 \pi \times 10^{-7} \text{ H/m}$
$J_v(x)$	Bessel function of the first kind of order v	μ_r	relative permeability
k	coupling	π	the ratio of a circle's circumference to its diameter, 3.14159265...
k	modulus of elliptic functions	q	filling factor
k'	complementary modulus of elliptic functions	θ	electrical length
$K(k)$	complete elliptic integral of the first kind	ρ	resistivity
$K(k')$	complete elliptic integral of the first kind and complementary modulus	ρ_{Cu}	resistivity of copper, $1.72 \times 10^{-8} \Omega\text{-m}$
$kei_v x$	Kelvin function	R	resistance
$ker_v x$	Kelvin function	R'	incremental resistance, resistance per unit length
L	inductance	R_s	surface resistivity
L'	inductance per length, incremental inductance	s	seconds
l	length	s	spacing
λ	wavelength	$\text{Si}(x)$	sine integral function
λ_g	wavelength in the guide	σ	conductivity
		t	thickness
		$\tan \delta$	tan delta, loss tan, loss tangent
		v_p	propagation velocity
		ω	angular frequency, $\omega = 2 \pi f$

W	watts
w	width
w_{eff}	effective width taking into account thickness or frequency
X	reactance
$y_{v,s}$	the s th zero of the Bessel function of the second kind of order v
$y'_{v,s}$	the s th zero of the derivative of the Bessel function of the second kind of order v
Y	admittance
$Y_v(x)$	Bessel function of the second kind of order v , Weber's function
$Y'_v(x)$	derivative of the Bessel function of the second kind of order v , Weber's function
$z_{v,s}$	the s th zero of the combined Bessel function equation of order v
$z'_{v,s}$	the s th zero of combined Bessel function derivatives of the first kind of order v
Z	impedance
Z_0	characteristic impedance
$Z_{0,e}$	even mode characteristic impedance
$Z_{0,o}$	odd mode characteristic impedance
Ω	ohm

Appendix B

Common Questions and

Answers

Does surface roughness matter?

At frequencies where the skin depth is of the same order as the surface roughness, the roughness will result in increased losses. Rolled copper (smoother) can be used instead of electroplated copper when this is important.

How close can a shield come to trace top before affecting Z_0 ?

Shielding lowers ϵ_{eff} and Z_0 and creates resonances. Neglect the effects of shield for $W_{wall} > 5 \times w$ and $h_2 > 5 \times h$. Equations for enclosure resonances agree pretty well with experience. When resonances occur, the box must be partitioned or the cavity loaded with a dissipative material.

What is the coupling between two traces? How far apart should they be spaced for no coupling?

Unless one conductor is completely shielded, it is impossible to have 0 dB of coupling. To calculate the coupling, use the equations of Chapter 4 for the correct situation. These will give the signal level in the second line relative to the first line (sometimes called the "backward coupling").

What happens when microstrip line is embedded? (v_p , capacitance, Z_0)

The capacitance increases, lowering v_p and Z_0 . To calculate this effect, see Chapter 3.

When should one use microstrip, when should one use stripline?

The advantage of stripline is that it almost completely confines the fields between two ground planes. This reduces coupling to adjacent traces and makes it less sensitive to packaging. On the negative side, the cost is increased and vias must be provided to connect components to the center conductor.

When does one use coplanar waveguide?

Coplanar waveguide is useful when components need to be mounted from the transmission path to ground. Plated-through holes to the other side of the PCB are not required. Coplanar waveguide also allows the trace to be narrowed to the width of a device's leads while keeping the impedance constant.

How does one choose a dielectric thickness for multilayer boards?

Determine the number of layers from the density. Programs are available to assist with this calculation. Determine the number of signal, ground, and power planes needed. Determine the overall board thickness required for the card cage and connectors.

Now the goal is to create a symmetric layout to minimize warpage. This means that the planes and signal layers are placed symmetrically on either side of the PCB center line. Corresponding planes on opposite sides of the center line should have similar amounts of metallization after etching. The dielectric thicknesses should also be symmetric about the center line. To minimize the effects of tolerances, controlled impedance signal layers should have C stage dielectric layers as thick as possible in a standard size. The other layers can be thinned to accomplish this down to the limitations tabulated in Chapter 9.

How does one choose a dielectric thickness for two-sided boards?

The main consideration here will often be the physical requirements of the card cage and connectors. The thicker the dielectric, the less sensitive you will be to tolerances in the manufacturing process. Thinner dielectrics reduce the trace dimensions for a given impedance and may make it easier avoid stray couplings.

How much effect does the trace thickness (inner layer vs. outer layer) have on the conductor width? (Z_0)?

To answer this, calculations were made using the microstrip line equations of Chapter 3 for 0.0625 in thick G-10 at 1.6 GHz. As shown in the table below, the line width for 50 Ω changes about 1% when the thickness is changed from 1 oz. to 2 oz. The trace length for an electrical 90° changed about 0.04% due to the change in ϵ_{eff} . This translates into about 0.03° or 0.1 ps length difference at 1.6 GHz over 1 in of length.

Trace Thickness (mil)	50 Ω Width (in)	90° length (in)
1.4	0.11984	1.01375
2.8	0.11878	1.01411
4.2	0.11785	1.01445
5.6	0.11699	1.01476

What happens at a T or crossing of lines?

The model is a series inductance in each arm and a shunt capacitance at their junction.

Why use a dielectric other than G-10 or FR-4?

The main reasons are loss characteristics, impedance tolerance, frequency range, and feature size. At high frequencies the loss ($\tan \delta$) of G-10 will cause the transmission paths to be lossy, which increases insertion losses and lowers the Q achievable in the circuit. PTFE dielectrics tend to have a more uniform dielectric constant and etch characteristics (at a greater cost) as well as lower loss. Dielectrics with ϵ_r 's greater than that of G-10 will have narrower lines for a given impedance resulting in less radiation, better match to device leads, and more dense circuitry.

What is the difference between G-10 and FR-4?

FR-4 is a fire-resistant version of G-10.

What effect does solder mask have on microstripline impedance?

Solder mask increases the capacitance and lowers v_p . See Chapters 3 and 9 for more information.

How far from a transmission line does the dielectric need to extend to guarantee Z_0 ?

Referring to the figure for microstripline with truncated dielectric (Figure 3.5.2.1): $T/w > 1.0$ will result in insignificant changes in the propagation characteristics of the line (Z_0 , v_p , α , etc.).

For striplines, ground extending beyond the trace $> 0.8h$ (for $w/h = 0.6$) will cause $< 0.2\%$ error in trace capacitance¹. Figure 7 of the reference indicates that extending the dielectric beyond the trace $> 3.0h$ (for $w/h = 0.6$) will give $< 0.2\%$ error in trace capacitance.

This effect is also frequency dependent. As frequency increases the fields concentrate more and more into the dielectric directly below the signal trace.

What are even and odd mode impedances?

There is only the characteristic impedance of a line whenever there are only two conductors involved; however, when there are three the signal can propagate in two different ways. The even mode corresponds to two of the conductors being at the same potential relative to the third ("common mode"). The odd mode corresponds to the two conductors having opposite potentials ("differential mode"). Usually one or the other is desired and is driven by the source. Source or load imperfections or discontinuities can cause the other mode to propagate.

What is the best way to make a constant impedance bend in a trace?

A mitered bend will take up less room than an equivalent VSWR rounded corner bend. Use Chapter 4 to calculate the required miter and the required minimum bend radius.

What should I do to achieve a very precise Z_0 ? What accuracy can be expected?

For tightest control over the characteristic impedance of a line:

- Use as thin a trace as possible. This minimizes etching tolerances.

¹Figure 6 of Ahmadouche, A., and J. Chilo, "Optimum Computation of Capacitance Coefficients of Multilevel Interconnecting Lines for Advanced Package," *IEEE Transactions on Components, Hybrids, and Manufacturing Technology*, Vol. 12, No. 1, March 1989, pp. 124–129.

- Keep controlled Z traces in the center of the PCB where etching is most uniform.
- Keep controlled Z traces on PCB inner layers. These layers are not plated and this eliminates tolerances caused by the plating and other processes.
- Use as thick a dielectric (and as wide a trace) as possible.
- Specify that you want controlled Z on the PCB fabrication documentation.
- Use a standard C stage dielectric thickness. Calculate line widths using the dielectric constant of the actual thickness dielectric used (applies to woven dielectrics).
- Include the effect of any overlaid dielectric. This includes solder mask and *skim coats*.
- Keep traces separated from each other and any adjacent ground plane by several line widths.
- Follow controlled impedance bend guidelines.
- Use as low a dielectric constant as possible (do not forget air!).
- Add thieves to keep trace density constant over the surface. This allows etching to be even over the board's surface.

Why 50 Ω ? Couldn't I use 60 Ω ?

The value 50 Ω comes from a trade-off in coaxial lines that was made between attenuation (loss) and power-handling capability. The maximum power point is at 30 Ω and the minimum attenuation occurs at 77 Ω . From these came the 75 Ω used for TV cabling and the compromise value of 50 Ω .¹ For maximum transmission of power from a source to a load, the source, the load, and the line should have the same impedance.

Some references give 4.3 and others give 4.8 as the dielectric constant of FR-4. Which is correct?

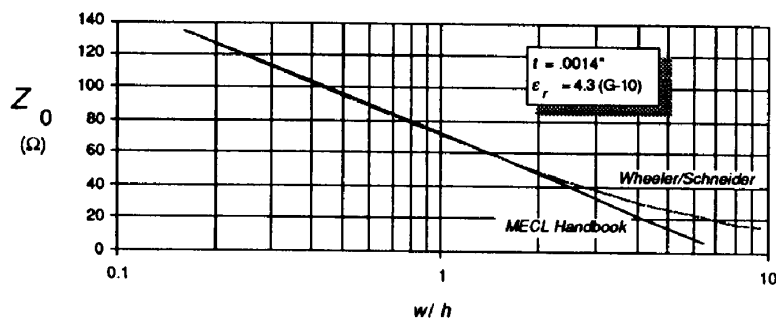
The amount of epoxy in the dielectric affects the dielectric constant. Thinner dielectrics tend to have a greater percentage of epoxy than thicker dielectrics. For thin dielectrics (less than ~30 mils), 4.3 is a good approximation to the permittivity. For thicker

¹Laverghetta, Thomas S., *Practical Microwaves*, Howard W. Sams, Indianapolis, IN, 1984.

dielectrics, 4.8 is a better number. See Chapter 9 for more detailed information.

How accurate are the MECL handbook equations relative to the Wheeler-Schneider ones?

As the graph below shows for 0.0625 in G-10 and 0.0014 in thick traces, the MECL Handbook equations are most accurate for impedances near and above 50 Ω :



MECL Handbook Compared to Wheeler-Schneider Equations

Appendix C

VSWR, Reflection Coefficient, and Impedance

There are a number of ways to express the impedance of a circuit. The terms VSWR, reflection coefficient, and impedance are all interrelated and convey roughly the same information.

The **voltage standing wave ratio** is physically the ratio of the maximum voltage along a line to the minimum. It can actually be measured this way—e.g., when we use a detector probe in the field of a slotted line. These maximum and minimum voltages along the line can also be written in terms of the waves incident on and reflected from the load:

$$VSWR \equiv \frac{E_{max}}{E_{min}} = \frac{E_{inc} + E_{refl}}{E_{inc} - E_{refl}} \quad (C.1)$$

The **reflection coefficient**, ρ , is also related to the incident and reflected waves:

$$\Gamma = \rho \equiv \frac{E_-}{E_+} = \frac{Z_L - Z_0}{Z_L + Z_0} \text{ unitless} \quad (C.2)$$

Similarly, we define the **transmission coefficient**:

$$\tau \equiv \frac{E_L}{E_+} = \frac{2 Z_0}{Z_L + Z_0} \quad (C.3)$$

The reflection and transmission coefficients have both magnitude and phase. Note that

$$\tau - \rho = 1 \quad (C.4)$$

We can now write the VSWR in terms of the reflection coefficient:

$$VSWR = \frac{1.0 + |\rho|}{1.0 - |\rho|} \quad (\text{unitless}) \quad (C.5)$$

The return loss is the ratio of the reflected power to the incident power:

$$\text{return loss} = -20 \log_{10} |\rho| \quad (\text{dB}) \quad (C.6a)$$

$$= -20 \log_{10} \frac{VSWR - 1}{VSWR + 1} \quad (C.6b)$$

In the above equations Z_0 is the characteristic impedance of the system.

EXAMPLE 1: PERFECT MATCH

For a perfect match, $Z_L = Z_0$ and we have no reflected power, so using (C1), (C2), and (C6):

$$\rho = 0$$

$$VSWR = 1.0:1$$

$$\text{return loss} = \infty$$

EXAMPLE 2: SHORT CIRCUIT

For a short circuit, $Z_L = 0$ and all of the power is reflected, so

$$\rho = -1$$

$$VSWR = \infty:1$$

$$\text{return loss} = 0 \text{ dB}$$

EXAMPLE 3: OPEN CIRCUIT

For an open circuit, $Z_L = \infty$ and all of the power is again reflected, so

$$\rho = +1$$

$$VSWR = \infty:1$$

$$\text{return loss} = 0 \text{ dB}$$

Table C.1 summarizes the relationships for a range of cases

Z_L	ρ	Return Loss	VSWR
0	-1.00	0 dB	$\infty:1$
.5 Z_0	-.33	9.6	2.0:1
1.0 Z_0	0.00	∞	1.0:1
1.1 Z_0	+.05	26.4	1.1:1
2 Z_0	+.33	9.6	2.0:1
10 Z_0	+.82	1.7	10.0:1
100 Z_0	+.98	.2	100:1
∞	+1.00	0	$\infty:1$

Table C.1: Impedance, ρ , Return Loss, and VSWR

Some other useful relations are:

$$|\rho| = 10^{-\text{return loss}/20} \quad (C.7)$$

$$VSWR = \frac{1 + 10^{-\text{return loss}/20}}{1 - 10^{-\text{return loss}/20}} = \begin{cases} \frac{Z_L}{Z_0} & \text{iff } Z_0 < Z_L \\ \frac{Z_0}{Z_L} & \text{iff } Z_0 \geq Z_L \end{cases} \quad (C.8)$$

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